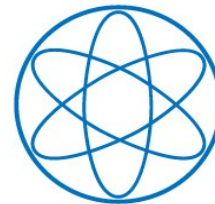


# Theoretical overview of dark matter

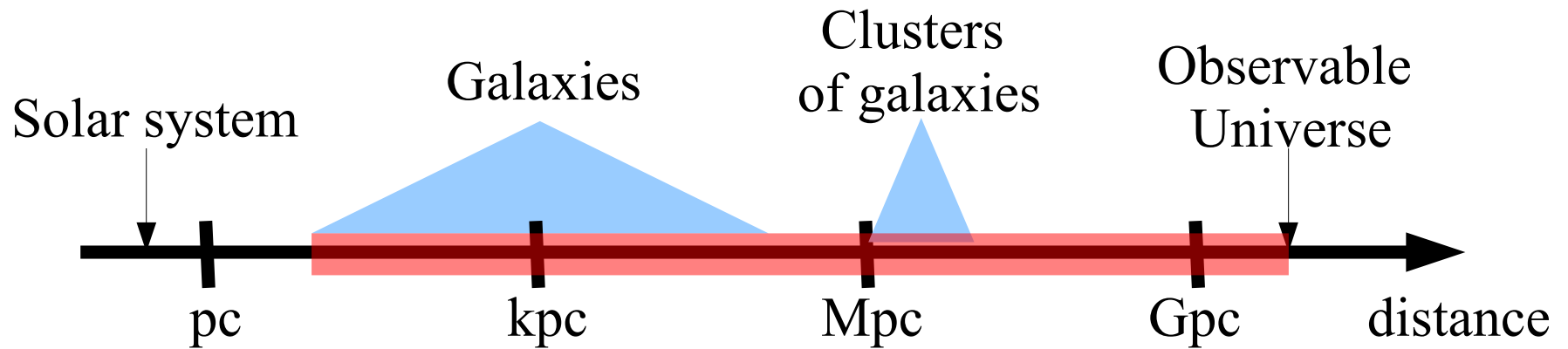
Alejandro Ibarra

Technische Universität München

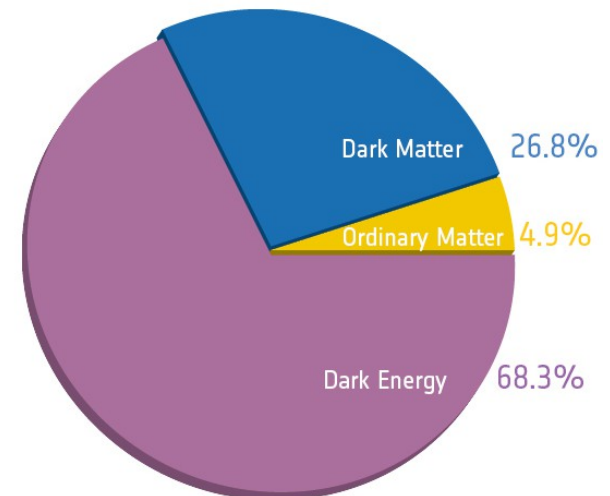
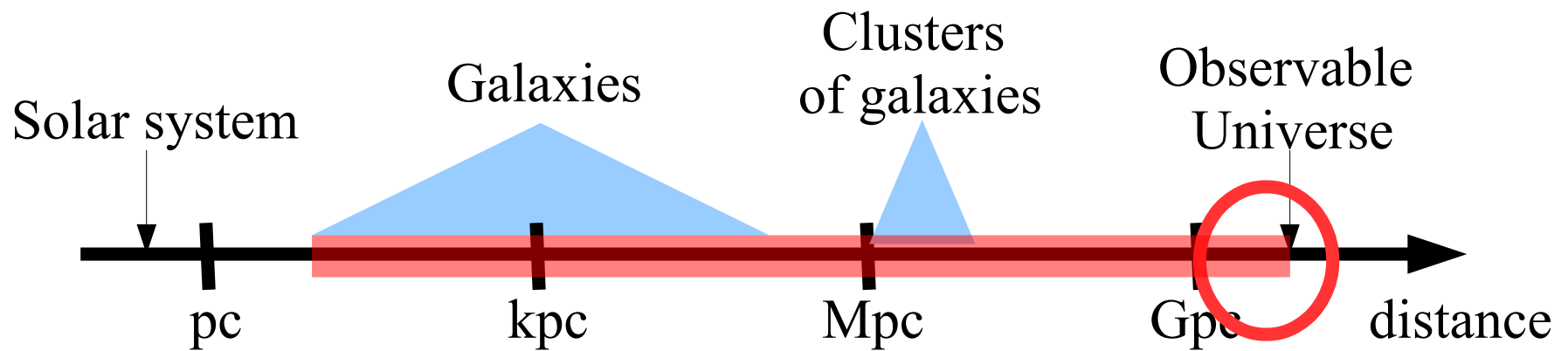


SUSY'17  
Mumbai  
14 December 2017

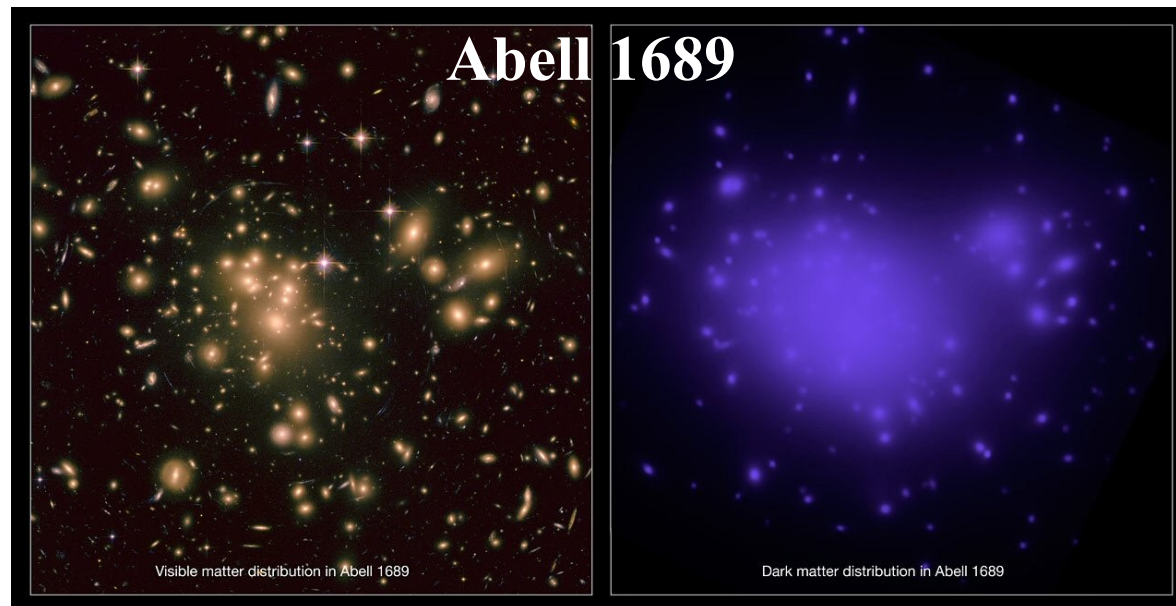
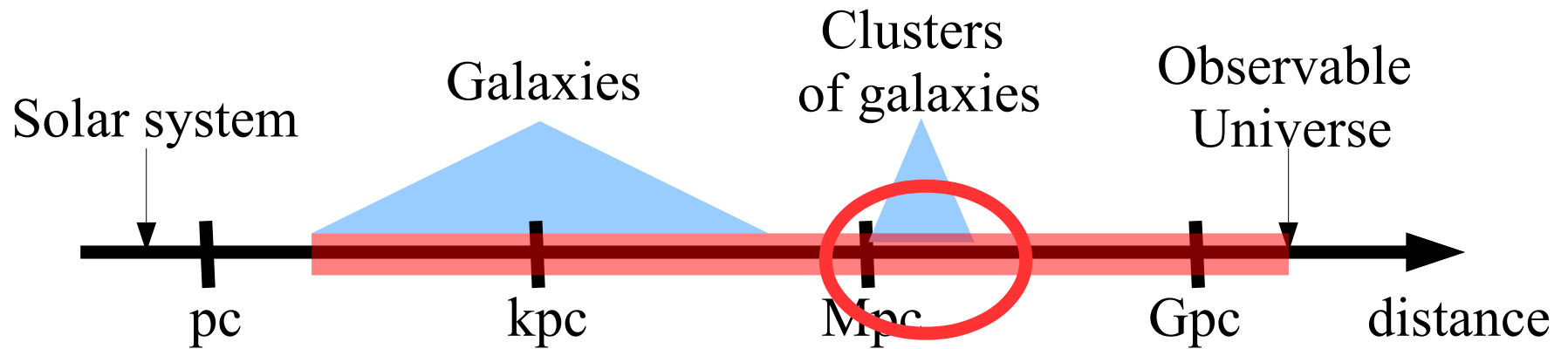
# There is evidence for dark matter in a wide range of distance scales



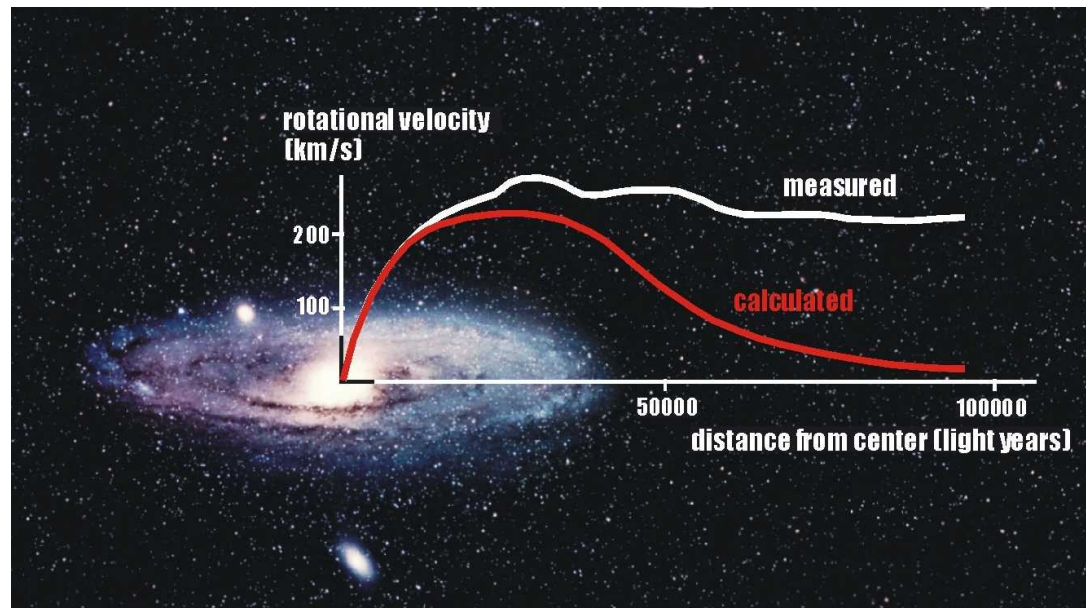
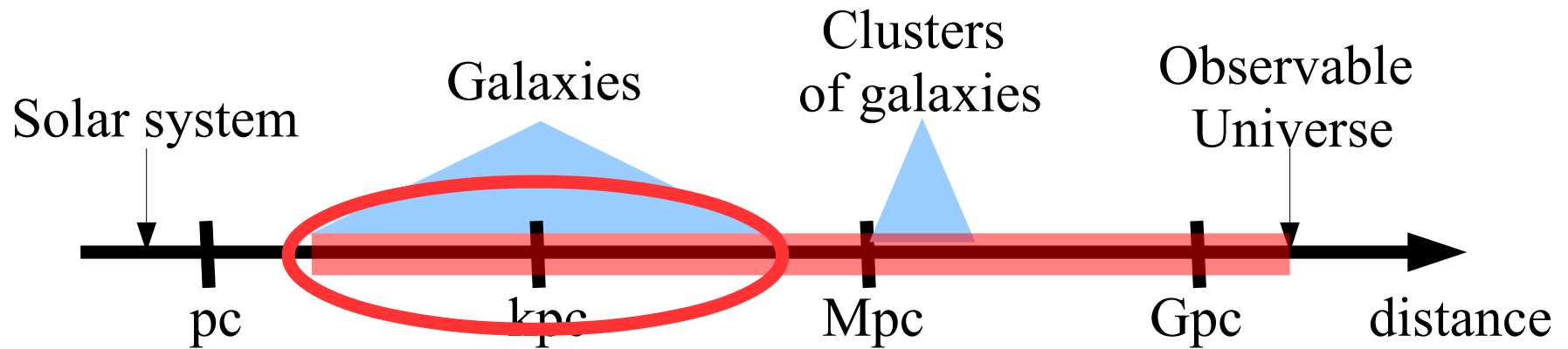
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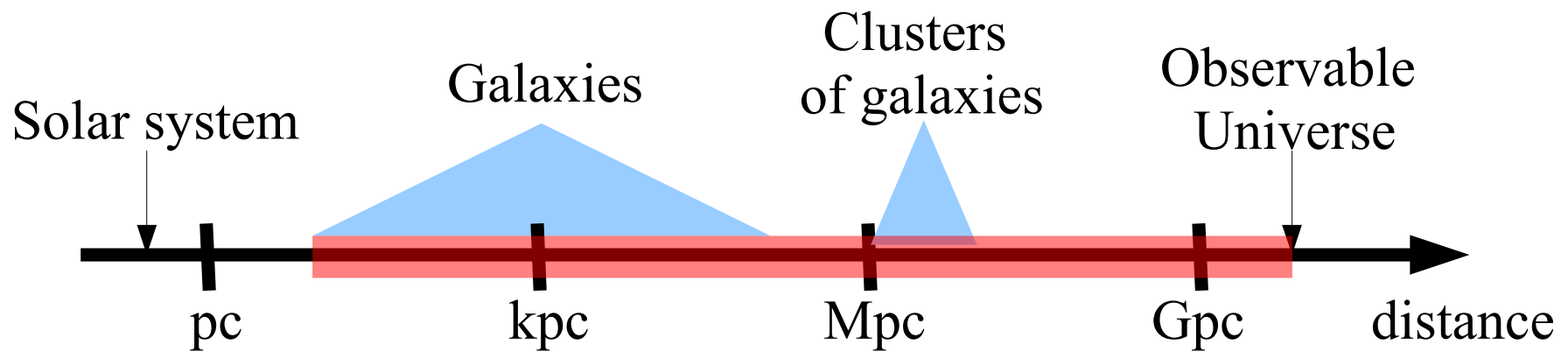
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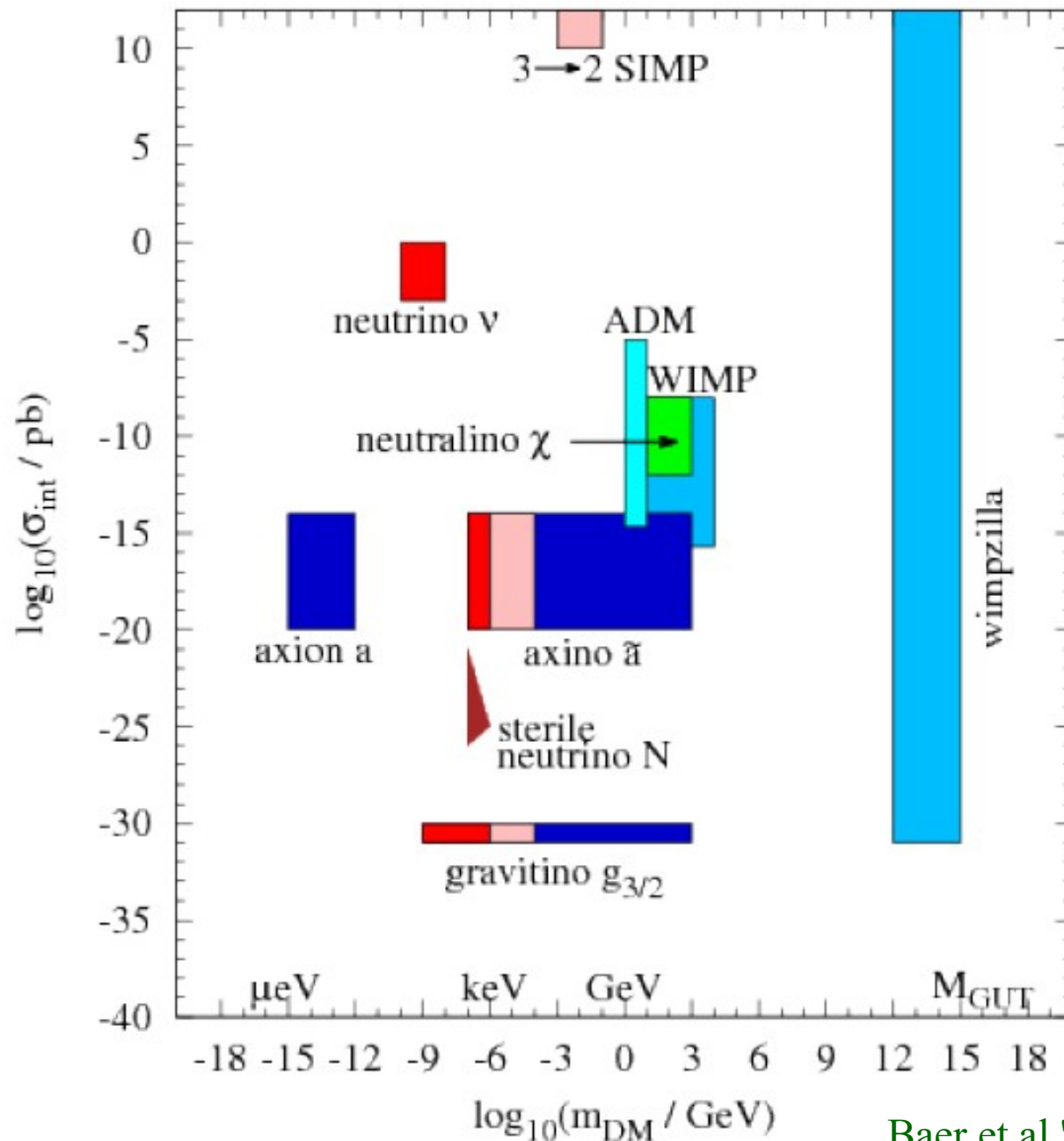


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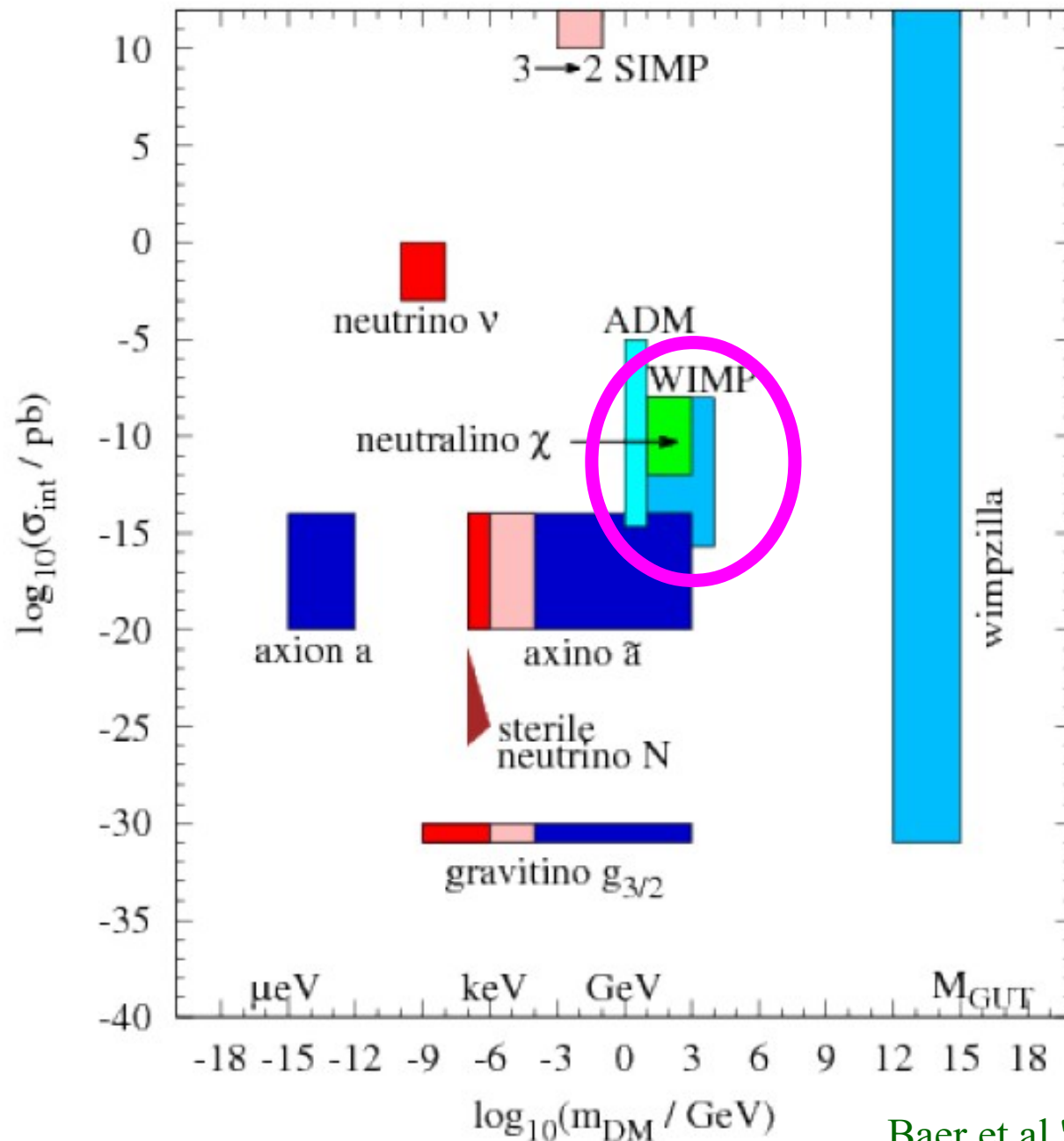


