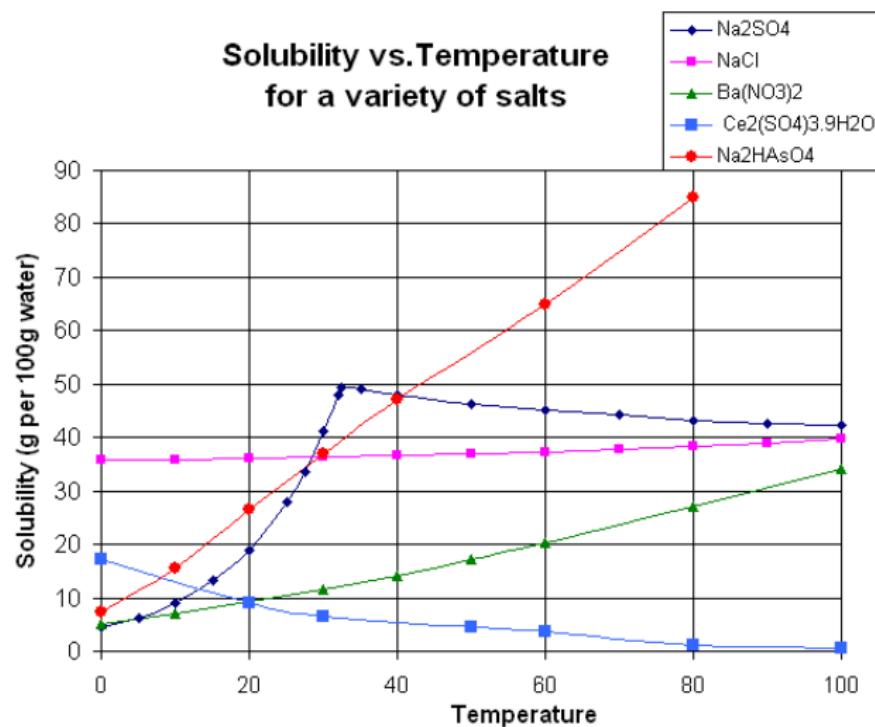


Flavored Freezeout

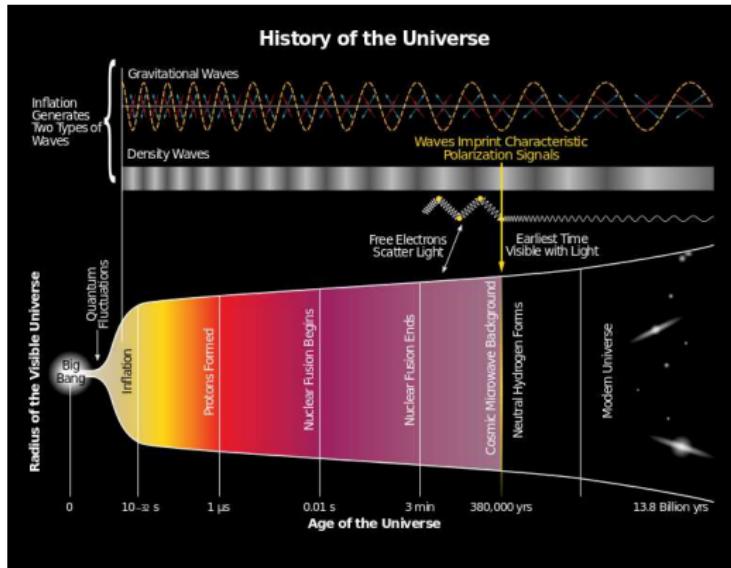
Sandeep Chatterjee (sandeepc@vecc.gov.in)
VECC, Kolkata

QCD at High Density, TIFR
27-30 January, 2015

Multiple Freezeout on your table top: Salt mixture in water



Multiple Freezeout in the Early Universe



Hypothetical Q: A more rapid expansion \rightarrow sudden single freezeout a la HIC ?

