

Non-standard annihilations in Scalar Dark Matter Models

Tarak Nath Maity

**IIT Kharagpur
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Based on :

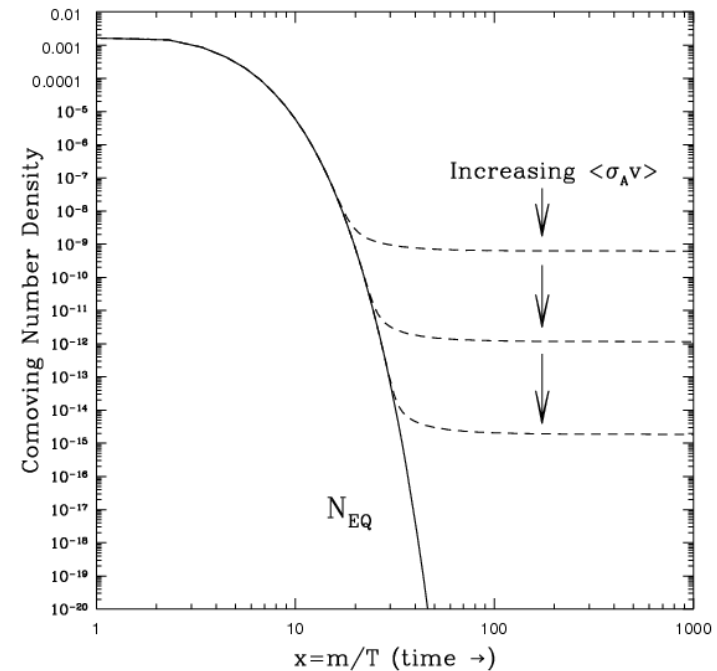
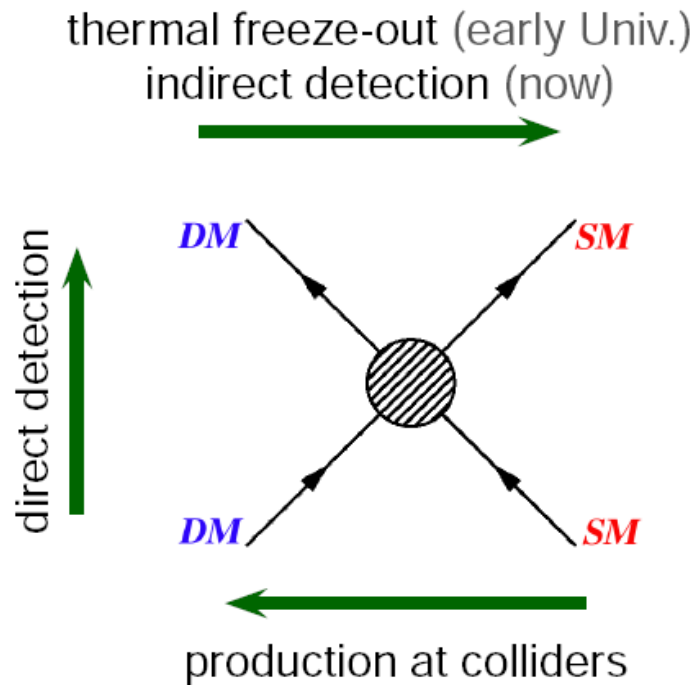
**JCAP 1703 (2017) 045 with U K Dey and T S Ray
and JHEP 10 (2017) 088 with S Bhattacharya, P Ghosh and T S Ray**

Outline

- Introduction : Status of Minimal Model
- Semi-annihilation, co-annihilation and multipartite feature
- Assisted Annihilation
- Summary

Standard framework and some limitations

➤ Weakly Interacting Massive Particle (WIMP).