Cosmological and astronomical observations are consistent with the dark matter being constituted by new particles not contained in the Standard Model

# Candidates of particle dark matter

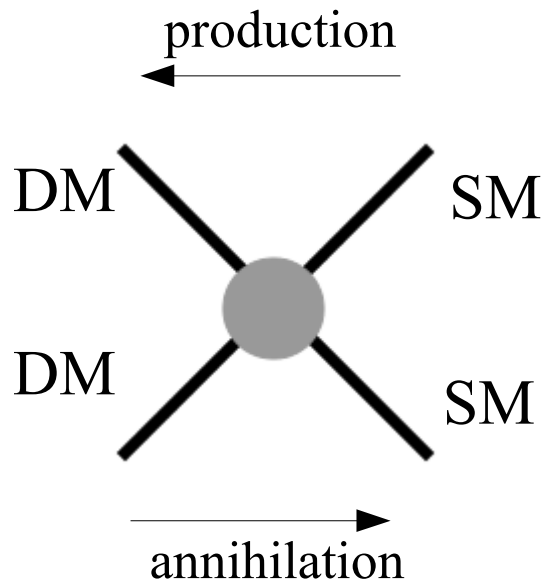


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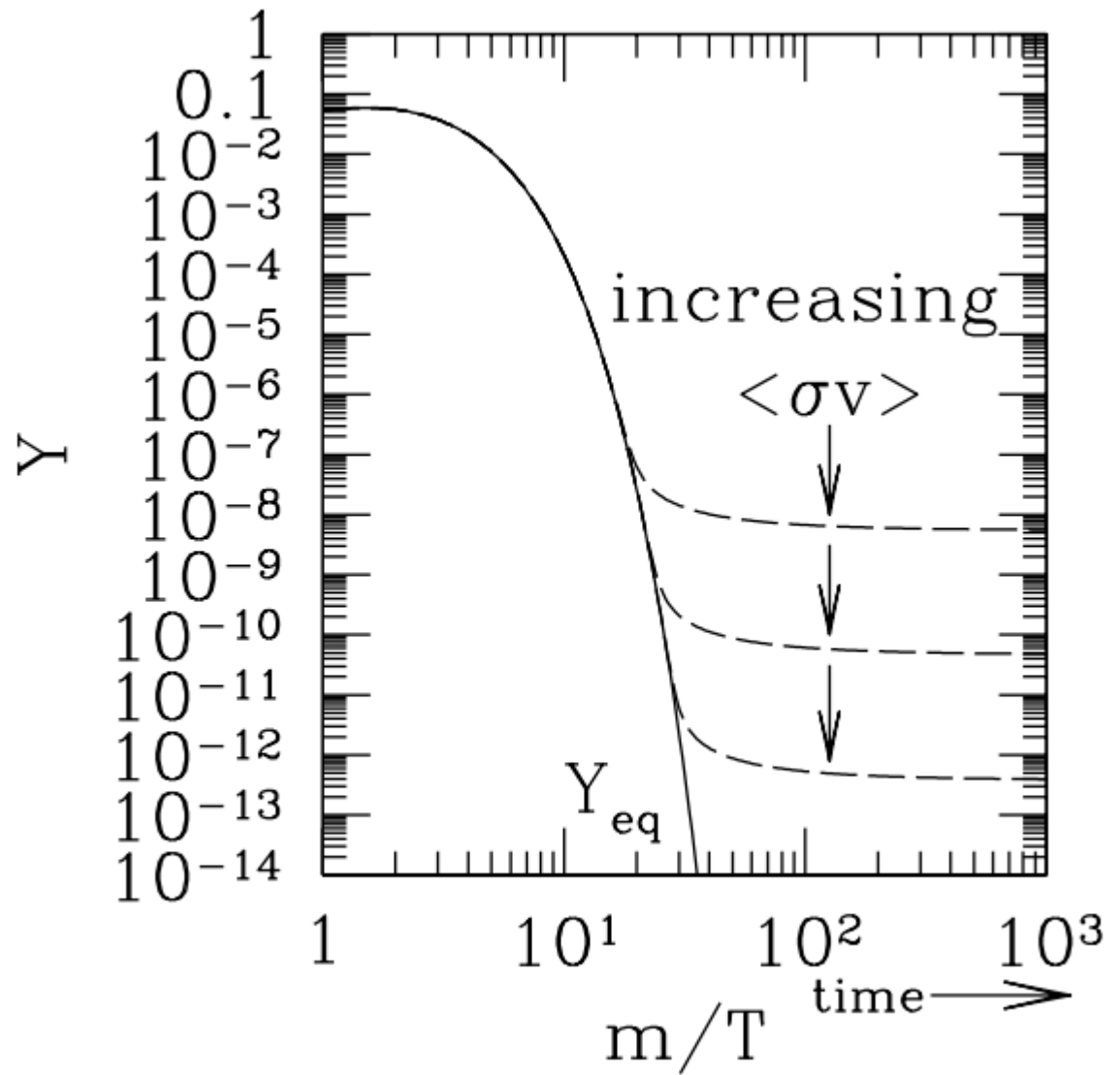
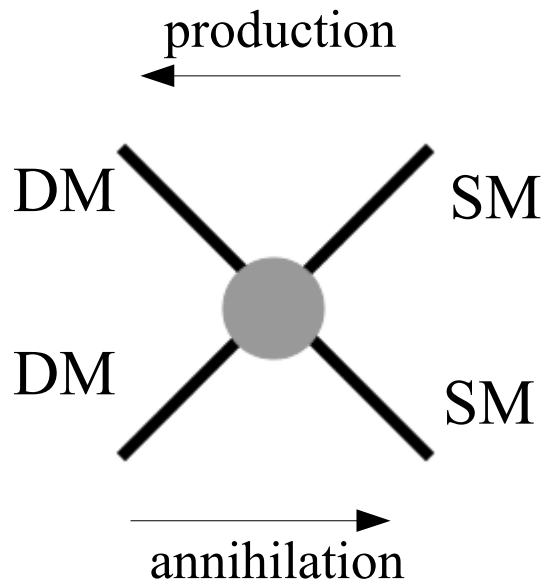




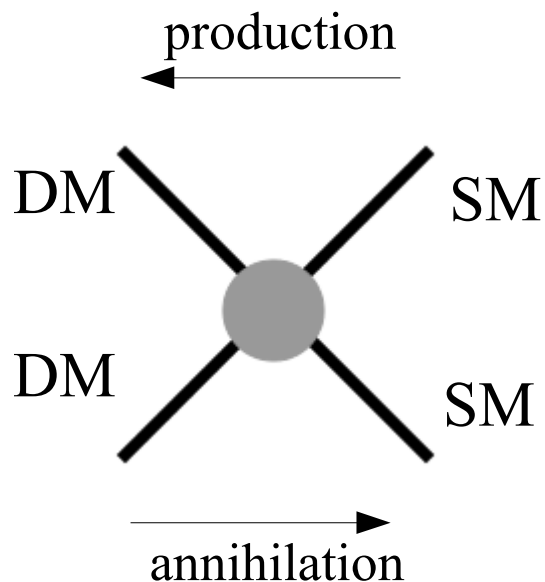
# WIMPs in the early Universe



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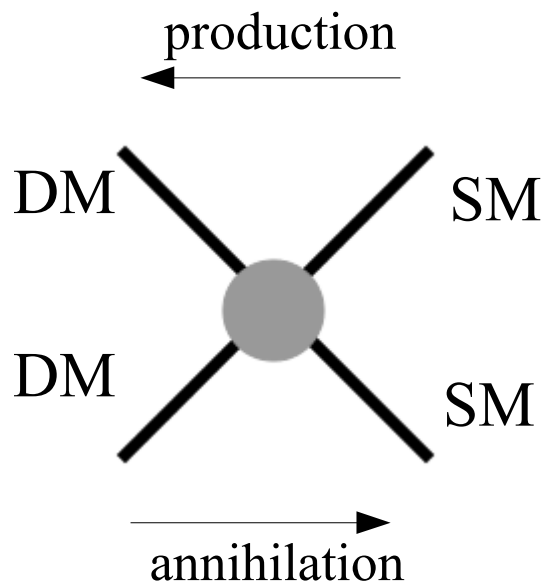
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Assuming that the dark matter particles were in thermal equilibrium with the SM in the Early Universe, their relic abundance reads:

$$\Omega h^2 \simeq \frac{3 \times 10^{-27} \text{ cm}^3 \text{ s}^{-1}}{\langle \sigma v \rangle}$$

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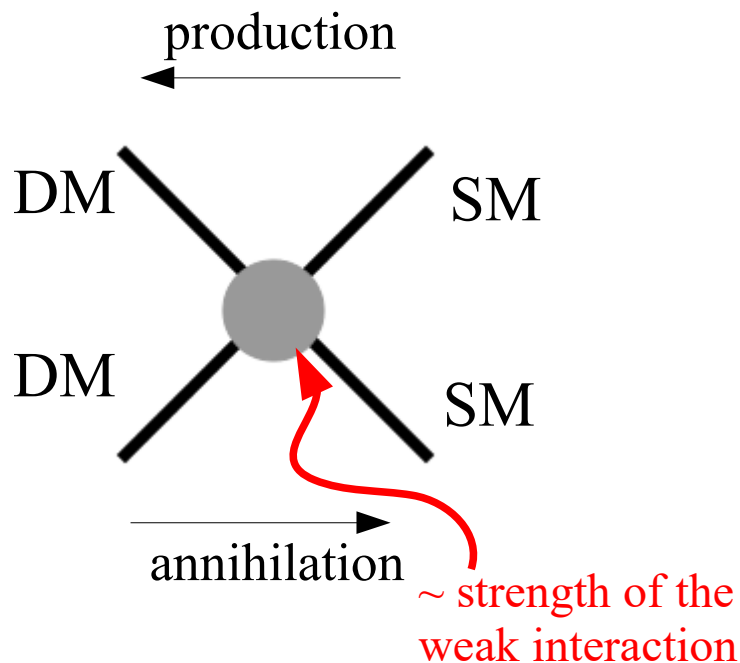
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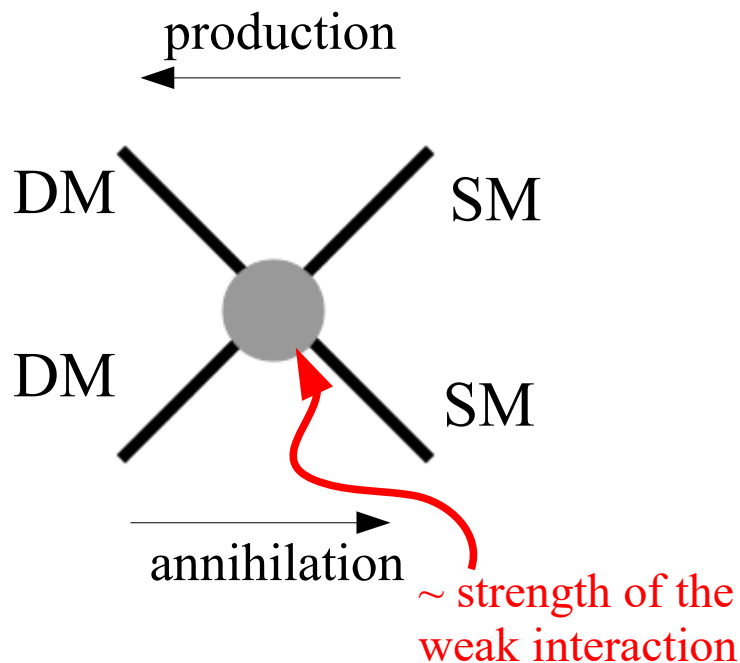
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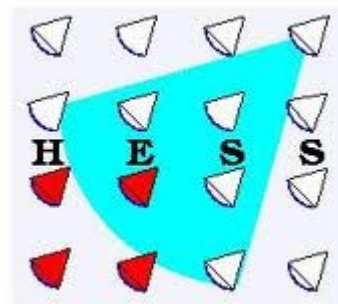
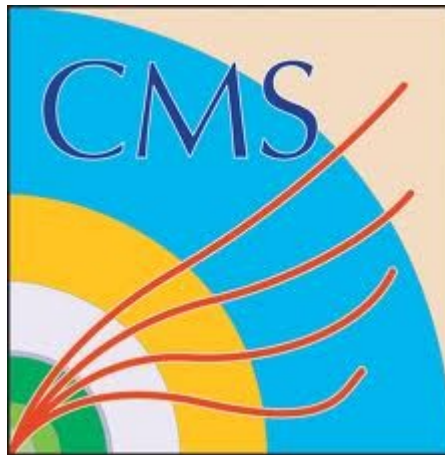
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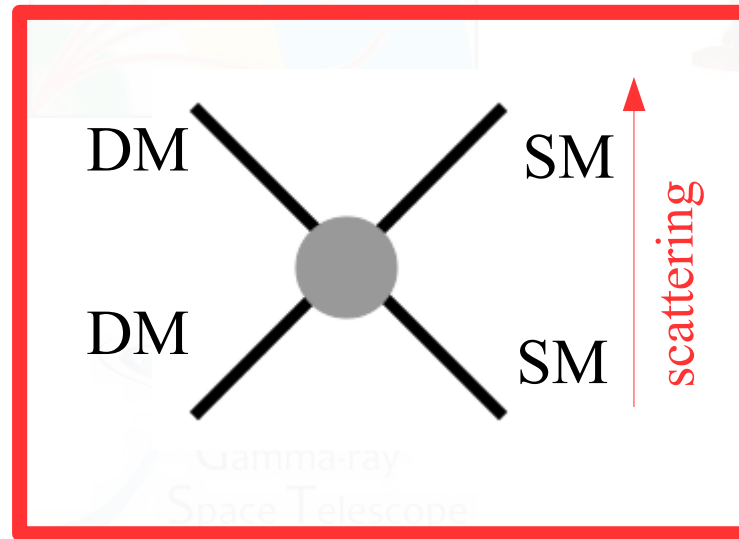
$$\langle \sigma v \rangle \simeq 3 \times 10^{-26} \text{ cm}^3 \text{ s}^{-1} = 1 \text{ pb} \cdot c$$

$$\sigma \sim \frac{g^4}{m_{\text{DM}}^2} = 1 \text{ pb}$$

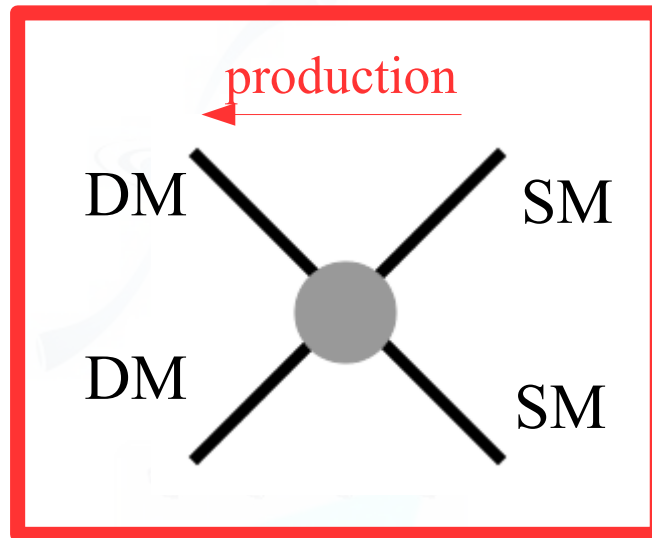
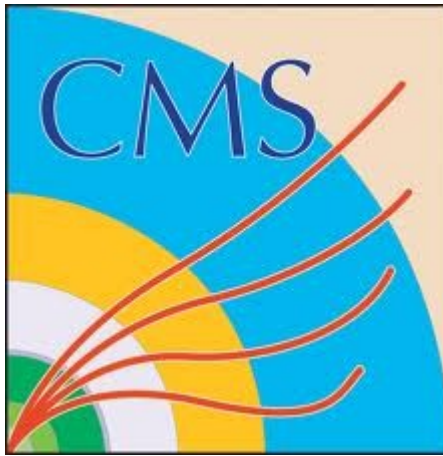
$$m_{\text{DM}} \sim \text{a few GeV} - \text{a few TeV}$$

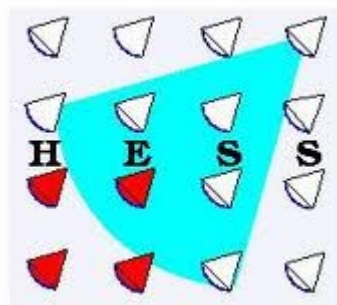
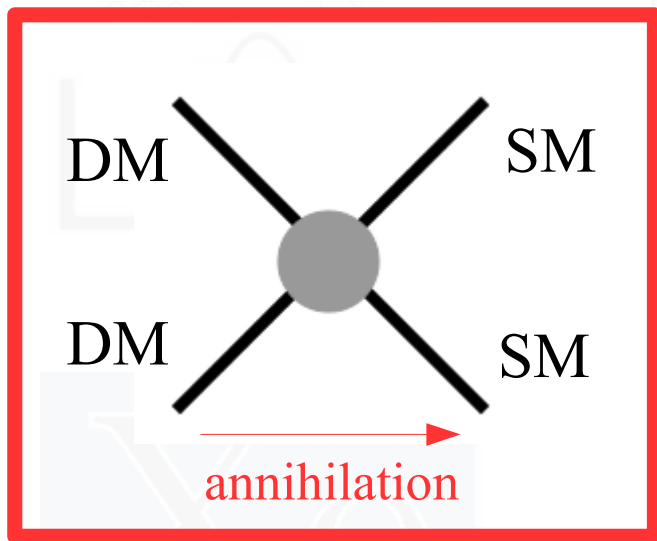
(for  $g \sim 0.01 - 1$ )



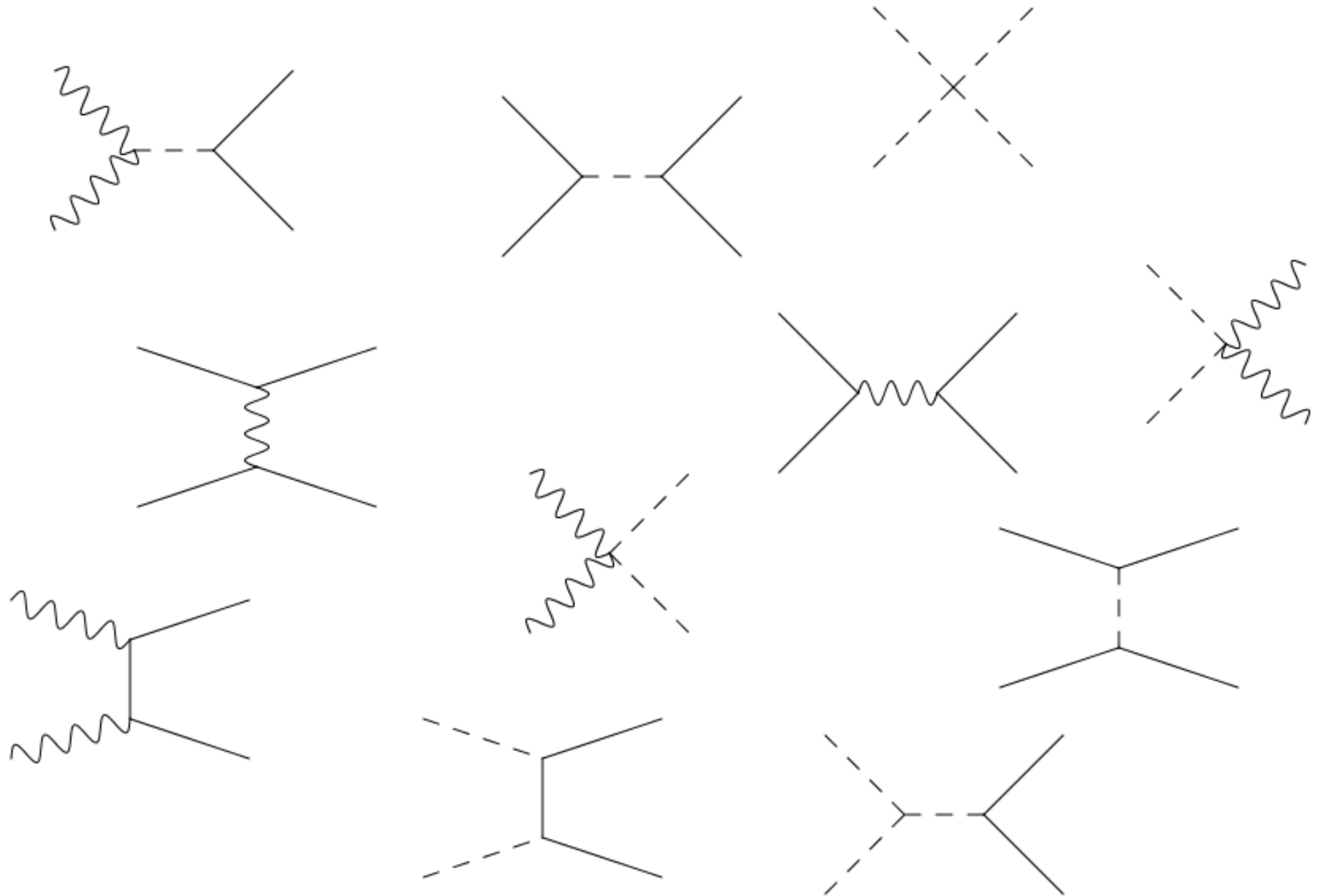








# Many possible realizations of the effective interaction



## Many possible realizations of the effective interaction

Which dark matter particle?