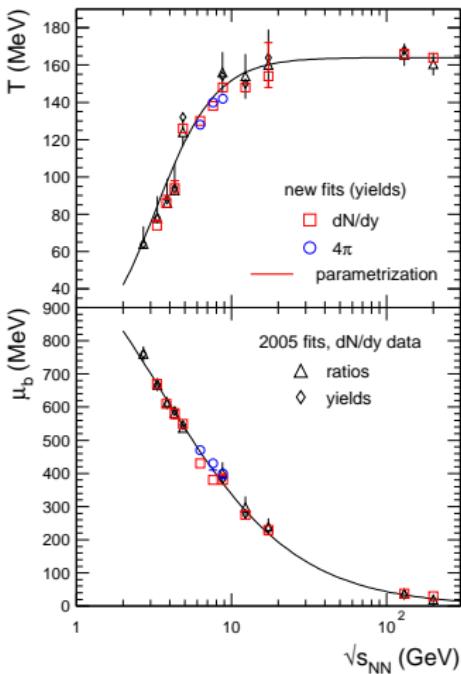
Freezeout in HIC

- Freezeout is a result of competition between 2 effects: constituent interactions and fireball expansion- Cross section vs Dilution
- In the late stage of a heavy ion collision (HIC), the rate of collisions between the constituents can no longer cope with the expansion rate. As a result, hadrons start freezing out.
- Simple assumption: All strong interaction rates are same. Hence single chemical freezeout (1CFO).

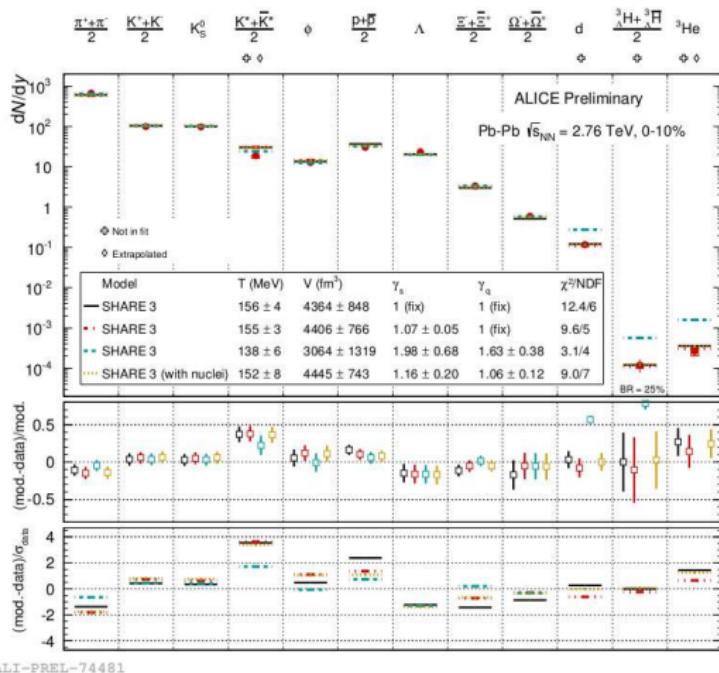
Single Chemical Freezeout: 1CFO

Standard practice:
All the hadrons CFO at
the same (T, μ_B) surface.
This provides an overall
good qualitative picture of
CFO at $\sqrt{s_{NN}} \sim 2 - 200$
GeV with ~ 4 params.

Andronic et al: 0812. 1186



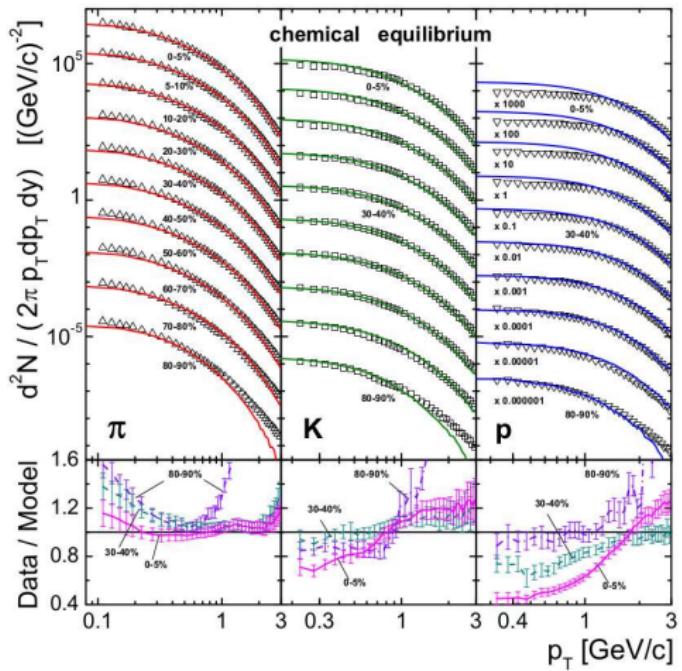
1CFO at LHC



no. of parameters 4+1 (μ_B set to 0 by hand.)