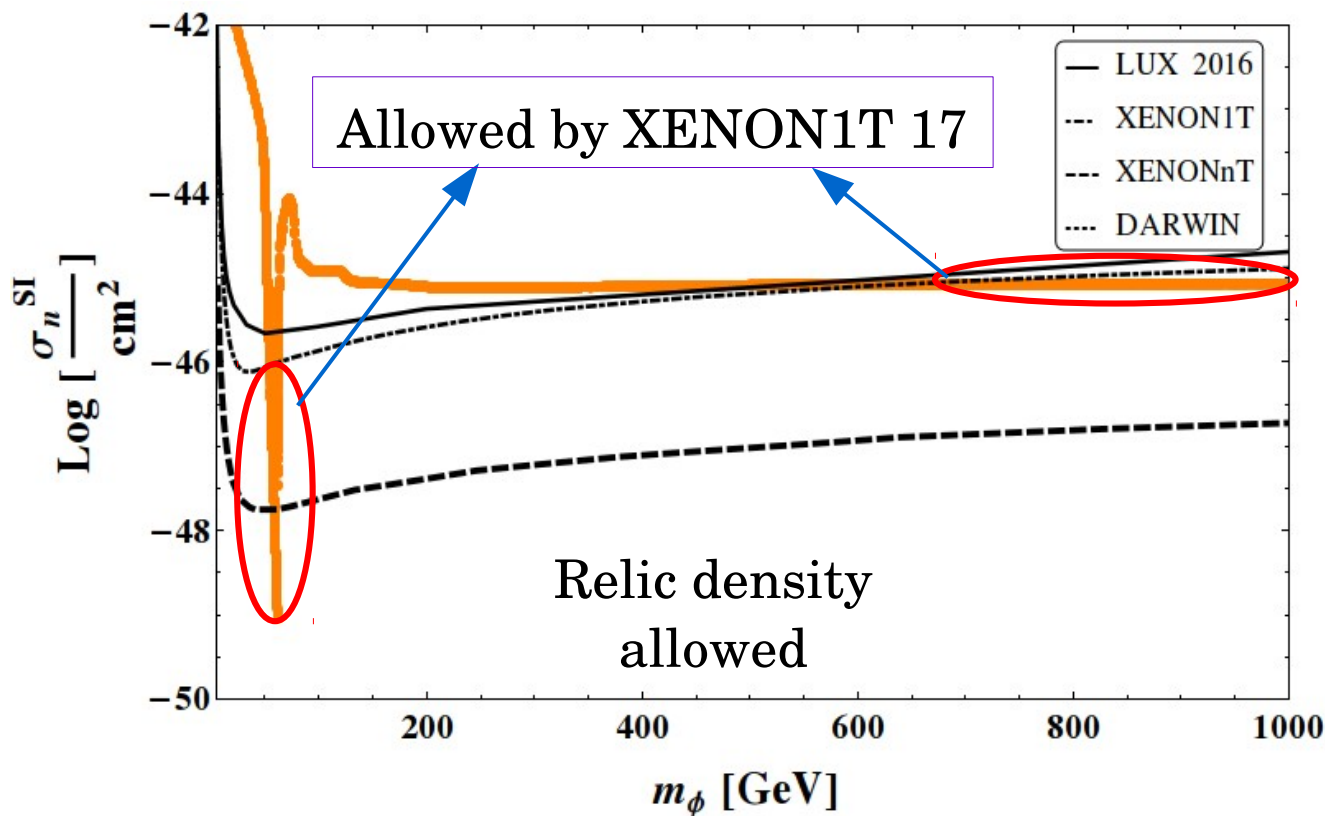


- ✗ No direct search signal.
- ✗ Small scale controversies.

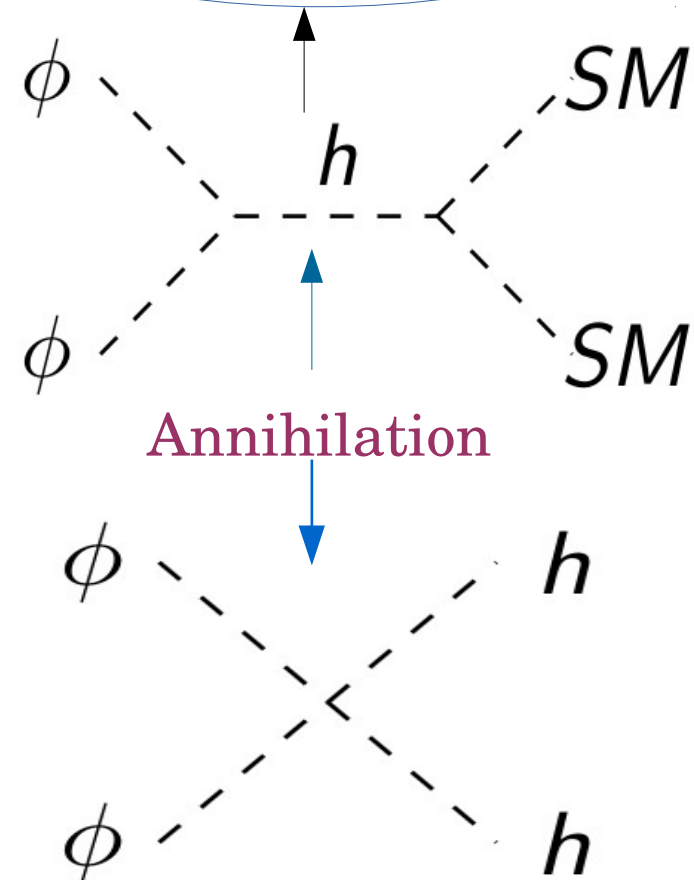
WIMP Higgs portal : Status of minimal model