Which mediator (if any)?


What is the role of the mediator in the phenomenology?

# Classification of models (according to minimality)

Number of new fields

		1	2	...
Number of new symmetries	0	“minimal DM”	“next-to-minimal DM” ...	
	1	Scalar singlet DM Inert doublet model Pure Wino Pure Higgsino ...	“simplified models” ...	
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Cirelli, Fornengo, Strumia '05

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Select the  $SU(2) \times U(1)$  representation such that:

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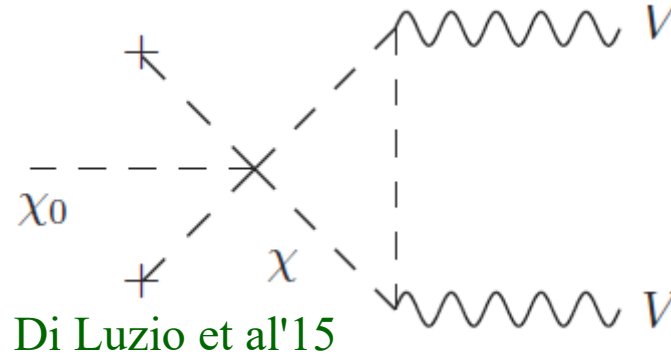
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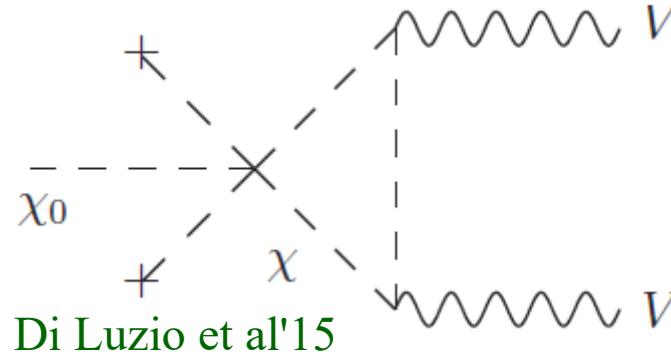
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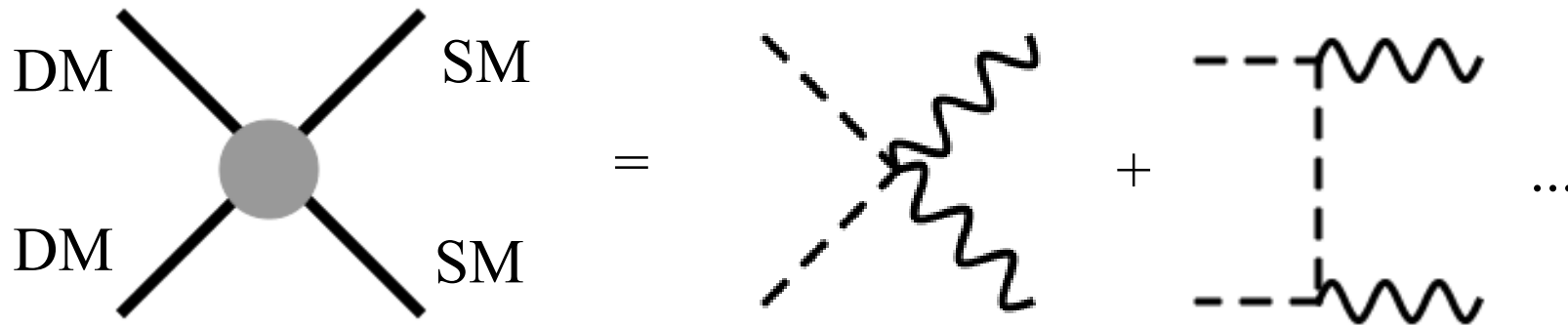
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Cirelli, Fornengo, Strumia '05



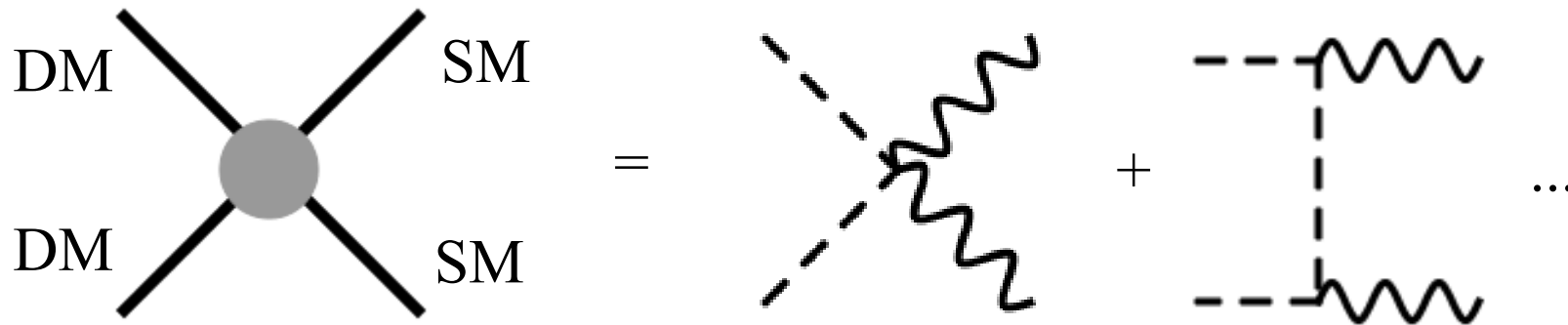
$$\langle \sigma v \rangle = \text{number} \times \frac{g^4}{m_{\text{DM}}^2}$$

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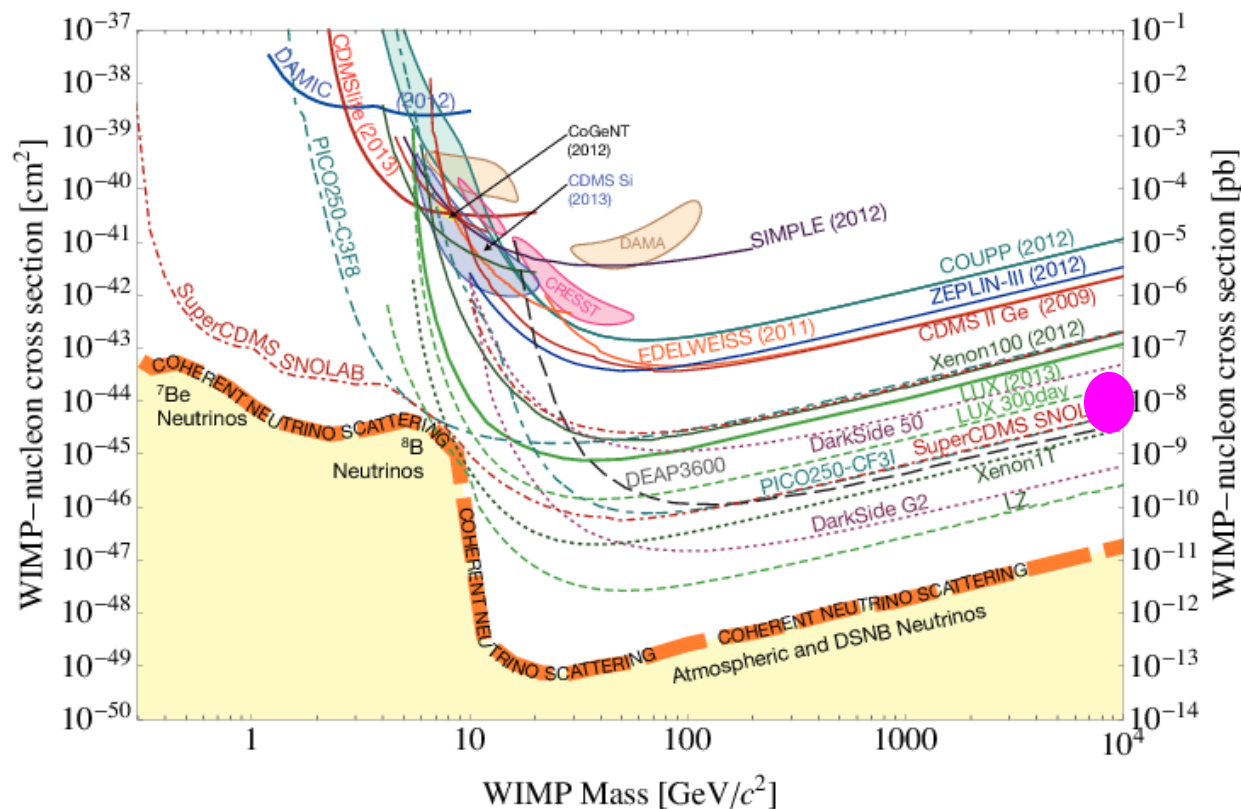


~~$$m_{\text{DM}} \approx 4.4 \text{ TeV}$$~~

$$m_{\text{DM}} \approx 10 \text{ TeV}$$

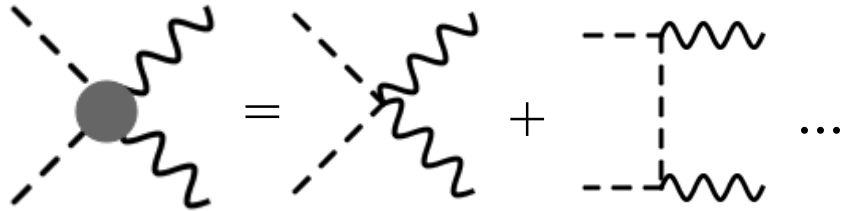
Including  
non-perturbative  
effects

Very predictive  
framework.

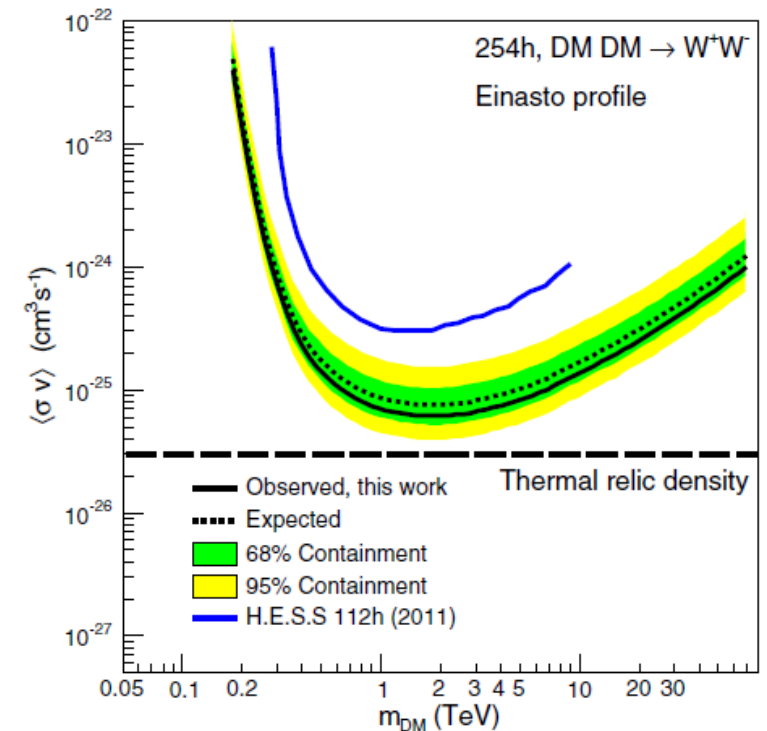


# “Minimal Dark Matter”: tests

## Indirect detection

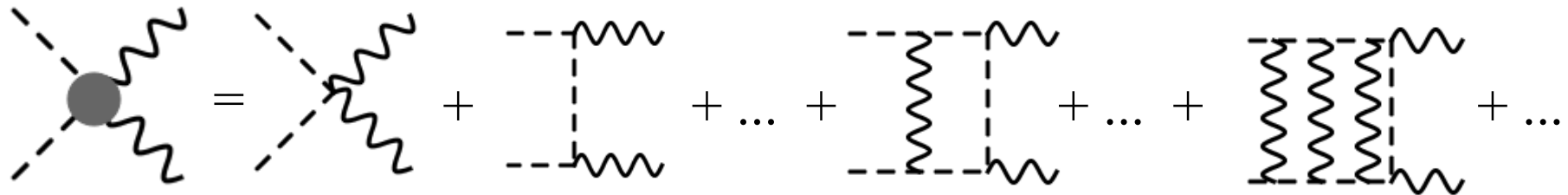


Naive estimate:  
 $(\sigma v) \simeq 3 \times 10^{-26} \text{ cm}^3 \text{ s}^{-1}$



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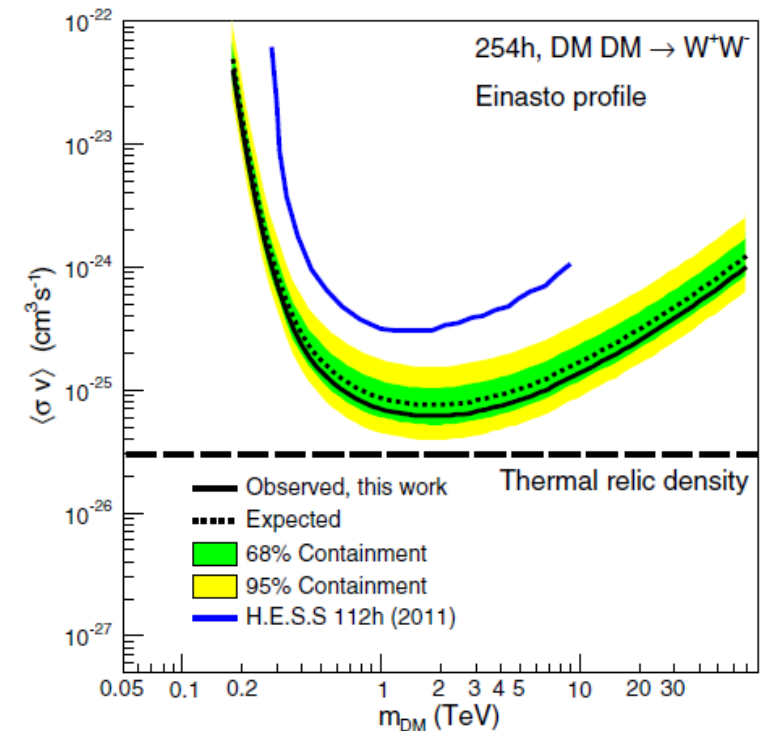
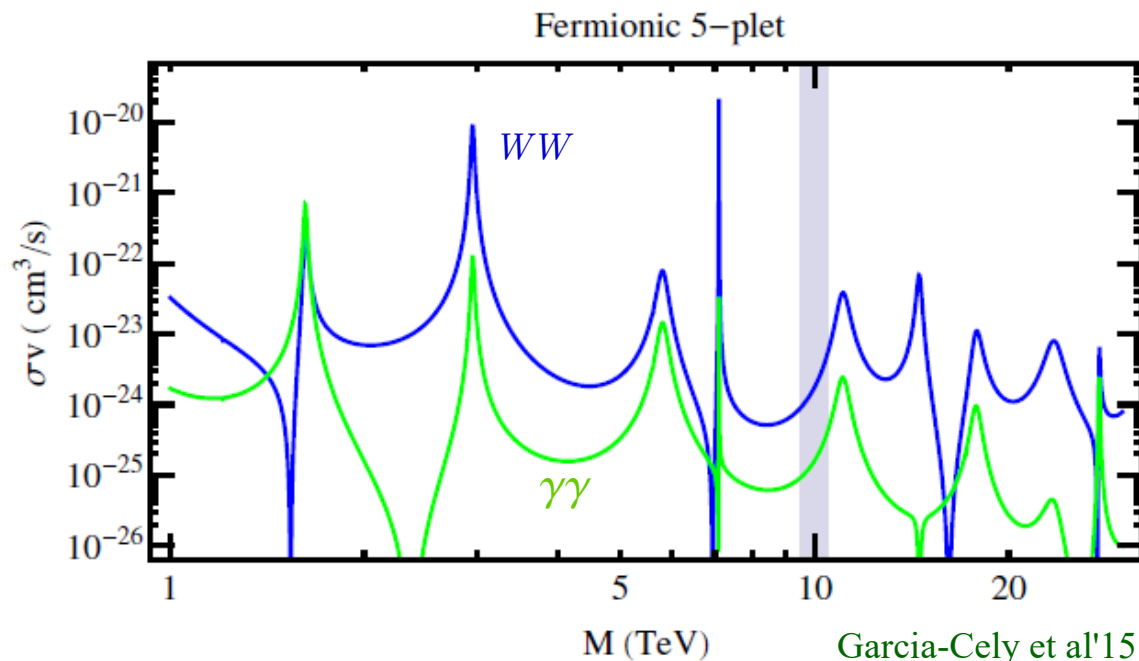
## Indirect detection



In models where the dark matter interacts with a lighter mediator, the naive calculation cannot be trusted  $\hookrightarrow$  Necessary to resum diagrams at all loops.

Hisano et al'03,04

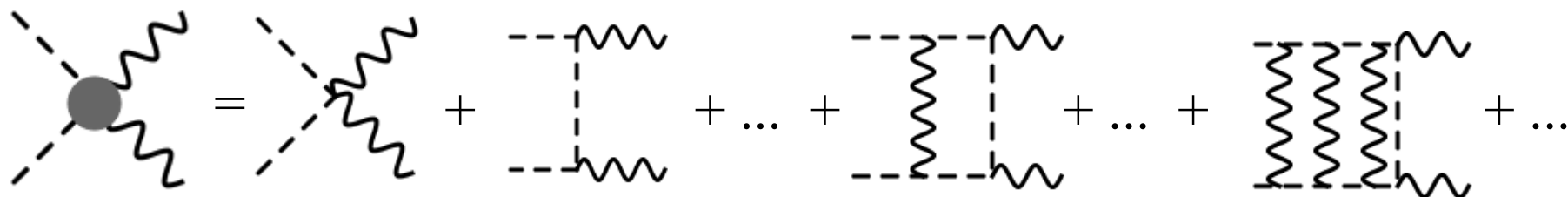
## “Sommerfeld enhancement” of the cross section