Floris: 1408.6403

1CFO at LHC



Begun et al: 1405.7252

Revisiting our 1CFO assumption: When does chemistry freeze out?

Basic observables are the spectra of identified particles; from this one gets yields. Relative yields of hadrons is the outcome of “chemistry”.

At early times, fireball is a reactive fluid. Reaction rates depend on local densities as well as rates of mixing.

Kinetic equations more complicated: need numerical treatment, but many parameters fed into the code.

When does isospin freeze out?

- The rates for processes $p + \pi^- \leftrightarrow n + \pi^0$, remain high at $\simeq 150$ MeV, because $m_n - m_p$ is small and the yield of pions is large. So the chemical freezeout of baryon isospin can be delayed. The $p \leftrightarrow n$ reaction may proceed without suppression right up to kinetic freezeout

Asakawa, Kitazawa, 2011

Can the K and π freeze separately?

- Indirect transmutations of K and π involve strange baryons in reactions such as $\Omega^- + K^+ \leftrightarrow \Xi^0 + \pi^0$. These have very high activation thresholds. There is no physics forcing K and π to freezeout together. But K and ϕ are resonantly coupled, so freeze out together.

SC, Godbole, Gupta, 2013

Double Chemical Freezeout: 2CFO

- 'Isospin changing' reactions are last to freezeout
 $(p + \pi^0 \leftrightarrow n + \pi^+)$ (Asakawa, Kitazawa 2011)
 - low activation energy
 - high pion density
- 'Strangeness changing' reactions can freezeout earlier
 $(\Omega^- + K^+ \leftrightarrow \Xi^0 + \pi^0)$
 - High activation energy
 - Ω and K densities much less compared to that of π ;
 $\Omega^- + K^+$ reactions much suppressed
- Motivates to propose separate CFO for (strange+hidden strangeness) and non strange hadrons: 2CFO
- T_s, V_s, μ_{B_s} characterise the strange surface
- $T_{ns}, V_{ns}, \mu_{B_{ns}}$ characterise the non-strange surface

Hadron Yields in Thermal Model

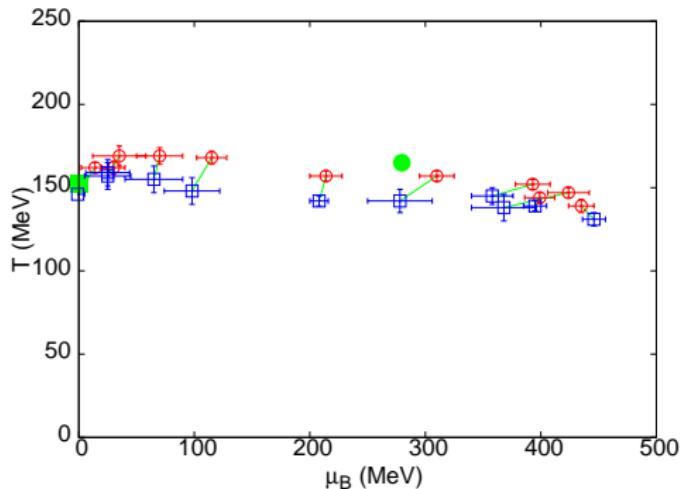
- The ideal hadron resonance gas (HRG) partition function Z in the grand canonical ensemble at the time of CFO at a particular beam energy $\sqrt{s_{\text{NN}}}$ is given as

$$\log [Z(\sqrt{s_{\text{NN}}})] = \sum_i \log [Z_i(T_i(\sqrt{s_{\text{NN}}}), \mu_i(\sqrt{s_{\text{NN}}}), V_i(\sqrt{s_{\text{NN}}}))]$$