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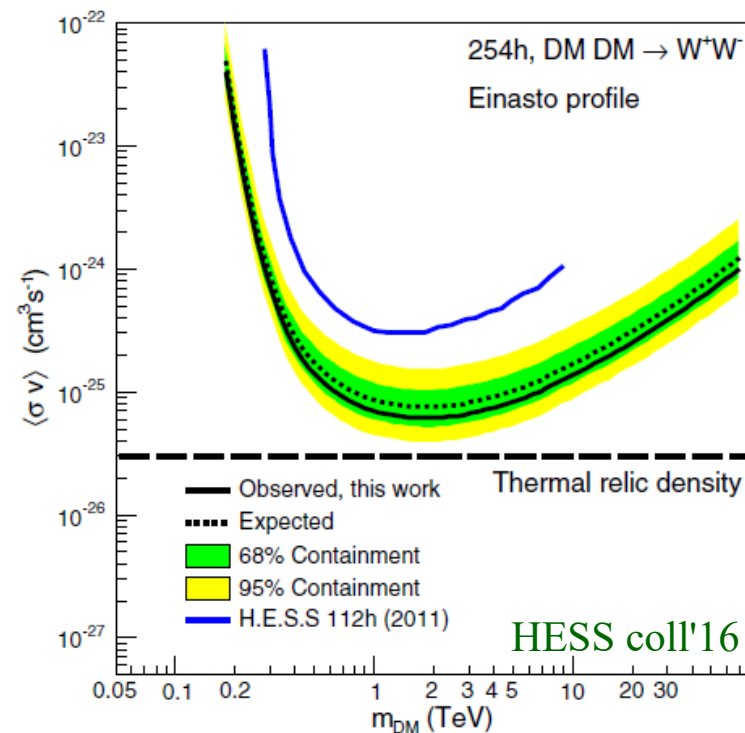
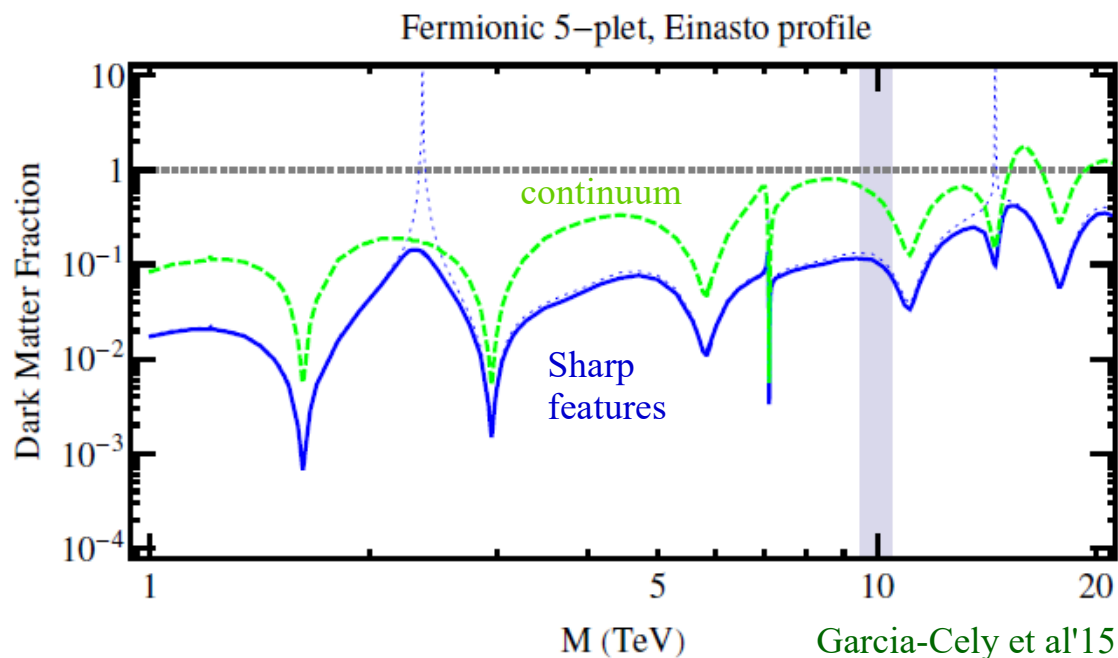
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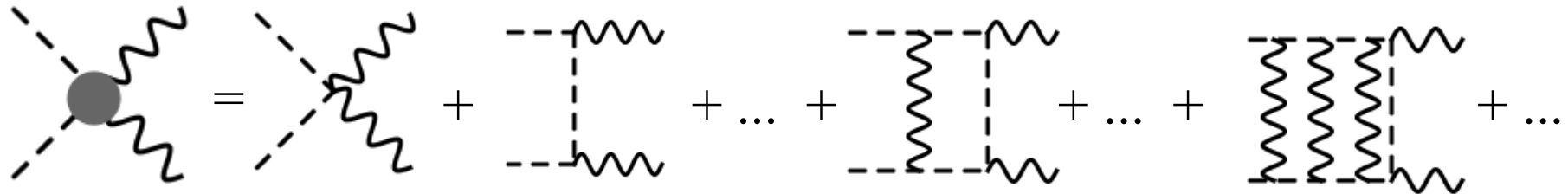
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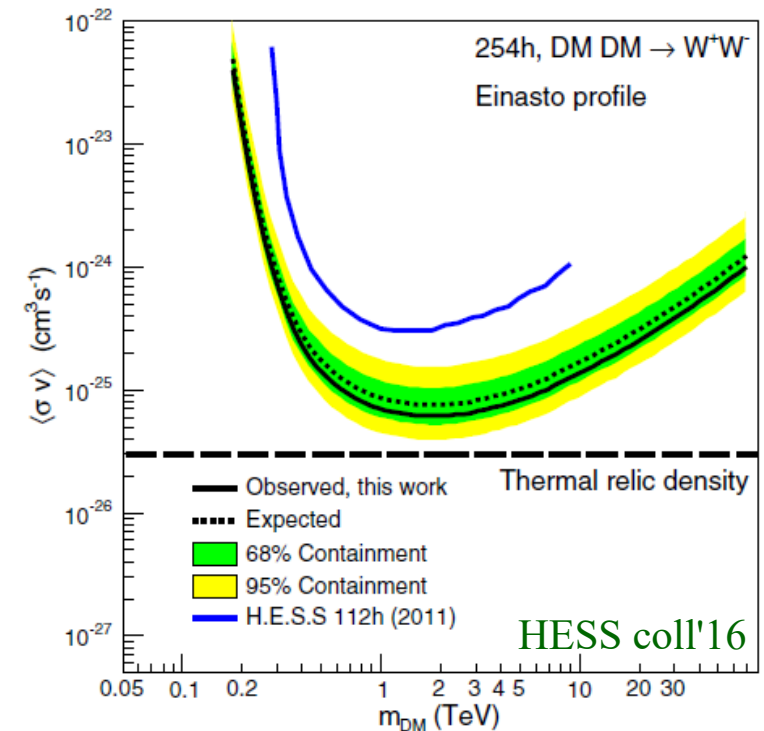
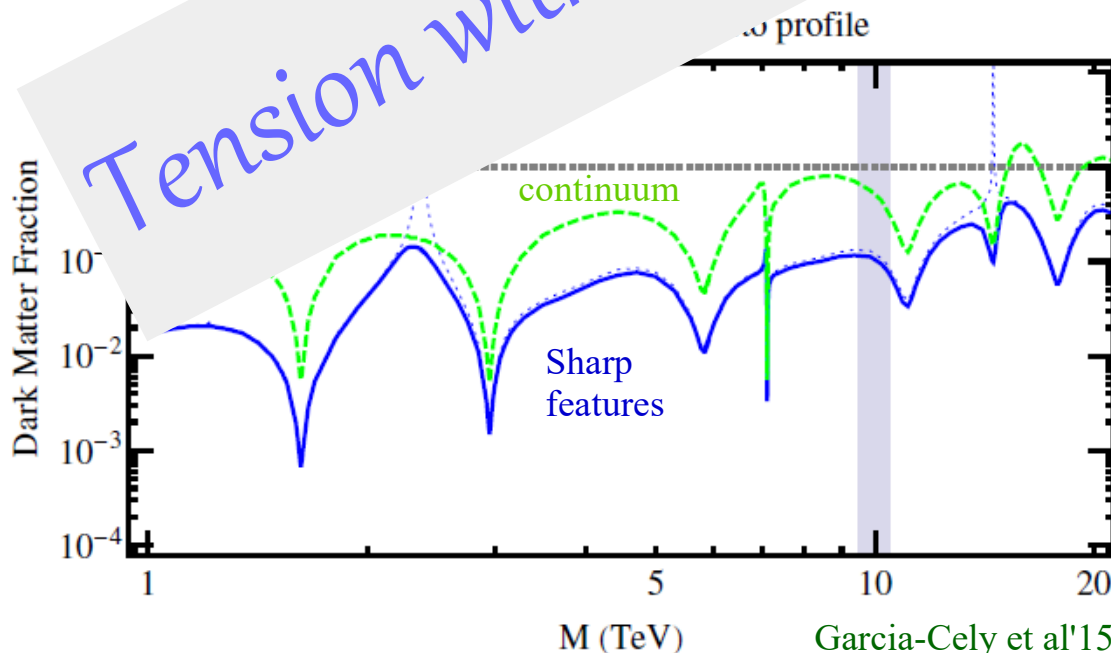
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“Sommerfeld enhancement” cross section



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# Singlet scalar dark matter

**Assumption:** Add to the SM a real scalar field,  $\phi$ , singlet under the SM gauge group. *Silveira, Zee '85, McDonald '07*

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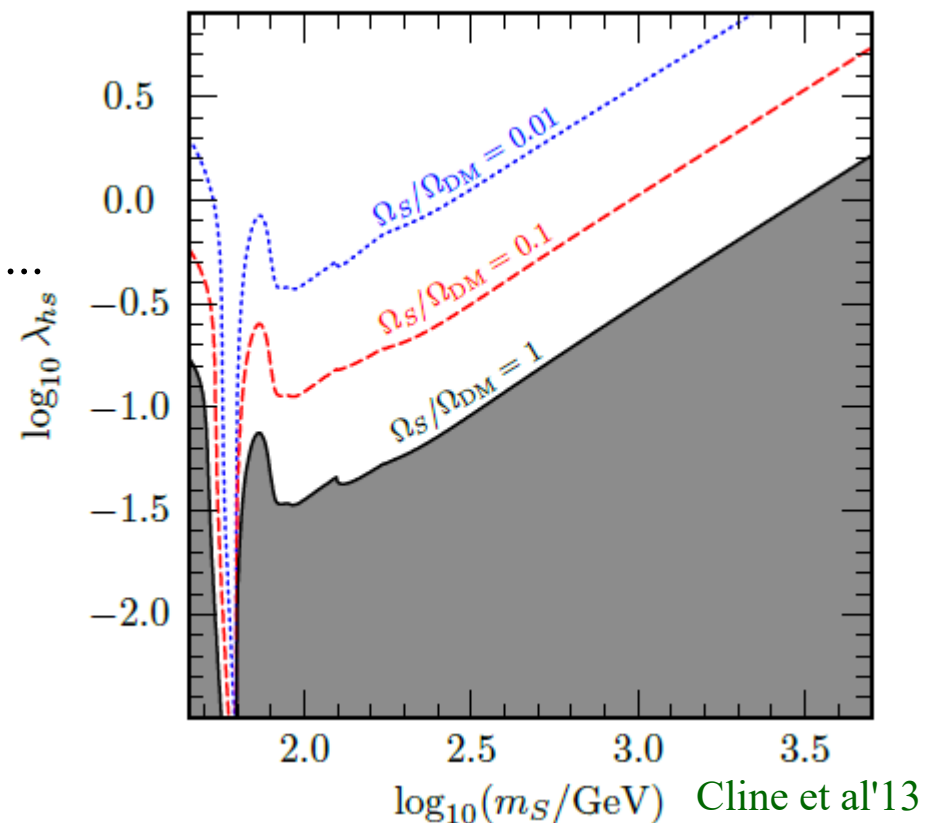
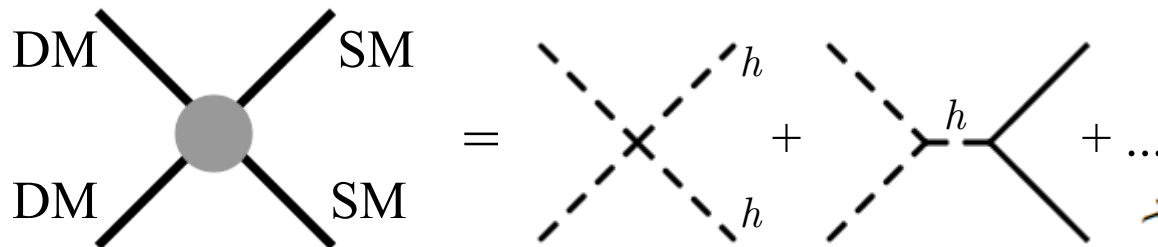
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**Assumption 2:** Impose a  $\mathbb{Z}_2$  symmetry  $\phi \rightarrow -\phi$ , SM  $\rightarrow$  SM

$\phi$  is stable and constitutes a dark matter candidate



# Singlet scalar dark matter: tests

## Collider searches

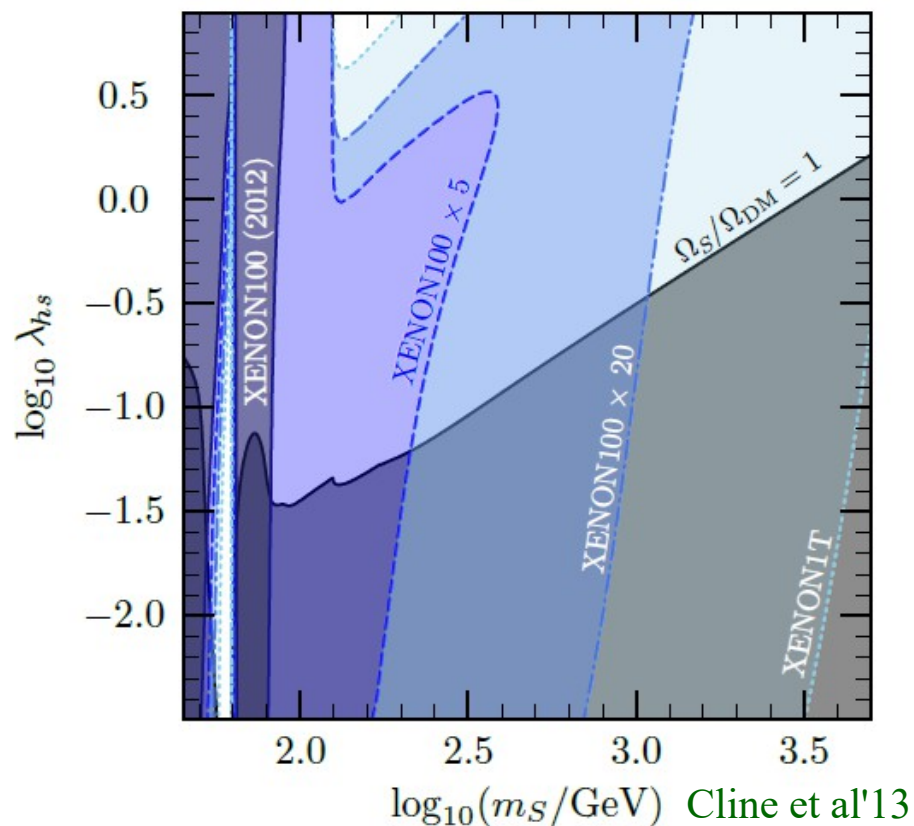
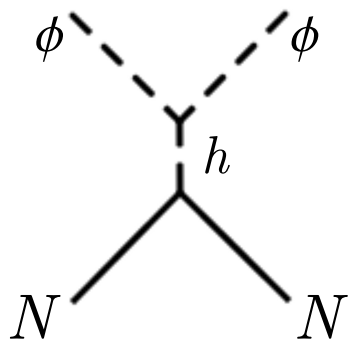
For  $m_\phi < m_h/2$ , measurements of the Higgs decay width exclude  $\lambda_{\phi H} \gtrsim 0.03$

## Indirect searches

Dwarf galaxy observations + CMB constrain the annihilation rate for low DM masses.

## Direct searches

Interactions with nucleons via Higgs exchange:





# Singlet scalar dark matter: tests

## Collider searches

For  $m_\phi < m_h/2$ , measurements of the Higgs decay width exclude  $\lambda_{\phi H} \gtrsim 0.03$

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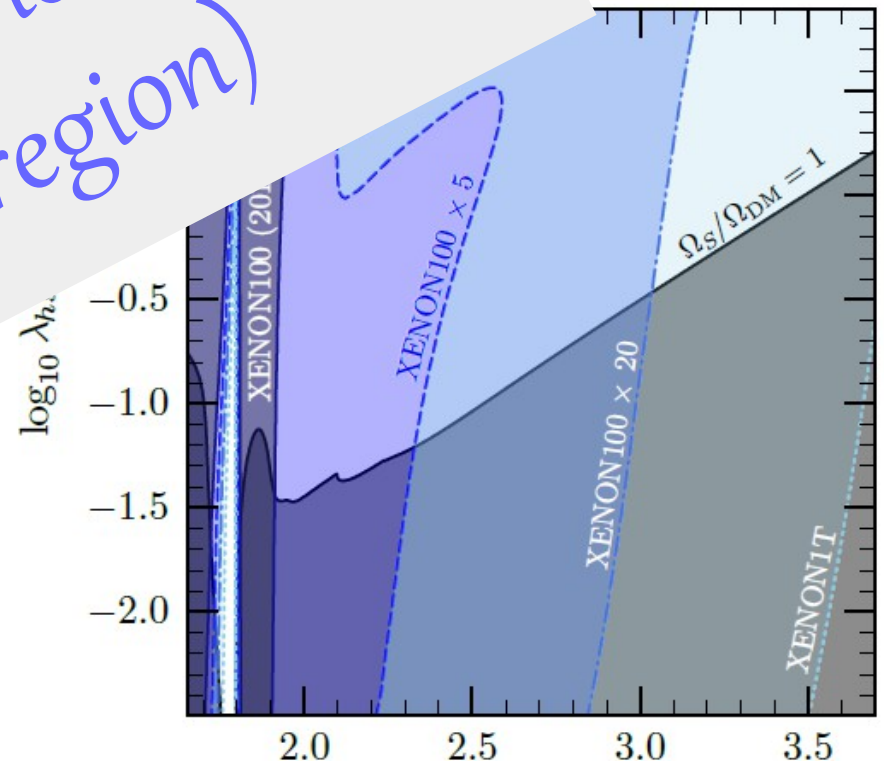
Dwarf galaxy observations + CMB constraints on annihilation rate for low DM masses.

## Direct searches

Interactions  
exchange

DM mass pushed to the TeV scale  
(or resonant region)

$N$



$\log_{10}(m_S/\text{GeV})$  Cline et al'13

# Classification of models (according to minimality)

Number of new fields

Number of  
new symmetries

		1	2	...
0		“minimal DM”	“next-to-minimal DM” ...	
		Secluded singlet DM <b>Inert doublet model</b> Pure W Pure Higgsino ...	“simplified models” ...	
1				
...				

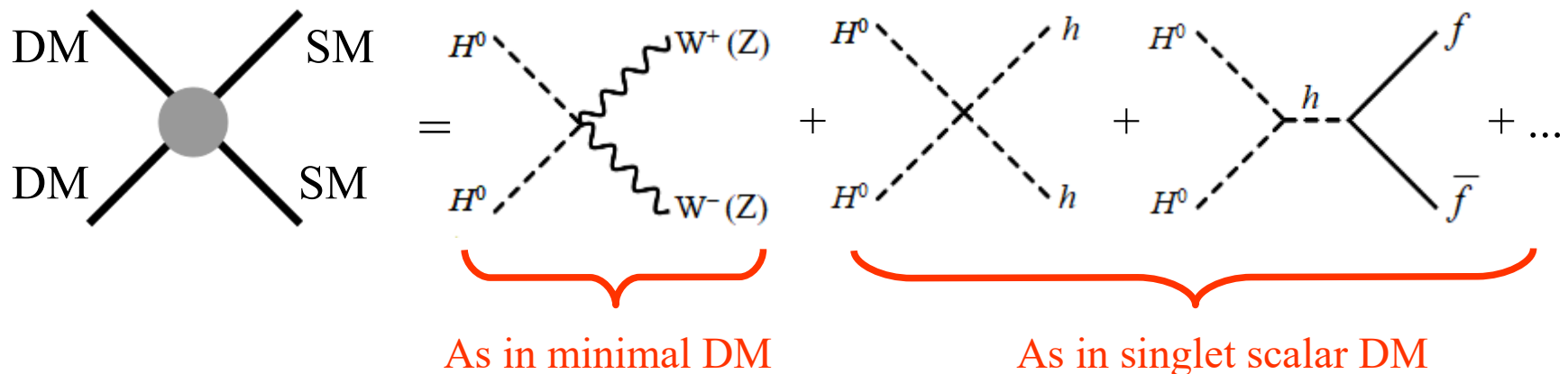
# Inert doublet dark matter model

**Assumption 1:** Add to the SM a scalar doublet with hypercharge 1/2,  $\eta$ , (copy of the SM Higgs)

**Assumption 2:** Impose a  $\mathbb{Z}_2$  symmetry  $\eta \rightarrow -\eta$ , SM  $\rightarrow$  SM

The lightest component of the doublet is stable and constitutes a dark matter candidate

$$\eta = \begin{pmatrix} H^+ \\ \frac{1}{\sqrt{2}} (H^0 + iA^0) \end{pmatrix}$$

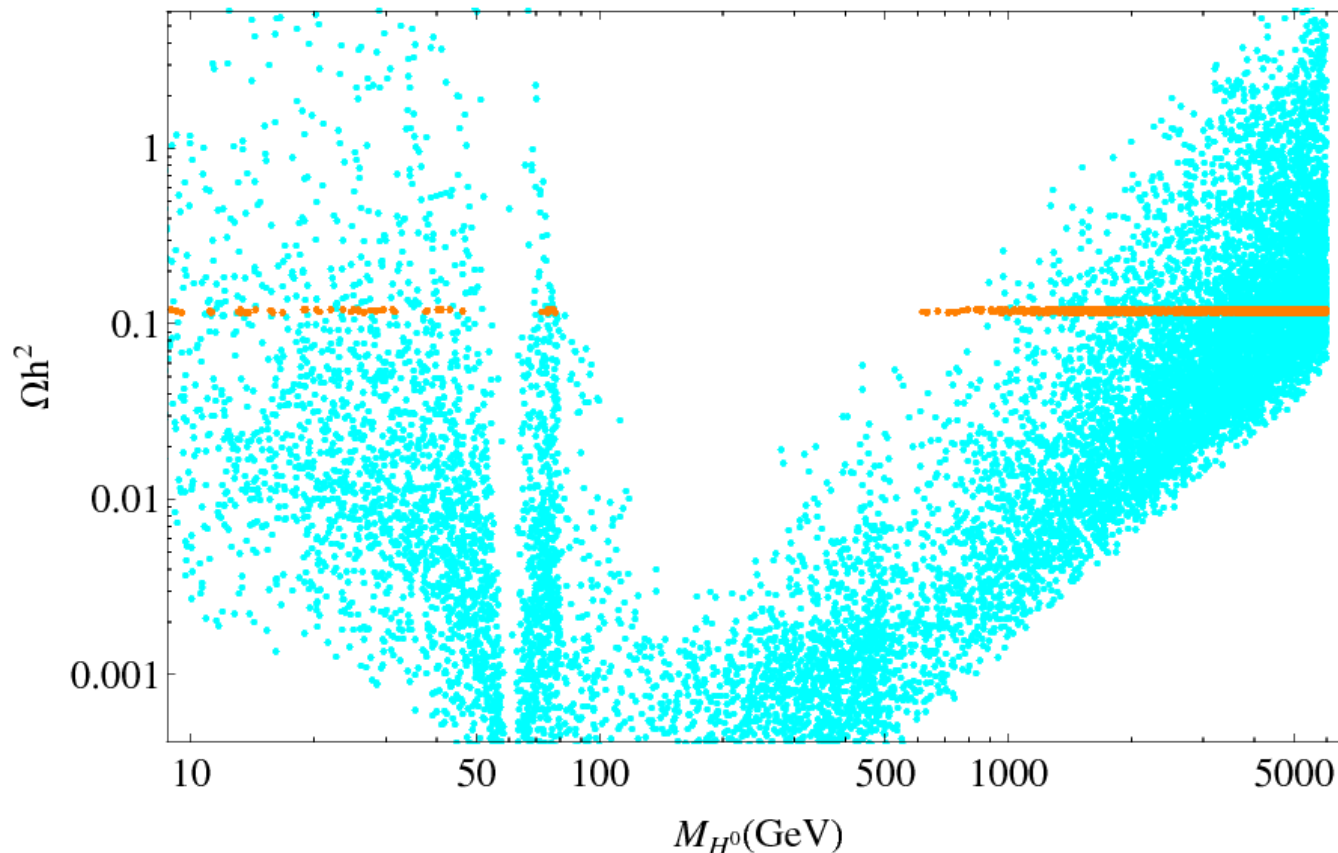


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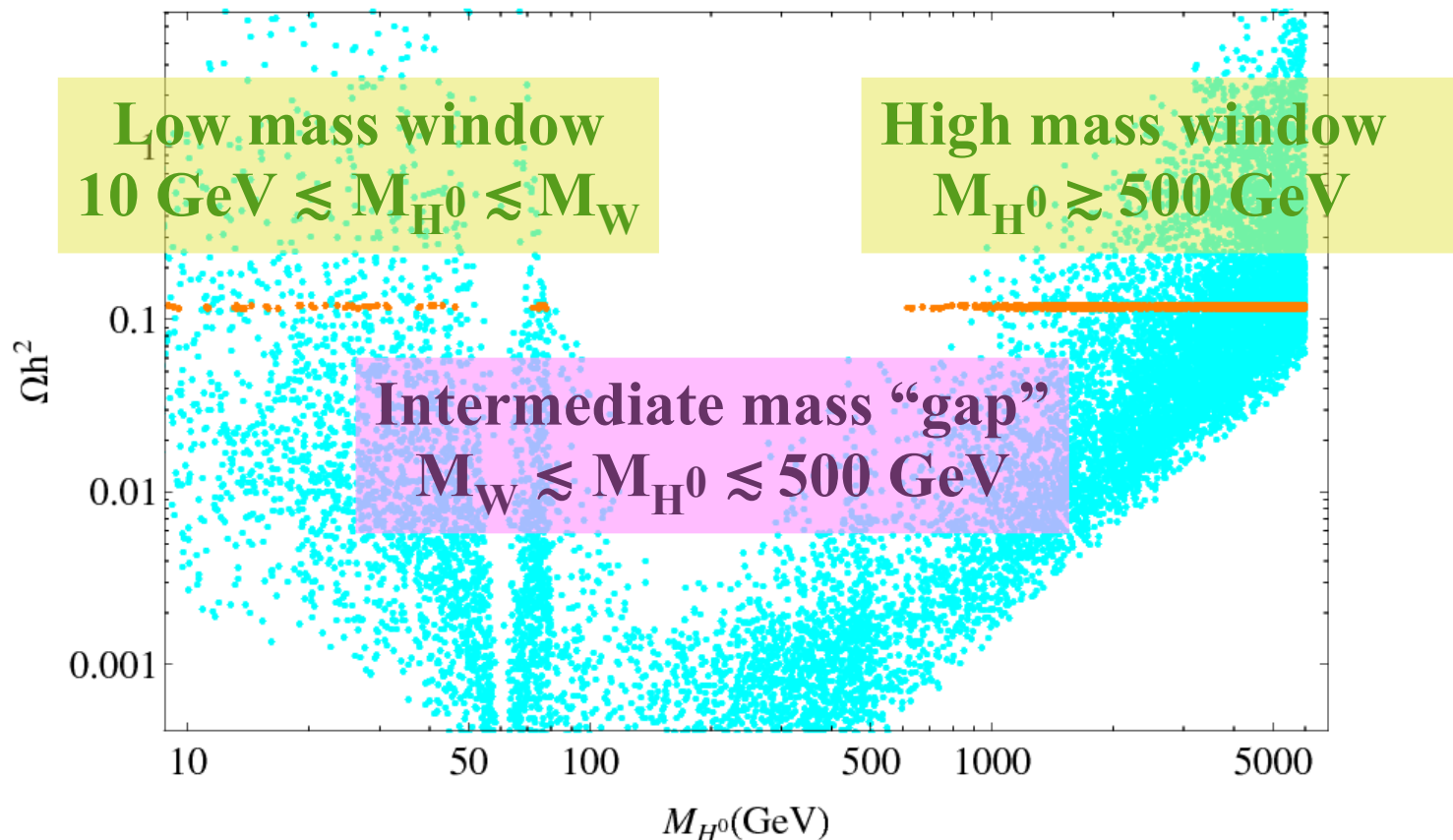


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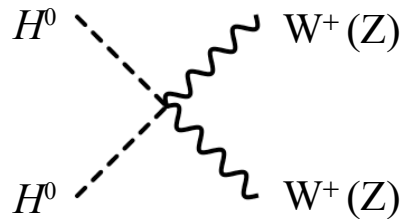
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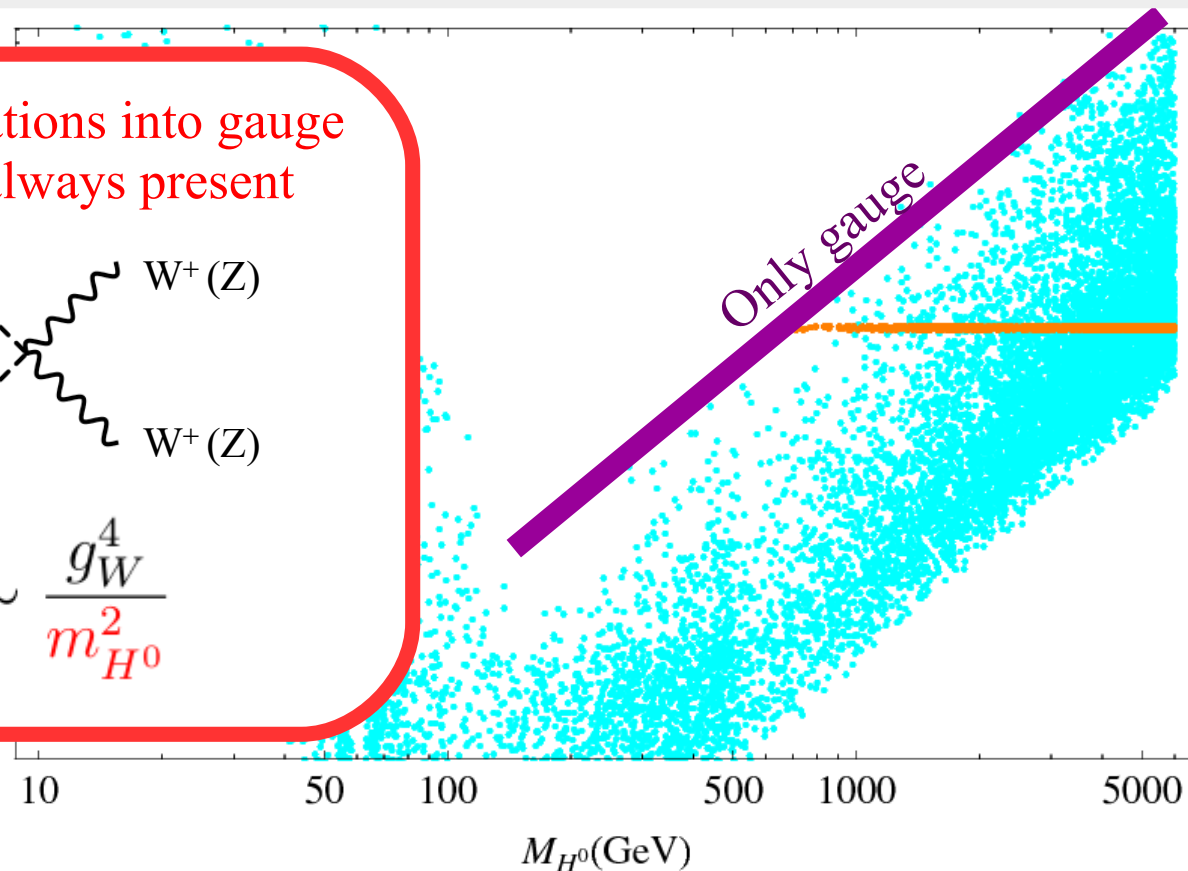
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Annihilations into gauge bosons always present



$$\sigma \sim \frac{g_W^4}{m_{H^0}^2}$$



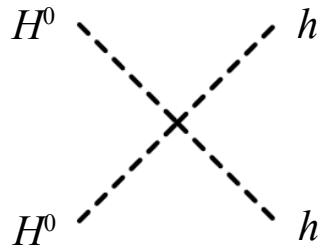
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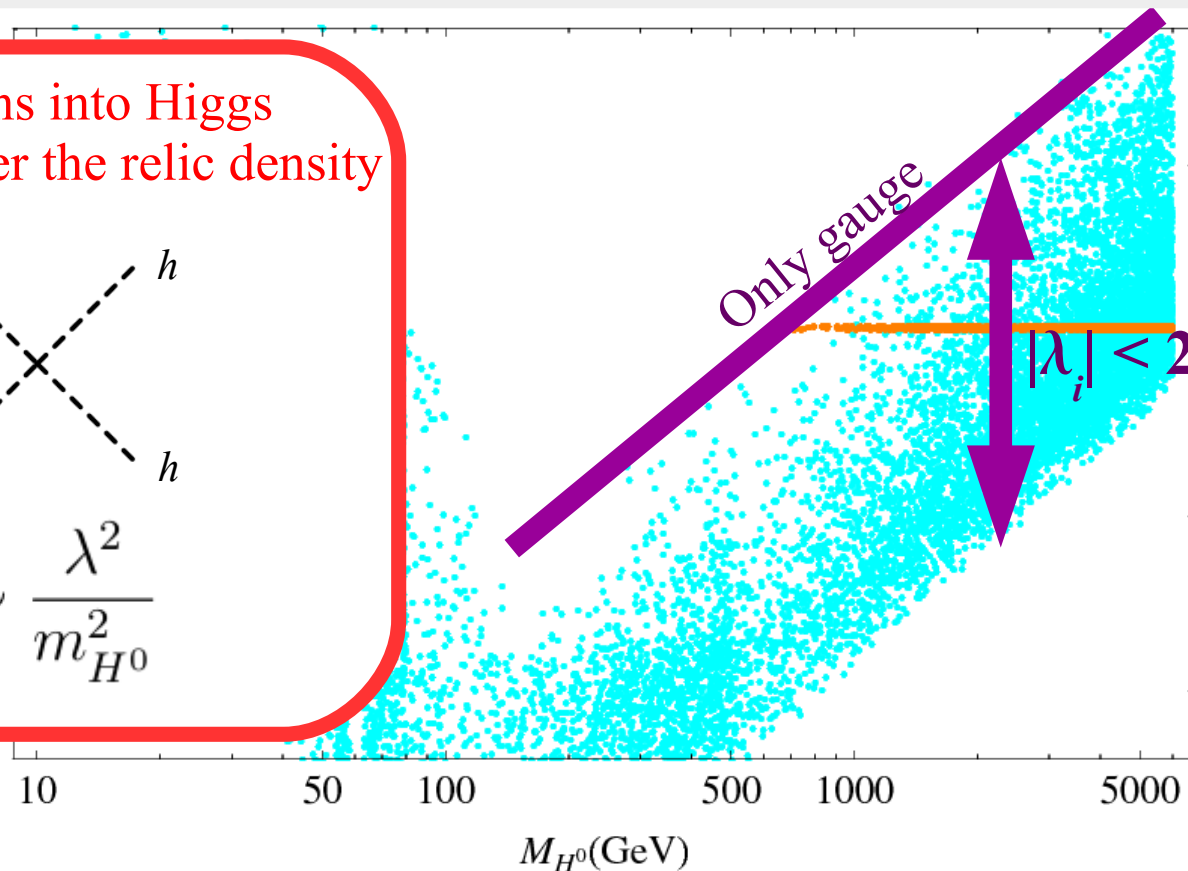
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The lightest component of the doublet is stable and constitutes a dark matter candidate

Annihilations into Higgs bosons lower the relic density



$$\sigma \sim \frac{\lambda^2}{m_{H^0}^2}$$

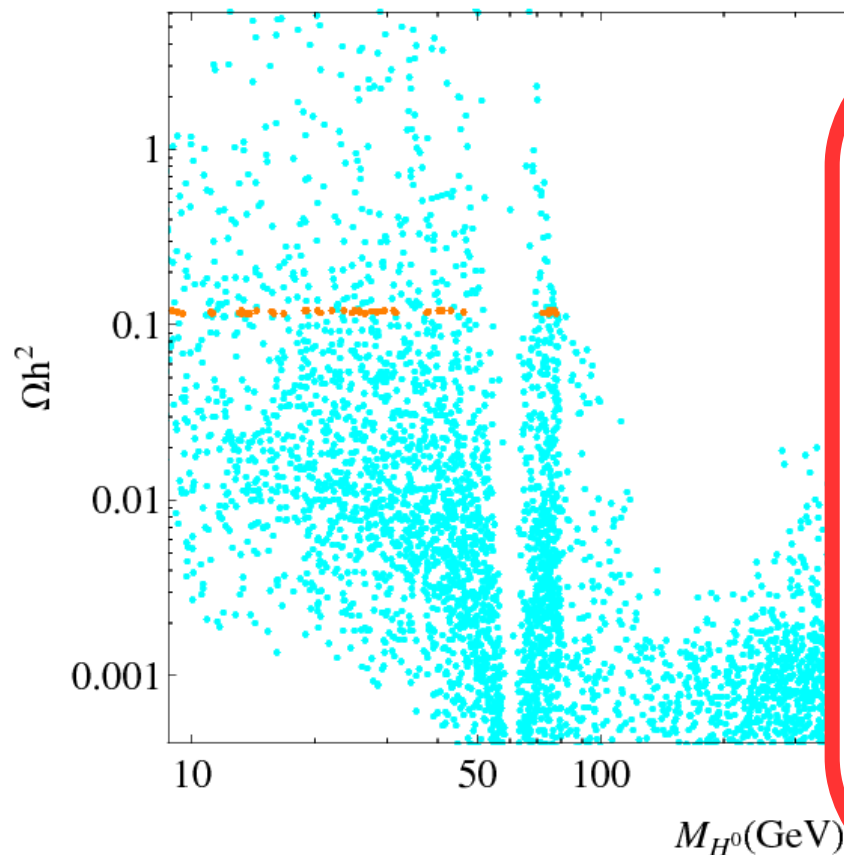


# Inert doublet dark matter model

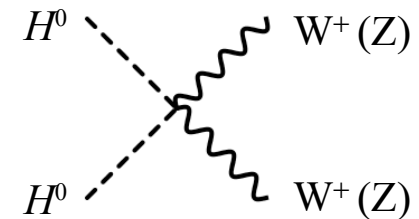
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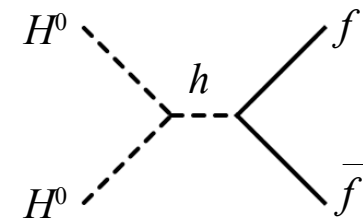
The lightest component of the doublet is stable and constitutes a dark matter candidate



Annihilations into gauge bosons kinematically forbidden



Annihilations into SM fermions possible. Adjust quartic couplings to generate the correct relic density



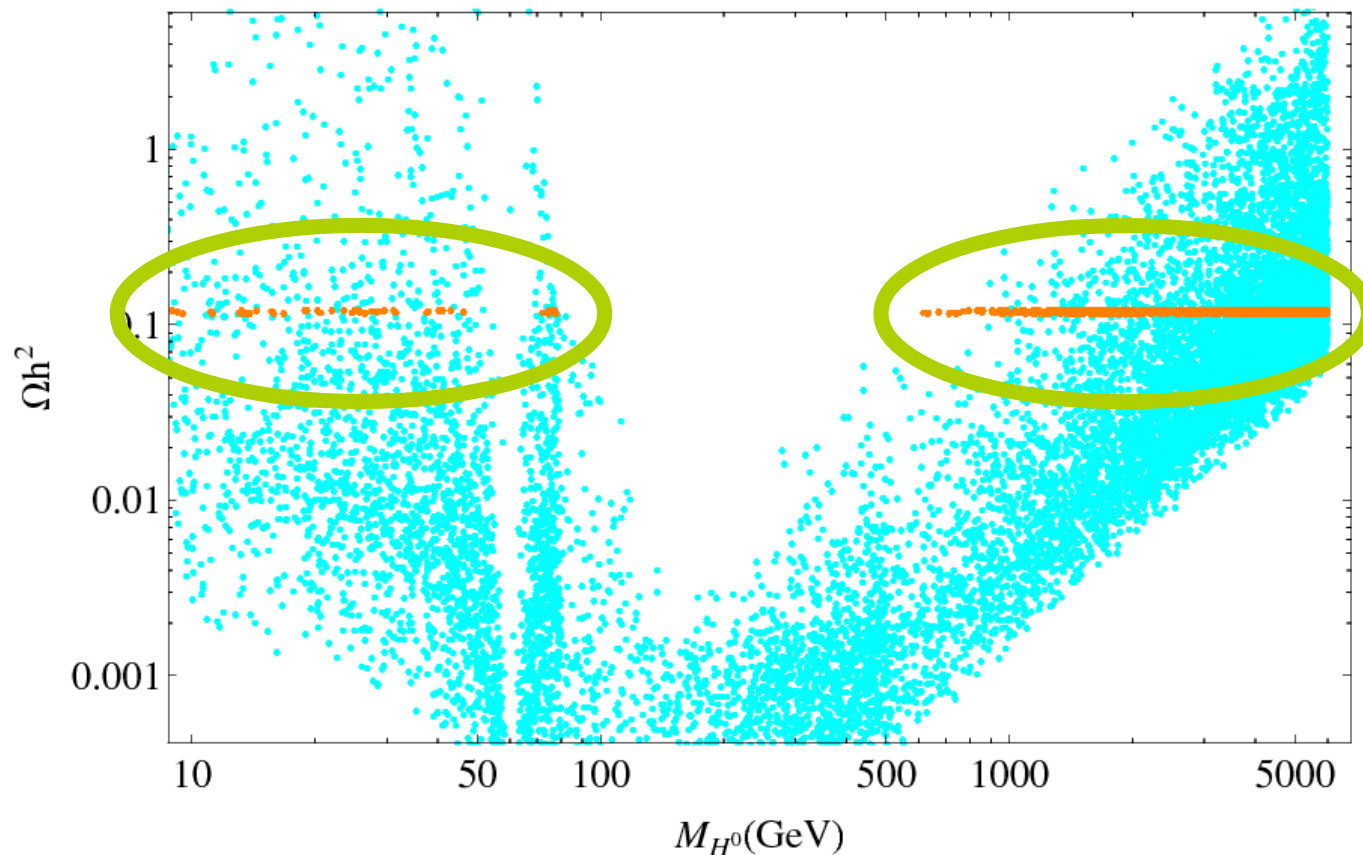


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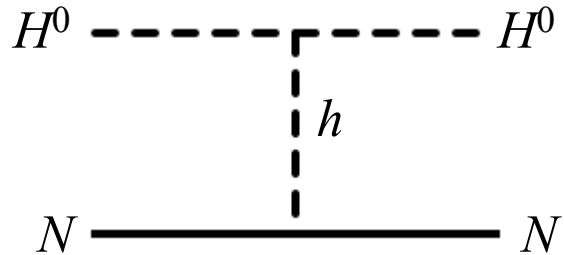
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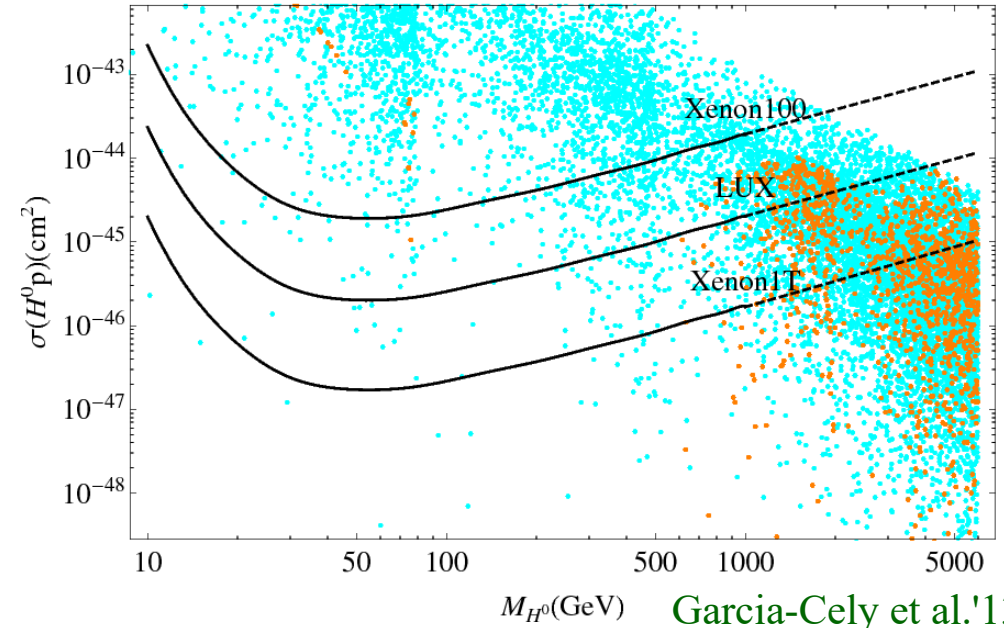