-

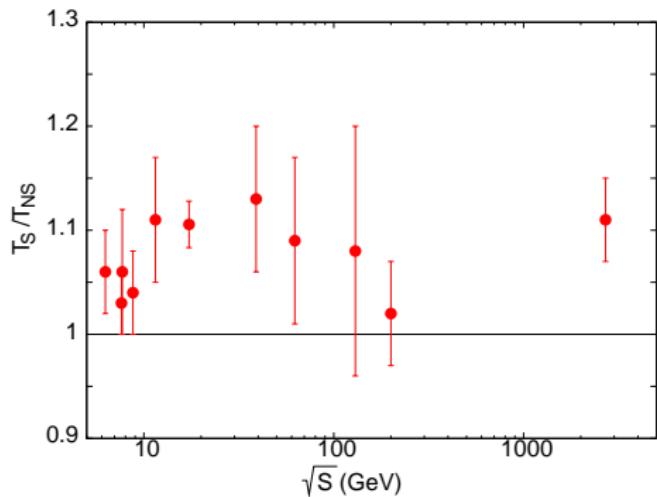
$$\begin{aligned} N_i^p &= \frac{\partial}{\partial \left(\frac{\mu_i}{T_i} \right)} \log [Z] \\ &= \frac{V_i T_i}{\pi^2} g_i m_i^2 \sum_{l=1}^{\infty} (-a)^{l+1} l^{-1} K_2(l m_i / T_i) \times \\ &\quad \exp(l(B_i \mu_{B_i} + Q_i \mu_{Q_i} + S_i \mu_{S_i}) / T_i) \end{aligned}$$

2CFO Freezeout Parameters



SC, Godbole, Gupta: 1306.2006

2CFO Freezeout Parameters



SC, Godbole, Gupta: 1306.2006

Nuclei Yields in Thermal Model

- Treat the nuclei on the same footing as other hadrons
- Thus the Boltzmann factor decides the nuclei yield:
Braun-Munzinger et al, 1996

Nuclei Yields in Simple Coalescence

- Nuclei formed by coalescence of hadrons near the KFO surface

$$E_A \frac{d^3 N_A}{d^3 P_A} = B_A \left(E_p \frac{d^3 N_p}{d^3 P_p} \right)^Z \left(E_n \frac{d^3 N_n}{d^3 P_n} \right)^{A-Z}$$

- At the level of yields

$$N_A = C_A (N_p)^Z (N_n)^{A-Z}$$

- For ratios of nuclei, for eg. anti-nuclei to nuclei

$$\begin{aligned} \frac{\bar{N}_A}{N_A} &= C_{\bar{A}A} (N_{\bar{p}}/N_p)^Z (N_{\bar{n}}/N_n)^{(A-Z)} \\ &\sim C_{\bar{A}A} (N_{\bar{p}}/N_p)^A \end{aligned}$$

Nuclei Yields in Simple Coalescence

- Nuclei production within the coalescence model is a combination of two distinct physics issues:
 - The physics of $C_{\bar{A}A}$. This is related to the correlation effects in the phase space that exist between the constituent hadrons at the time of the KFO
 - The abundances of the constituent hadrons at the time of KFO which is already fixed at the CFO surface obtained from fits to the hadron yields.
- Here we are interested in the role played by the latter physics in determining the nuclei yields. Hence take $C_{\bar{A}A} = 1$. Sufficient for our purpose to demonstrate the dependence of the nuclei production on the CFO scheme.