# Inert doublet dark matter model: tests

## Direct detection

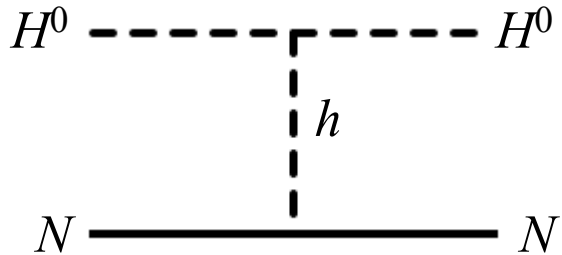


$$\sigma(H^0 p) \simeq 4 \times 10^{-44} \lambda_{H^0}^2 \left( \frac{1 \text{ TeV}}{M_{H^0}} \right)^2 \text{ cm}^2$$

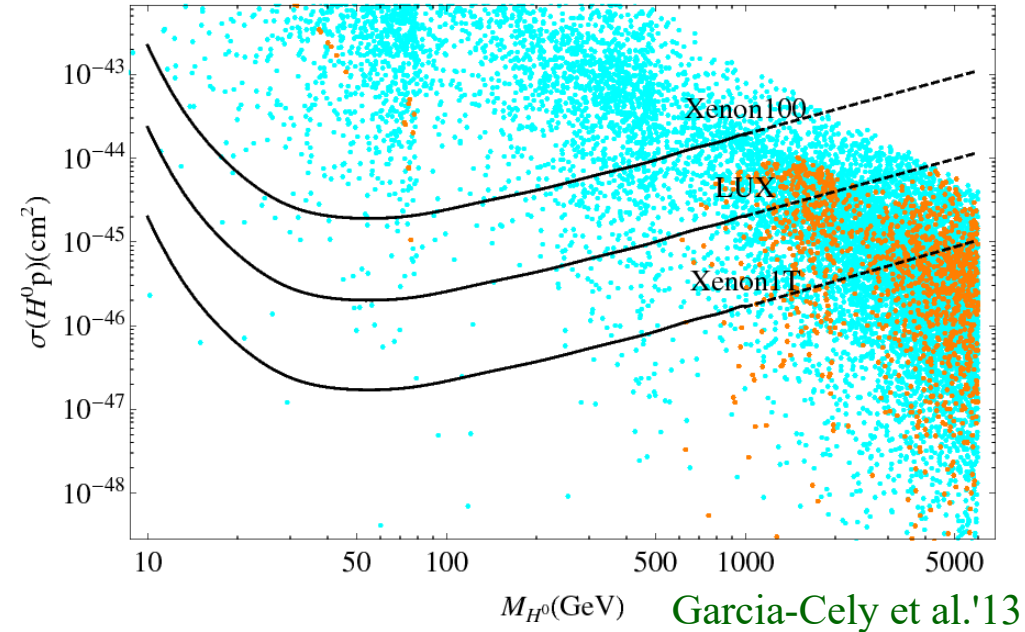


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## Direct detection

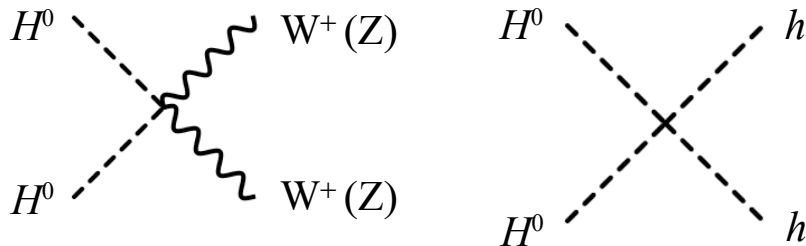


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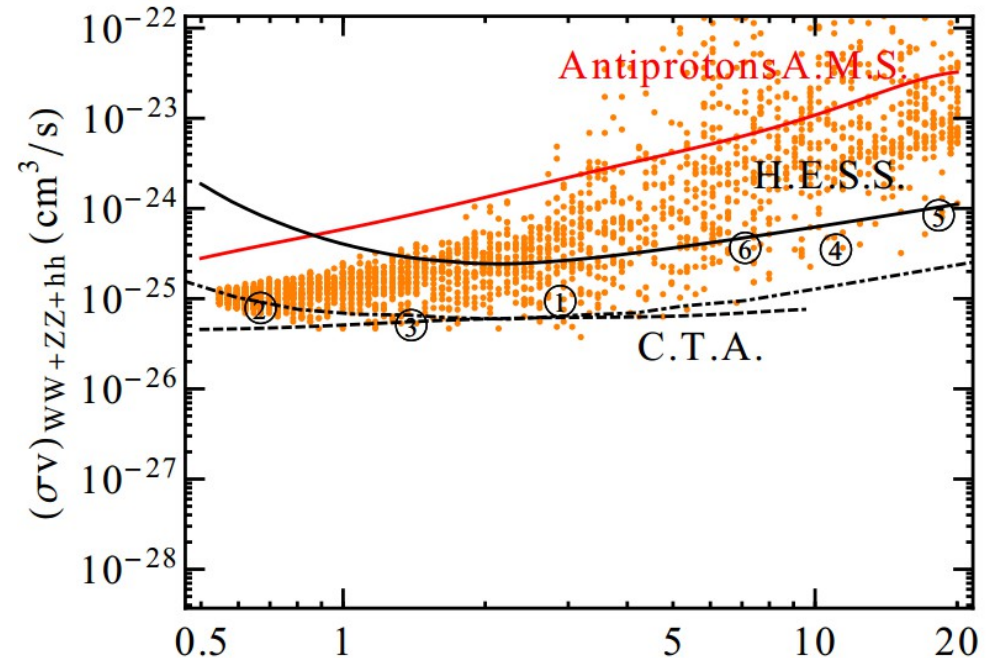


## Indirect detection

Annihilations into  $W, Z$  and  $h$



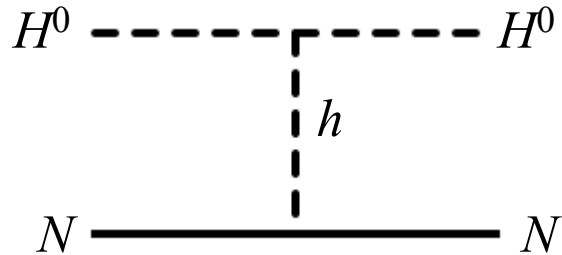
$H^0$  has electroweak charge  
→ Sommerfeld enhancement  
of the cross section



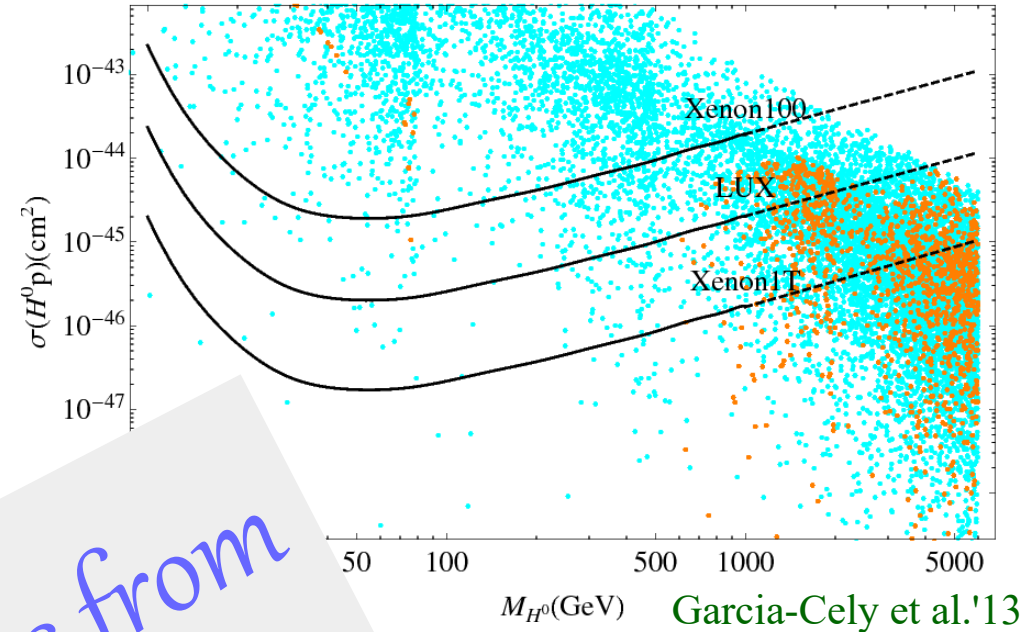
$M_{H^0}$  (TeV) Garcia-Cely et al.'17

# Inert doublet dark matter model: tests

## Direct detection

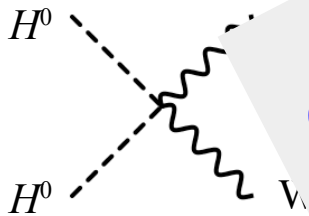


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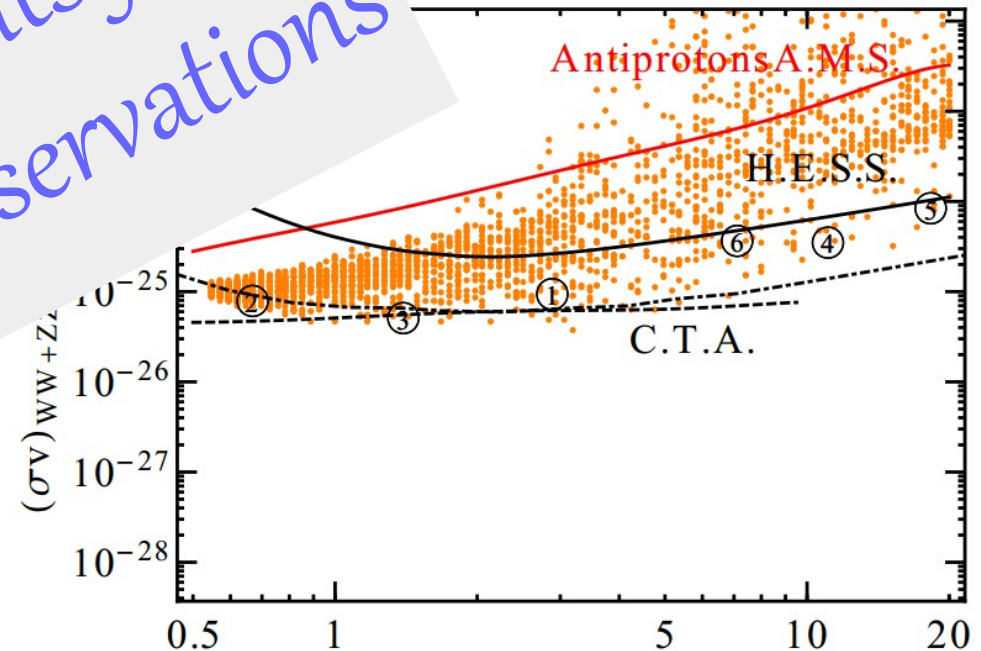


## Indirect detection

Annihilations into



Strong constraints from gamma-ray observations



$H^0$  has weak charge  
→ Sommerfeld enhancement  
of the cross section

$M_{H^0}$  (TeV) Garcia-Cely et al.'17

# Classification of models (according to minimality)

Number of new fields

Number of  
new symmetries

		1	2	...
0		“minimal DM”	“next-to-minimal DM” ...	
		Scalar singlet DM Inert doublet model Pure Wino Pure Higgsino ...	“simplified models” ...	
1				
...				

# “Simplified models”

**Assumption 1:** Add to the SM a dark matter candidate,  $\chi$ , and a mediator,  $\eta$ , heavier than the DM candidate.

**Assumption 2:** Impose a  $\mathbb{Z}_2$  symmetry  $\chi \rightarrow -\chi$ ,  $\eta \rightarrow -\eta$ ,  $\text{SM} \rightarrow \text{SM}$

$\chi$  is the lightest particle of the  $\mathbb{Z}_2$ -odd sector and is stable

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Many possible realizations:

- DM spin/parity
- DM gauge quantum numbers
- t-channel or s-channel mediation
- Coupling to the left- and/or right-handed electron, to the up-quark, to the bottom-quark...



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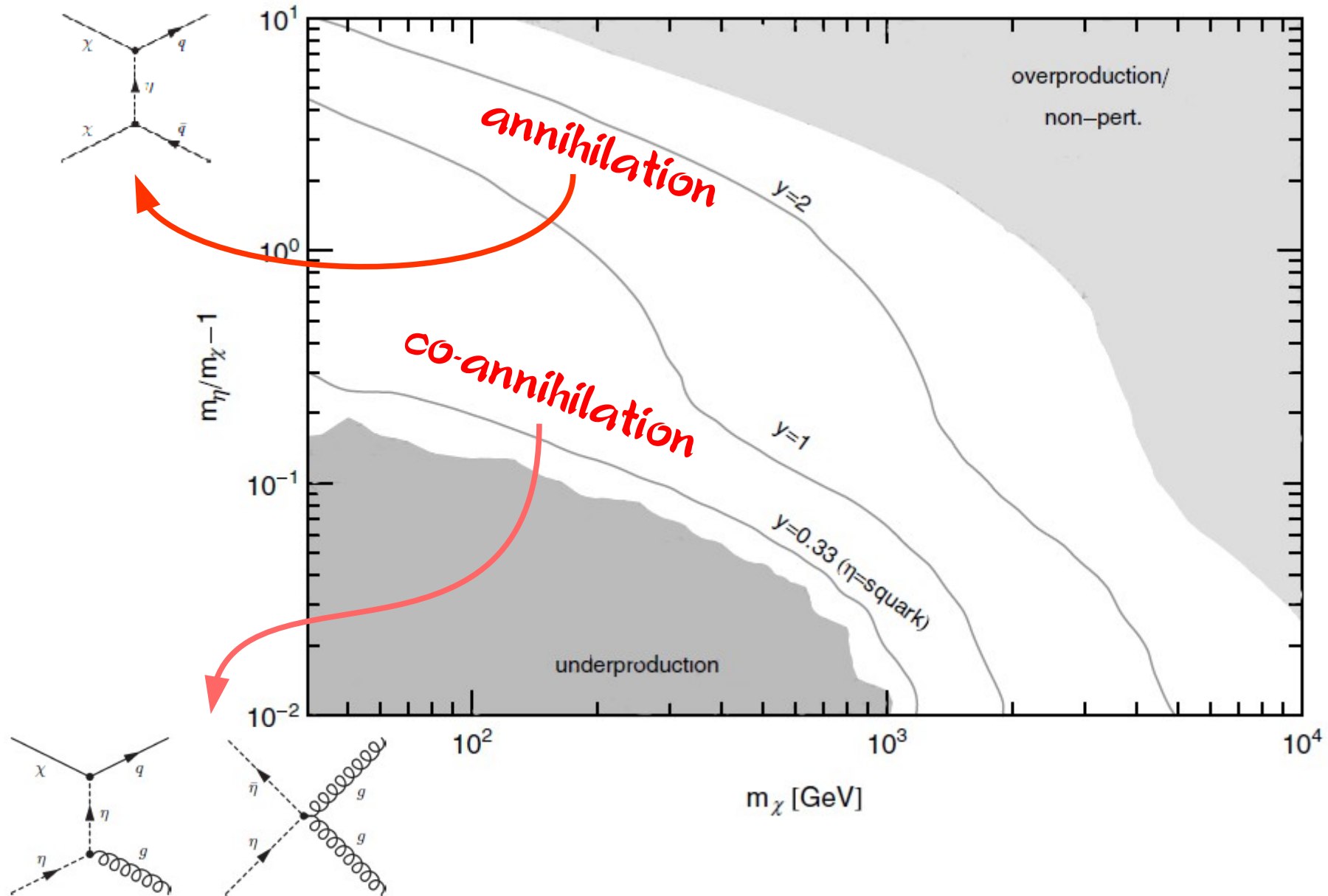
e.g. Majorana fermion dark matter, coupling to a right-handed up-quark via a scalar mediator.

Free parameters: DM mass ( $m_\chi$ ), mediator mass ( $m_\eta$ ), Yukawa coupling ( $y$ )



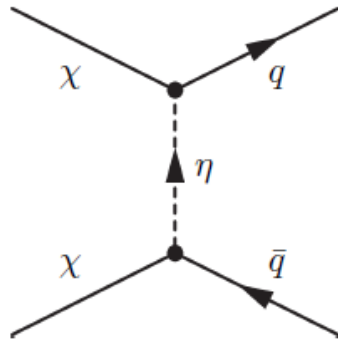
# “Simplified models”

Dark matter production fixes one of the parameters of the model.



# “Simplified models”: tests

## Indirect detection

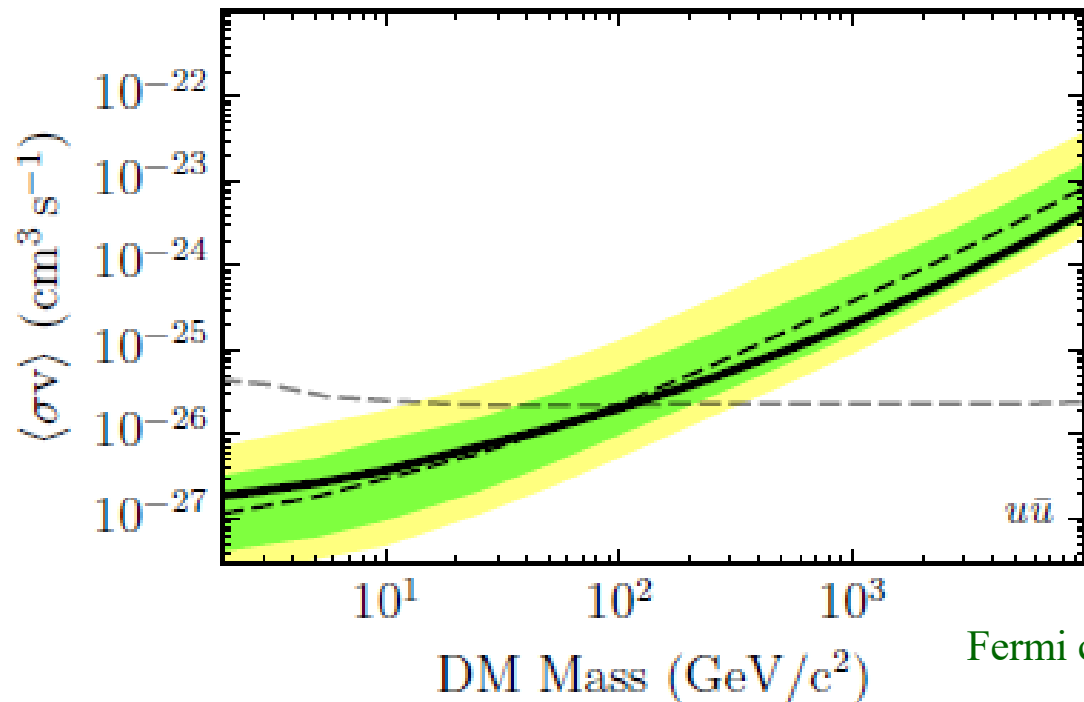


p-wave suppressed, helicity suppressed.

Expected annihilation cross-section in our Galaxy  
 $\sim 5$ -6 orders of magnitude smaller than  
the thermal cross section.



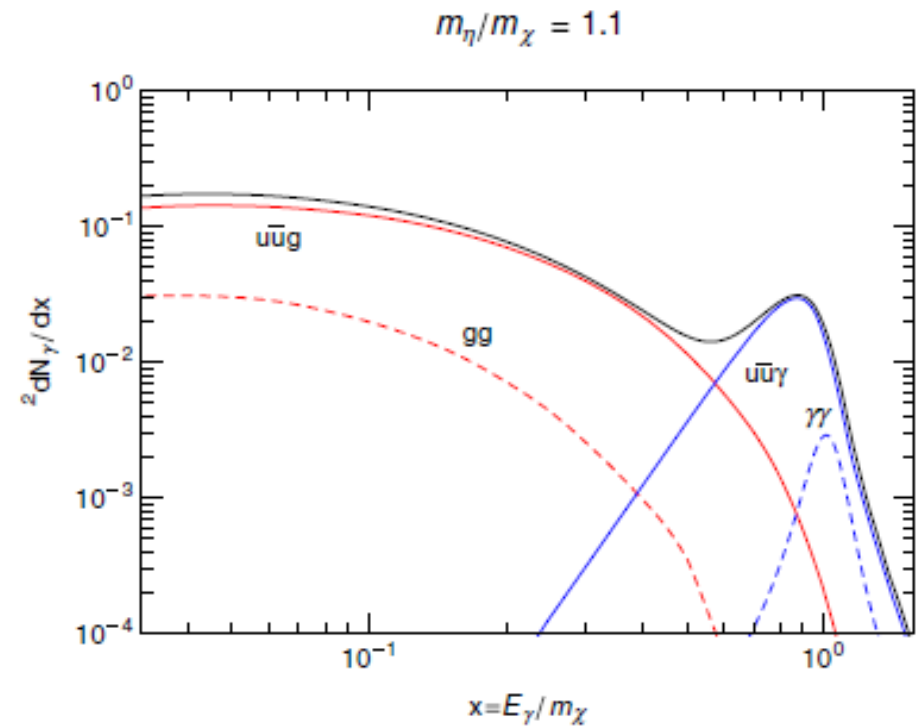
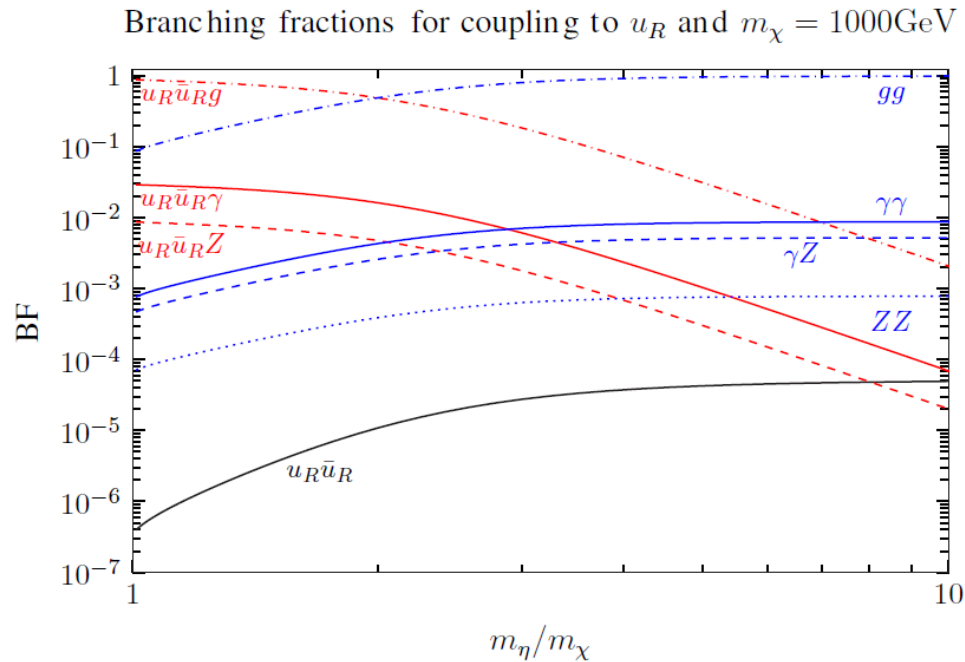
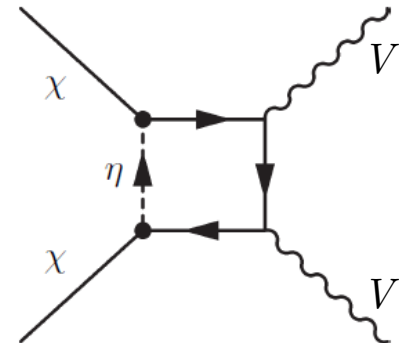
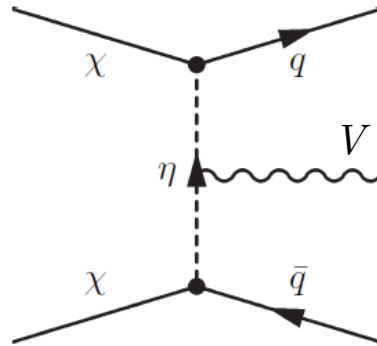
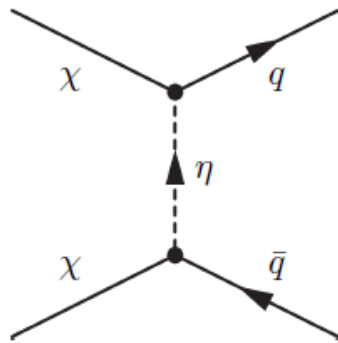
Limits from dwarf galaxy  
observations on the  $u\bar{u}$   
final state do not constrain  
this model.



Fermi coll.'15

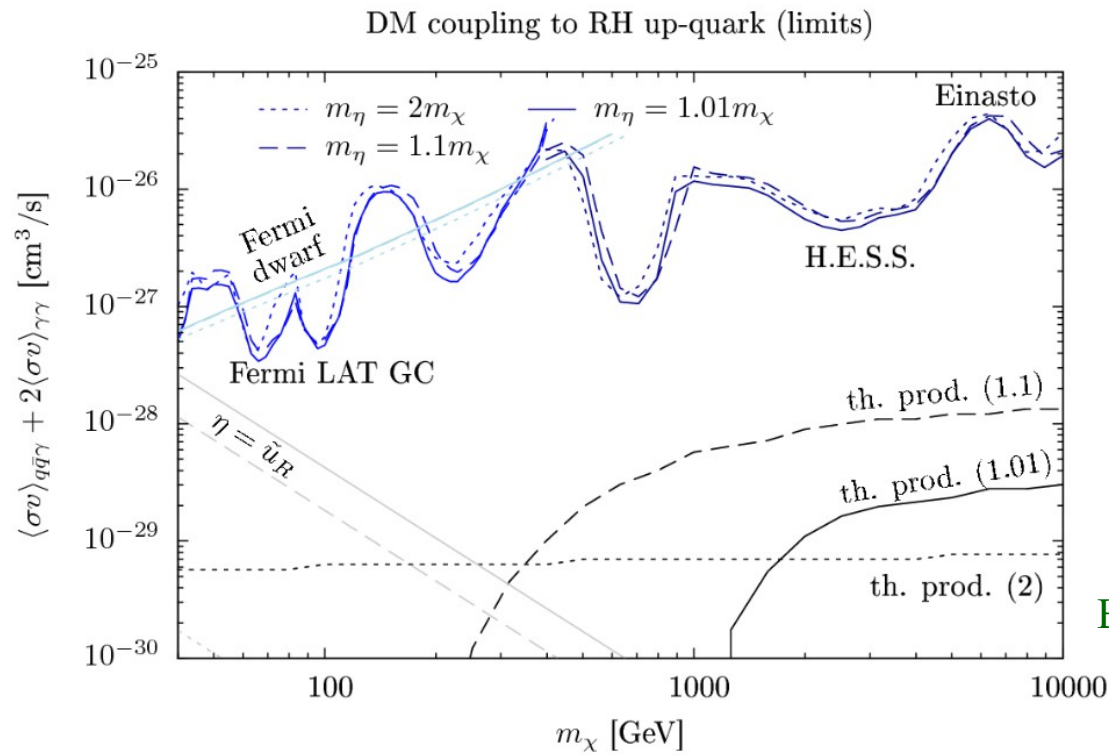
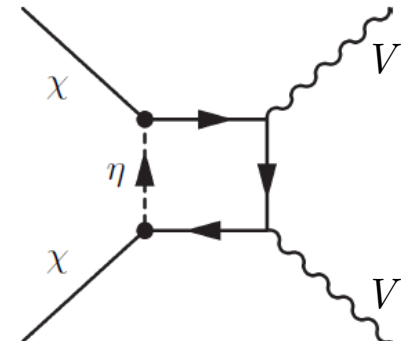
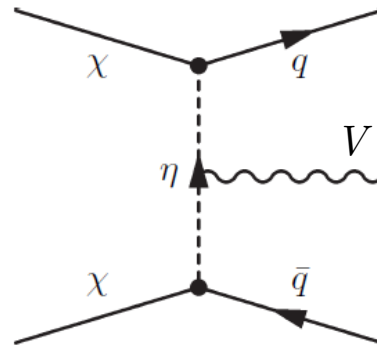
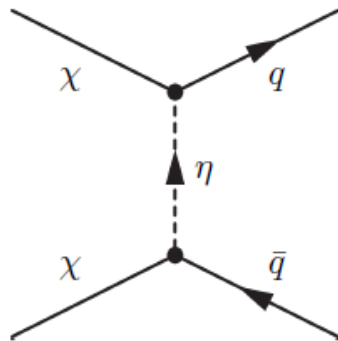
# “Simplified models”: tests

## Indirect detection



# “Simplified models”: tests

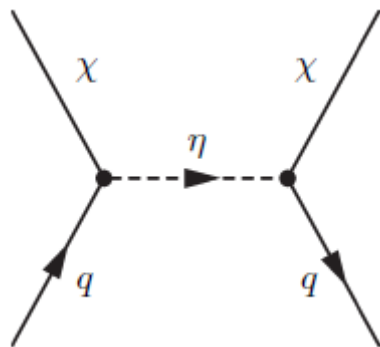
## Indirect detection



Bringmann et al.'12

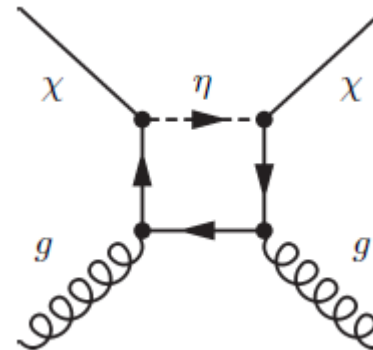
# “Simplified models”: tests

## Direct detection



SI at dim-8  
SD at dim-6

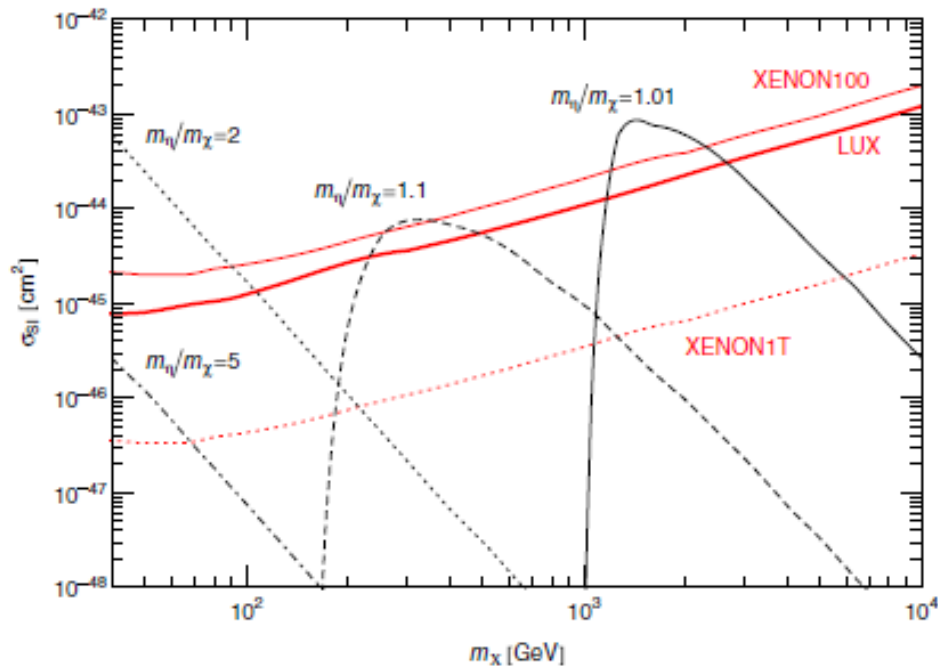
tree-level  
resonant



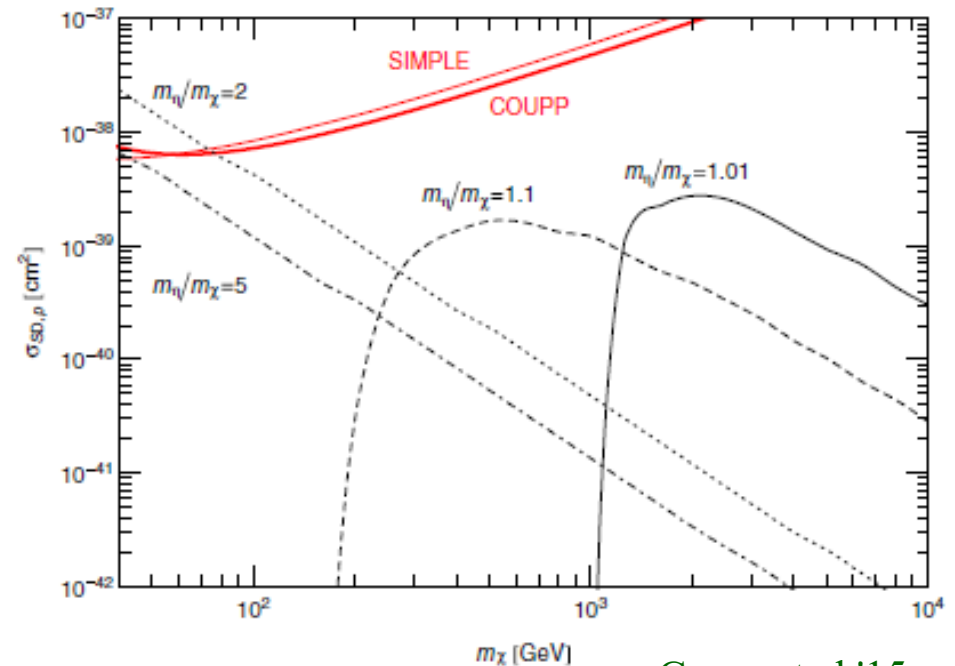
SI at dim-7

one-loop  
non-resonant

DM coupling to u-quark (spin independent)



DM coupling to u-quark (spin dependent proton)

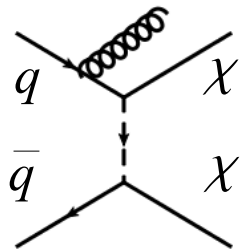


# “Simplified models”: tests

## Collider searches

$$m_\eta \gg m_\chi$$

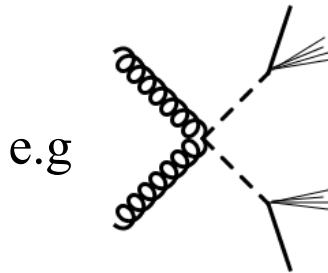
The scalar mediator cannot be produced at the collider, but only the DM.



Mono-X + missing  $p_T$

$$m_\eta = O(m_\chi)$$

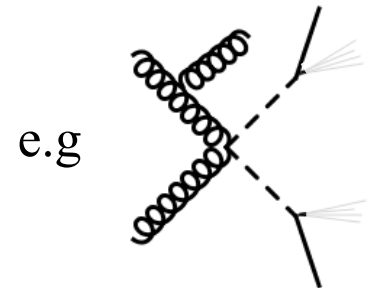
The scalar mediator can be produced at the collider, and then decays into the DM plus a quark.



jets + missing  $p_T$

$$m_\eta \simeq m_\chi$$

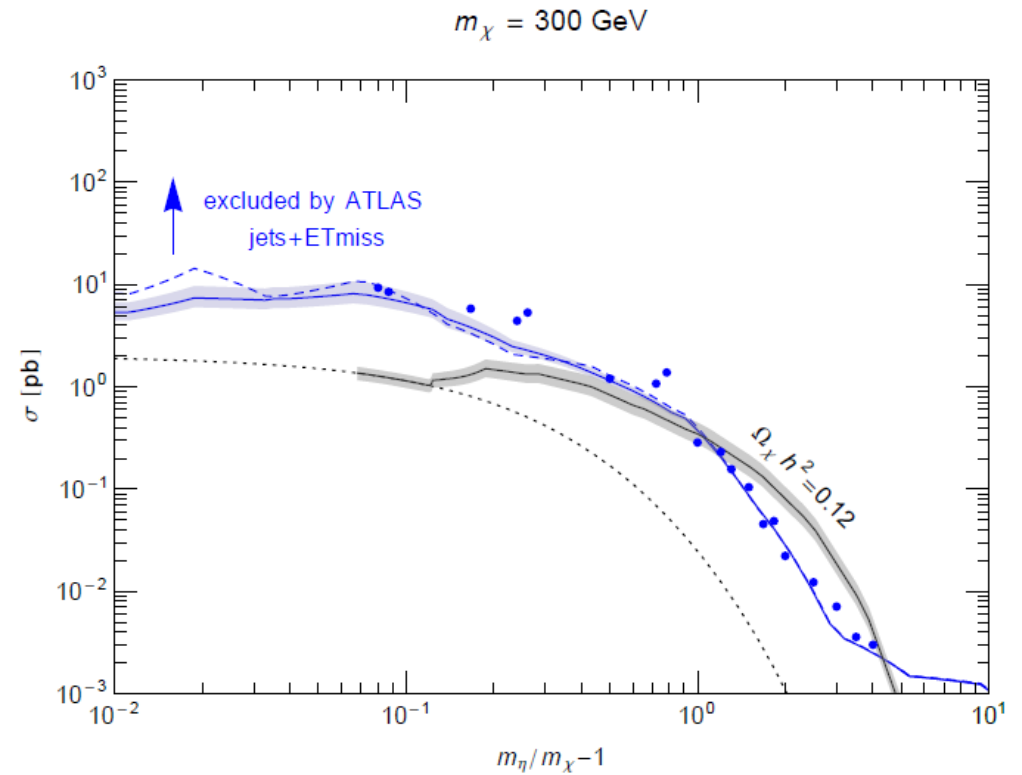
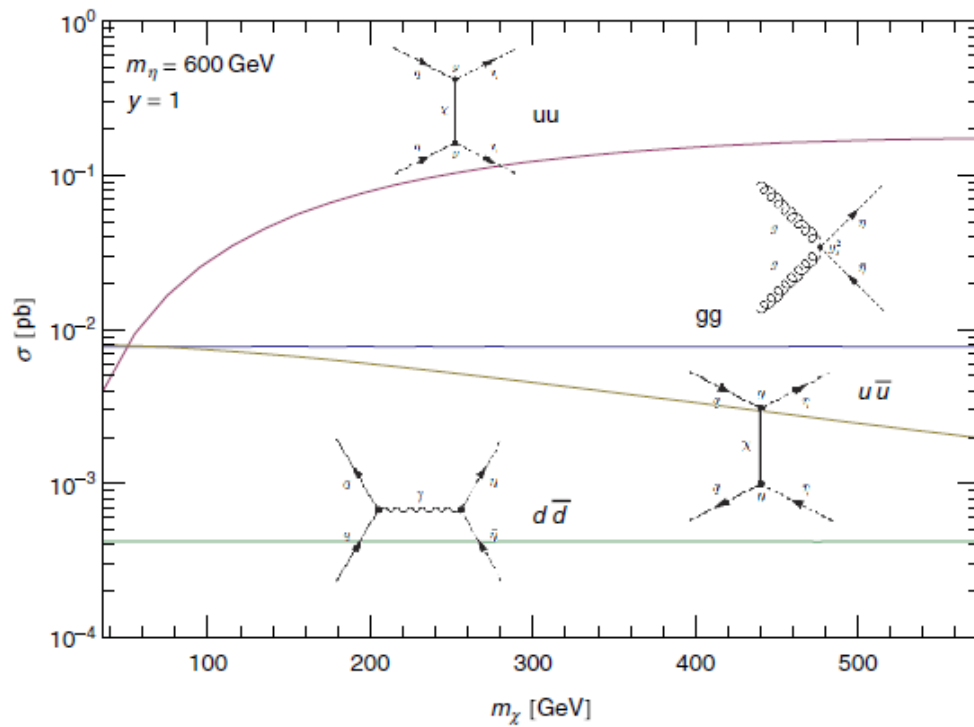
The scalar mediator can be produced at the collider, and then decays into the DM plus a quark. However the jet is too soft to be detected.