Ratios

- Unlike Flavor Ratio (R^{UF}):

$$\left(\frac{N_s^t}{N_{ns}^t} \right)^{\text{th}} = \exp \left(S \mu_S / T_s \right) \frac{g_s V_s}{g_{ns} V_{ns}} \left(\frac{T_s m_s}{T_{ns} m_{ns}} \right)^{3/2} \times \\ \exp \left(m_{ns} / T_{ns} - m_s / T_s \right) \times \\ \exp \left(\mu_{B_s} / T_s - \mu_{B_{ns}} / T_{ns} \right)$$

- Hence,

$$R_{\text{2CFO}}^{\text{UF}} \sim \left(\frac{T_s}{T_{ns}} \right)^{3/2} \left(\frac{V_s}{V_{ns}} \right) R_{\text{1CFO}}^{\text{UF}}$$

Ratios

- Like Flavor Ratio (R^{LF}):

$$N_i^t / N_j^t = \left(\frac{g_i}{g_j} \right) \left(\frac{m_i}{m_j} \right)^{3/2} \times \exp \left(\left((m_j - m_i) + (B_i - B_j) \mu_B \right) / T \right)$$

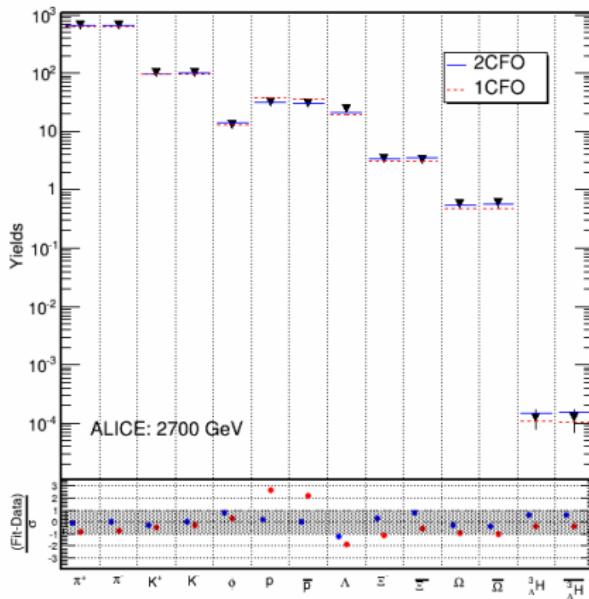
- Hence,

$$R_{2\text{CFO}}^{\text{LF}} \sim R_{1\text{CFO}}^{\text{LF}}$$

- Anti-particle to particle ratios simplifies even further.

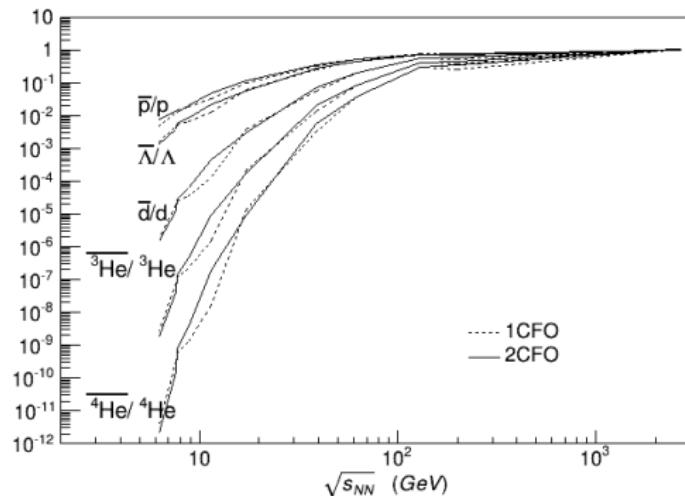
$$\left(\overline{N_i^t} / N_i^t \right)^{\text{th}} = \exp \left(-2 (B_i \mu_B + Q_i \mu_Q + S_i \mu_S) / T \right)$$

2CFO at LHC



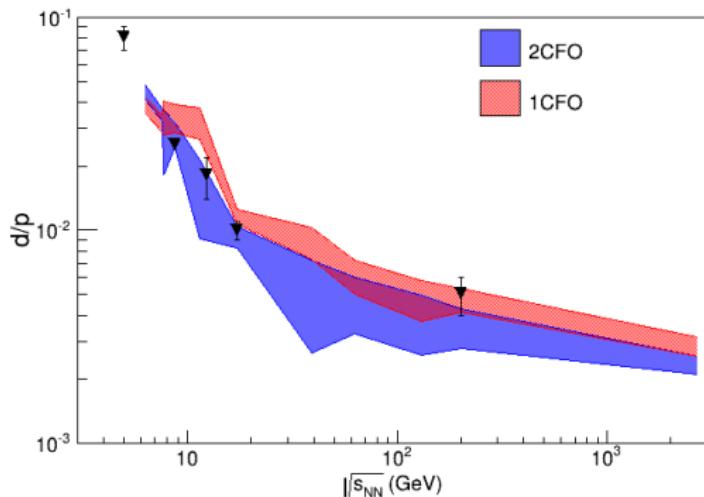
SC, Mohanty: 1405.2632

Antiparticle to Particle Ratio



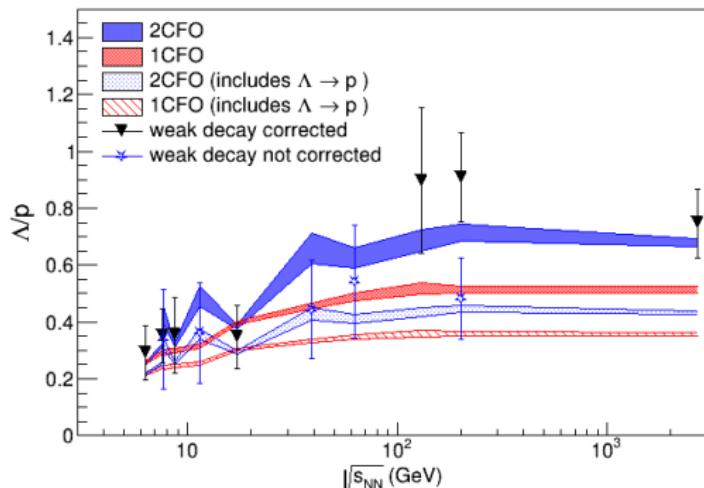
SC, Mohanty: 1405.2632

Like Flavor Ratio



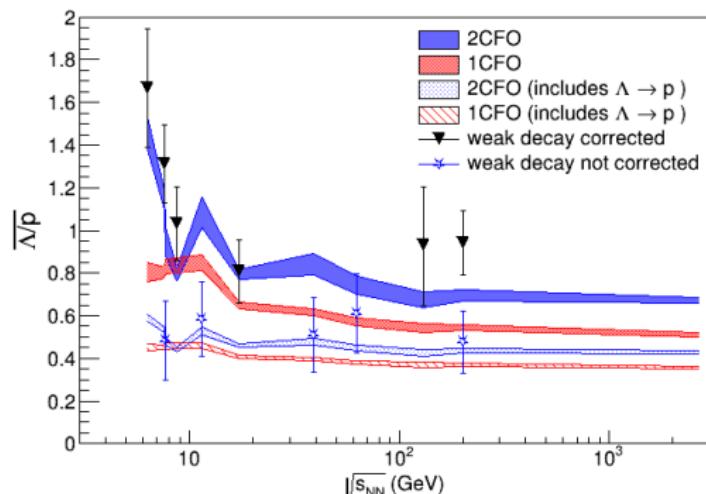
SC, Mohanty: 1405.2632

Unlike Flavor Ratio



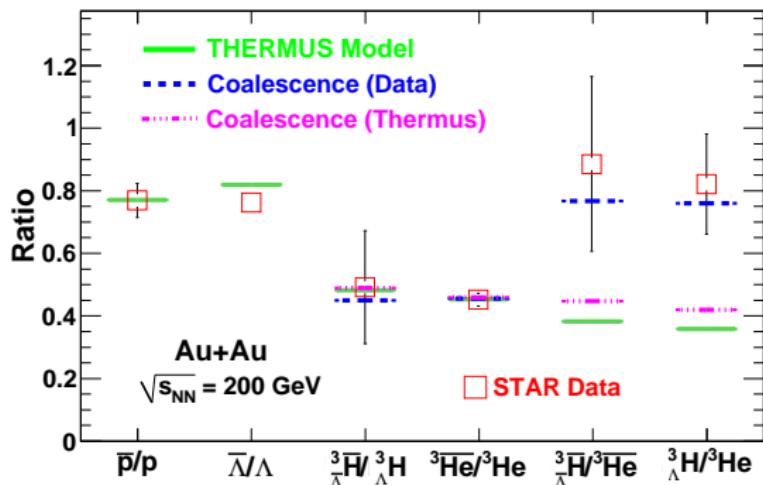
SC, Mohanty: 1405.2632

Unlike Flavor Ratio



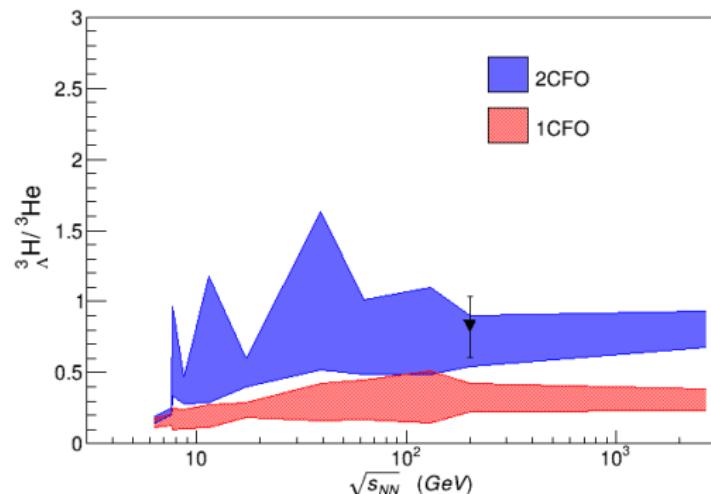
SC, Mohanty: 1405.2632

Unlike Flavor Ratio: Nuclei



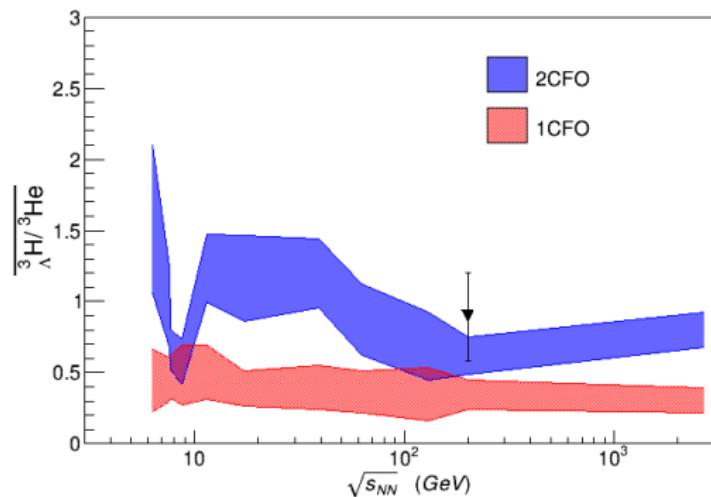
Andronic et al 2011, Cleymans et al 2011, Pal et al 2013

Unlike Flavor Ratio



SC, Mohanty: 1405.2632

Unlike Flavor Ratio



SC, Mohanty: 1405.2632

Other approaches

- Post 1CFO employ hadronic afterburner: Microscopic Transport Approach (UrQMD Model). Baryon-antibaryon annihilation main source of correction. Systematics ?- Steinheimer et al, 2013
- Introduce additional light and strange chemical non-equilibrium fugacity factors- Petran et al, 2013

Summarising..

- Multiple freezeout is a common occurrence in nature: from a cooling salt mixture in water to the cooling early universe. A multi-component system naturally freezes over a range in the relevant parameter space.
- Freezeout in the cooling fireball in HIC- Is the freezeout gradual enough to leave an imprint on the data ?
- 1CFO provides an overall good description of the hadrons(nuclei) yields across a wide range of $\sqrt{s_{\text{NN}}}$
- Does closer/careful inspection of the data reveal details in freezeout ? Which observables are most sensitive?
- Strange to non strange hadron/nuclei ratios are most sensitive to flavor dynamics at freezeout
- Anomaly with data of Λ/p at LHC, $^3\Lambda/\text{He}$ at top RHIC have a common origin: flavor dynamics at freezeout
- Influence of additional resonances ?- they will affect the above strange to non strange ratios. On including them, can the above anomalies with data be addressed within 1CFO ?
Need to check- require input on their branching ratios
- Data from FAIR, BES-II can throw more light

Take home