jets + missing  $p_T$

# “Simplified models”: tests

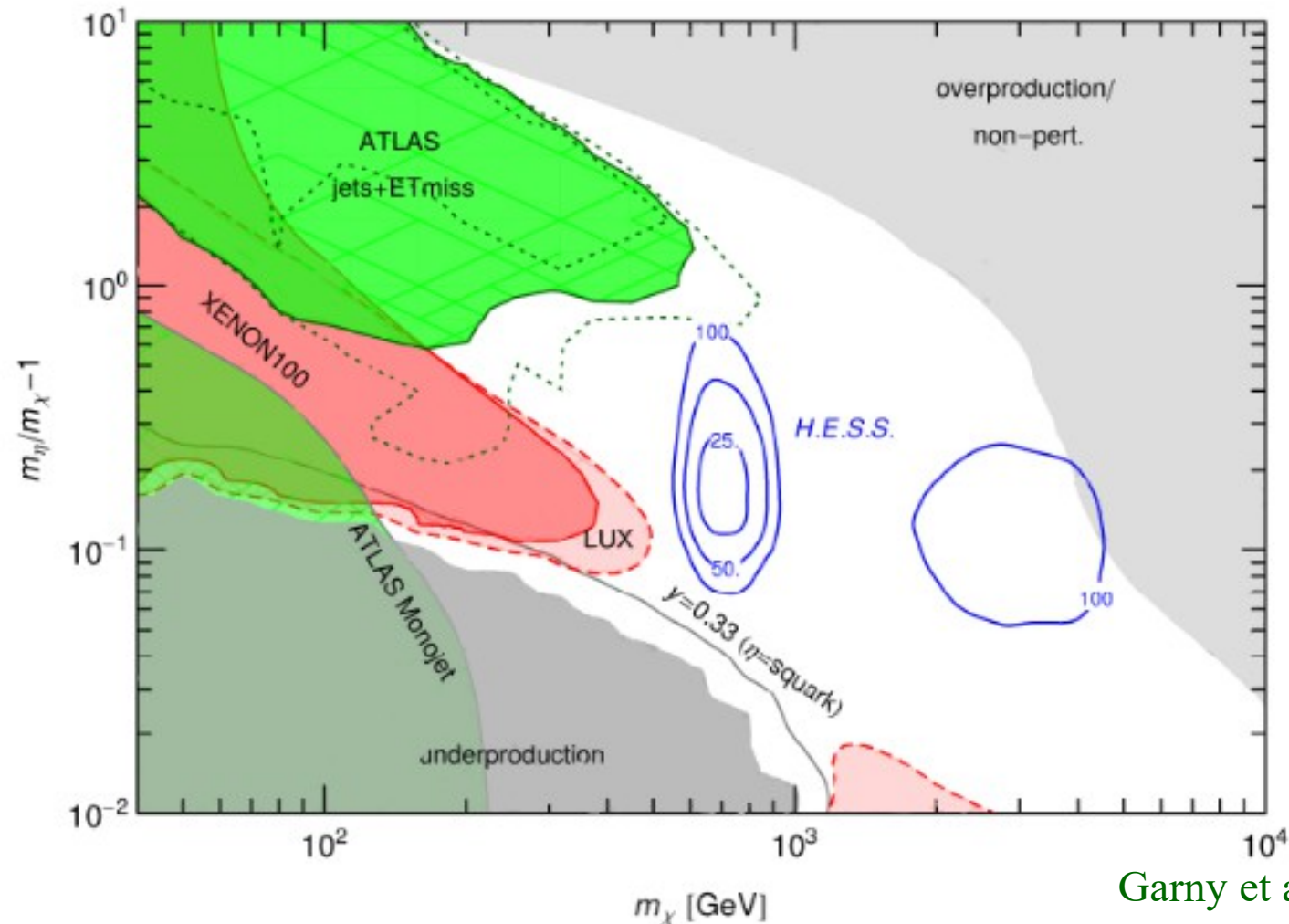
## Collider searches



# “Simplified models”: tests

## Complementarity of searches

DM coupling to RH up-quarks



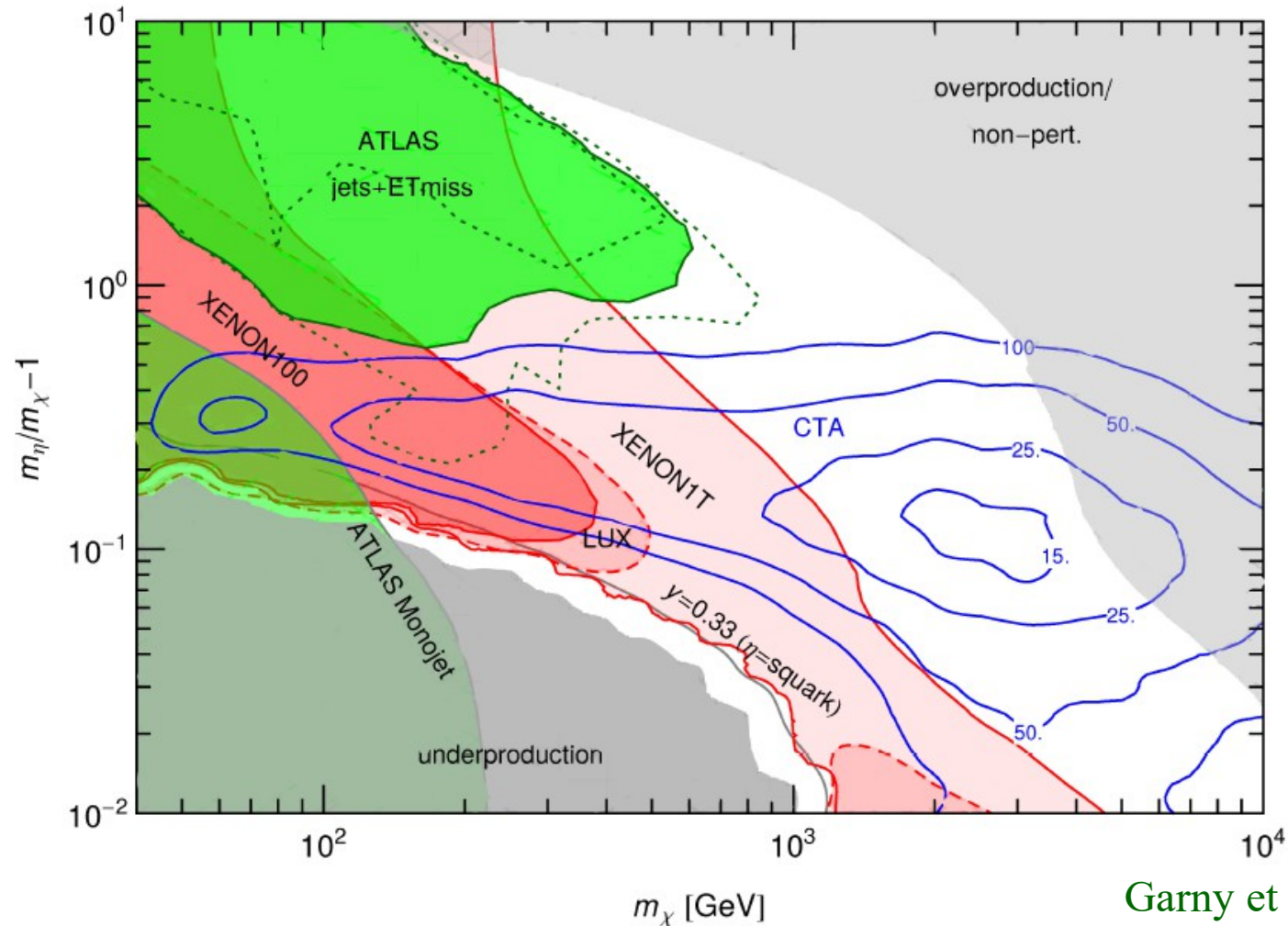
Garny et al.'15



# “Simplified models”: tests

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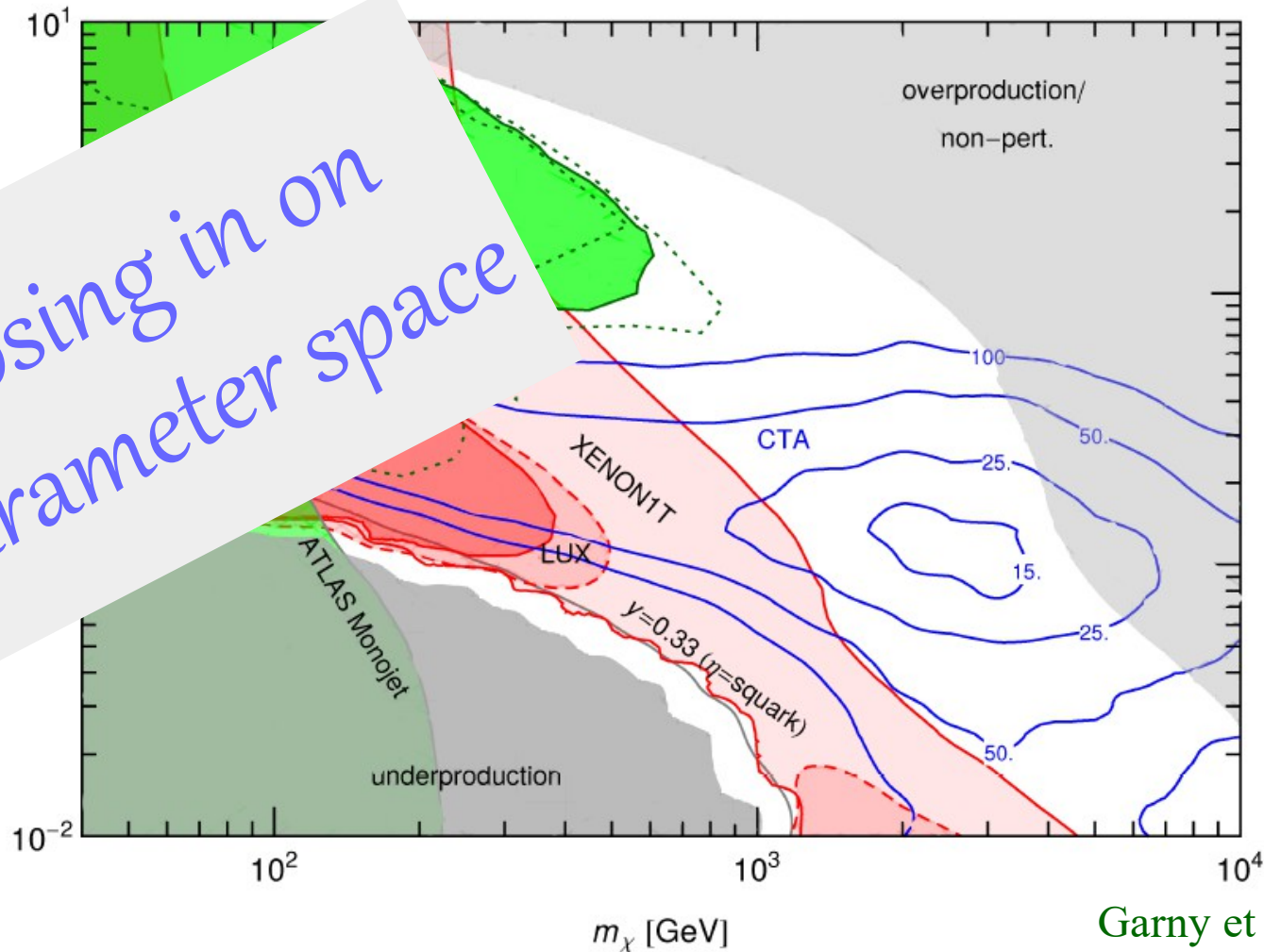
Garny et al.'15

# “Simplified models”: tests

## Complementarity of searches

DM coupling to RH up-quarks

Closing in on  
parameter space

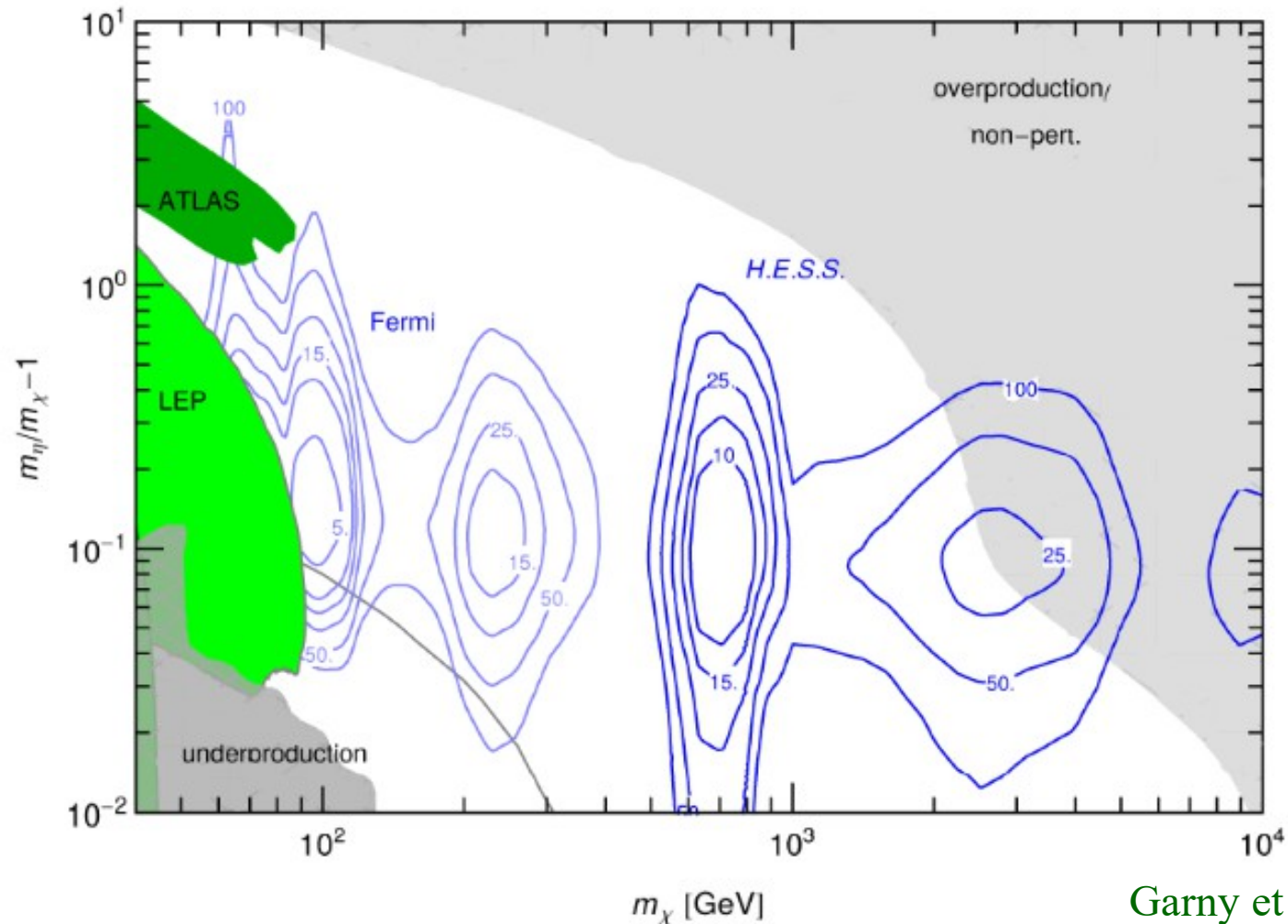


Garny et al.'15

# “Simplified models”: tests

## Complementarity of searches

DM coupling to RH muons

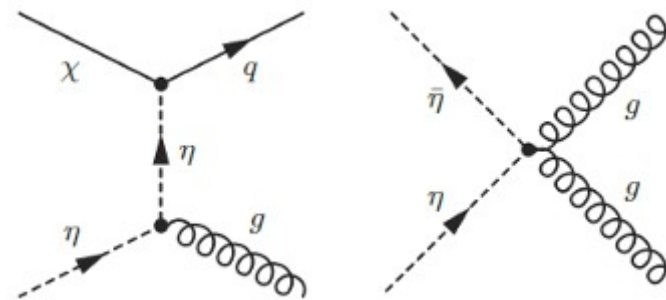
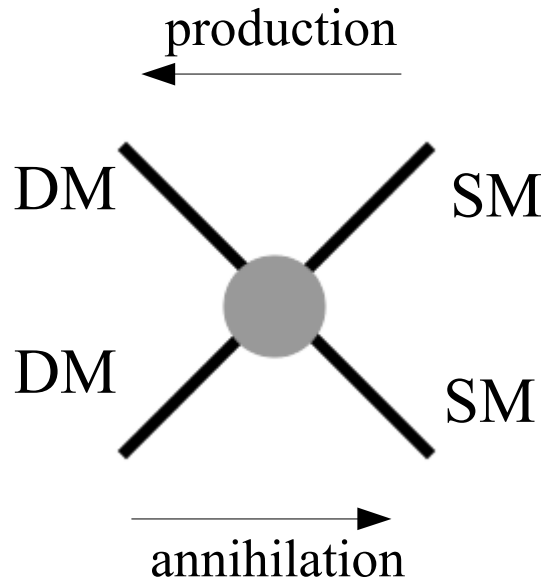


Garny et al.'15

# “Simplified models”: take home lesson

The mediator plays a crucial role in the phenomenology, when its mass is close to the dark matter mass.

## Thermal production

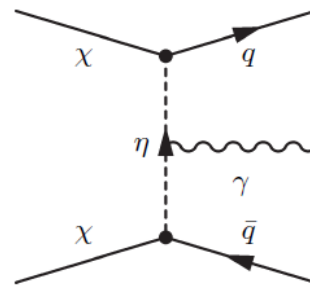
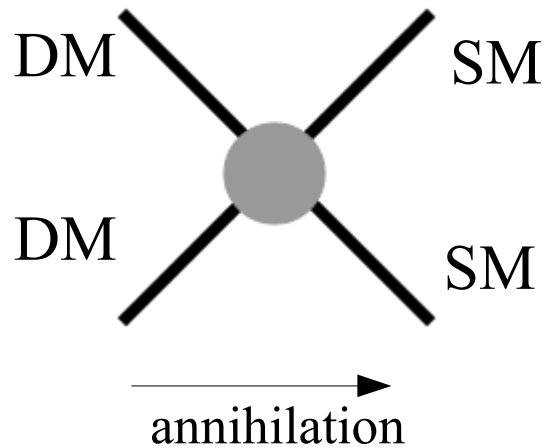


coannihilations

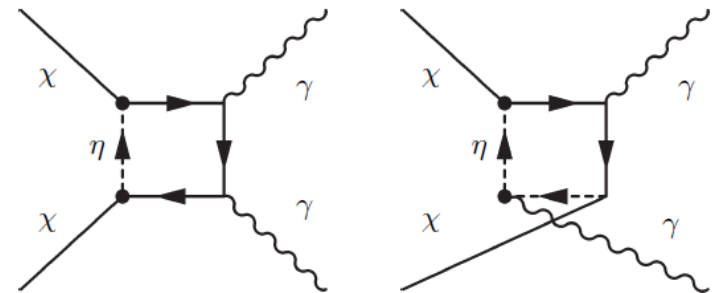
# “Simplified models”: take home lesson

The mediator plays a crucial role in the phenomenology, when its mass is close to the dark matter mass.

## Indirect detection



Internal  
Bremsstrahlung

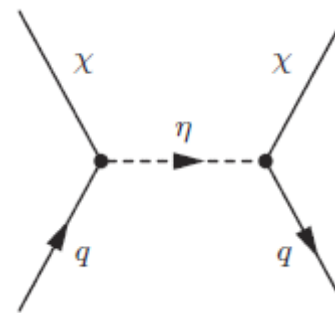
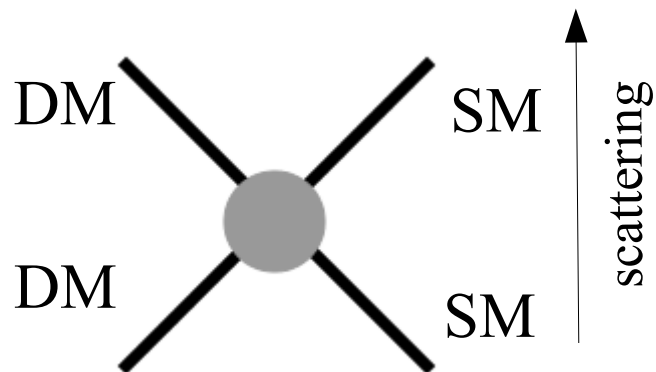


Loop annihilations

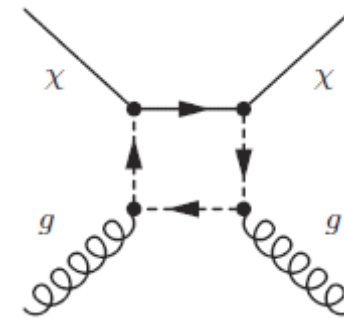
# “Simplified models”: take home lesson

The mediator plays a crucial role in the phenomenology, when its mass is close to the dark matter mass.

## Direct detection



resonance

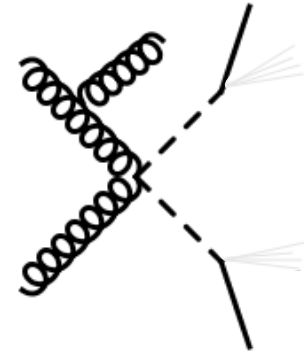
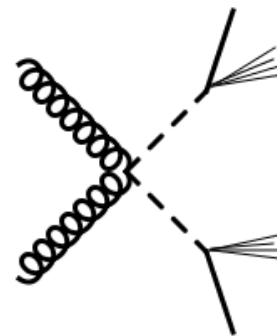
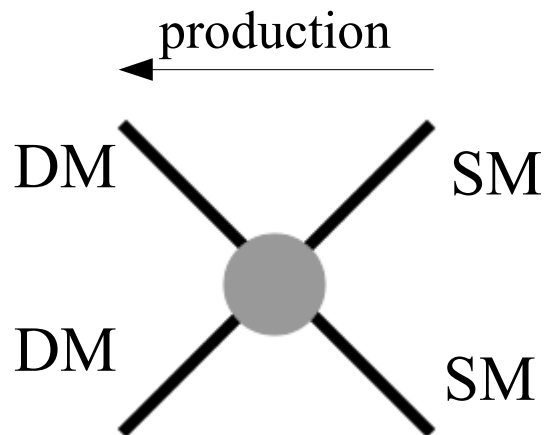


loops

# “Simplified models”: take home lesson

The mediator plays a crucial role in the phenomenology, when its mass is close to the dark matter mass.

## Collider searches



# Conclusions

- Simple particle physics frameworks contain a dark matter candidate with an abundance today which is in qualitative good agreement with observations. These models contain very few parameters and have some predictive power.
- Current experiments have achieved sufficient sensitivity to put tension on some of the simplest frameworks. From an optimistic point of view, discovery could be around the corner.