➤ Real singlet scalar charged under $\mathbb{Z}_2$

$$-\mathcal{L}_{\text{DM-Higgs}} = -\mu_H^2 H^\dagger H + \lambda_H (H^\dagger H)^2 + \frac{1}{2} m_\phi^2 \phi^2 + \frac{\lambda_s}{4!} \phi^4 + \frac{1}{2} \lambda_h H^\dagger H \phi^2.$$



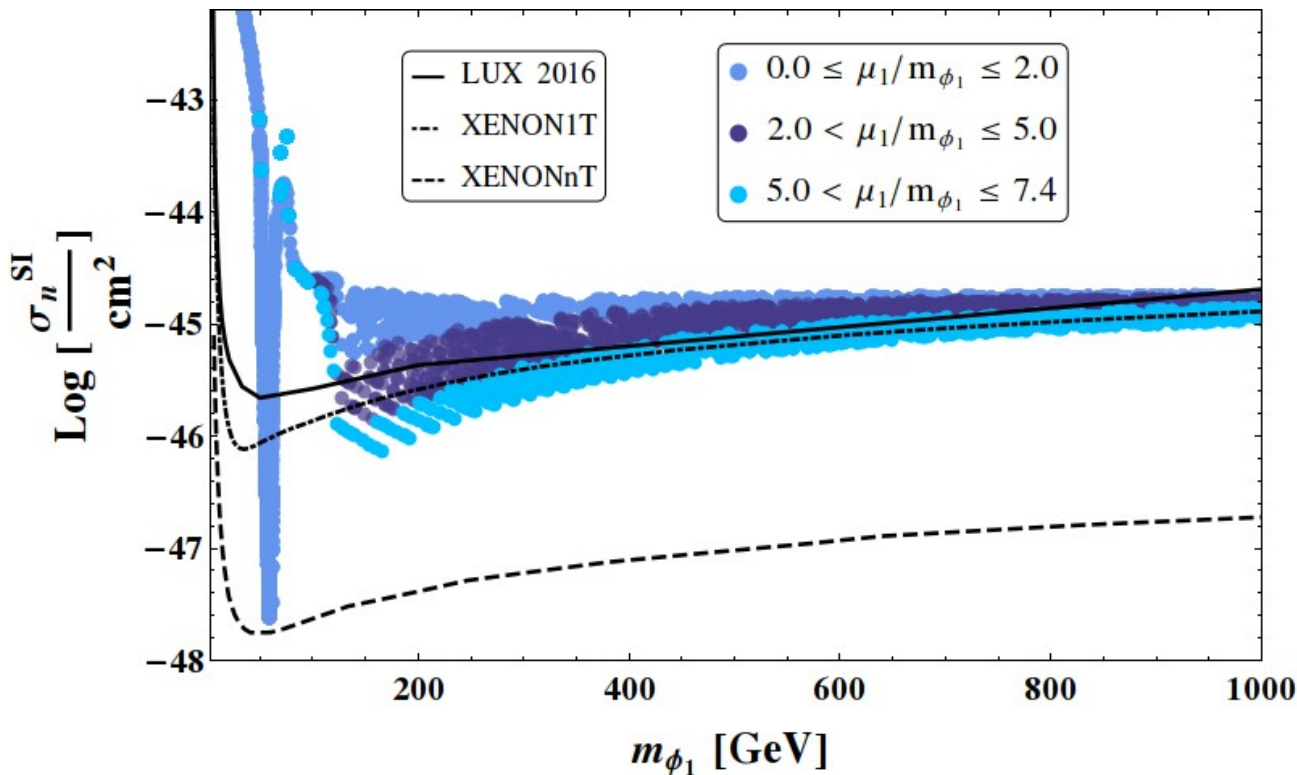
t-channel of this topology sensitive to direct detection



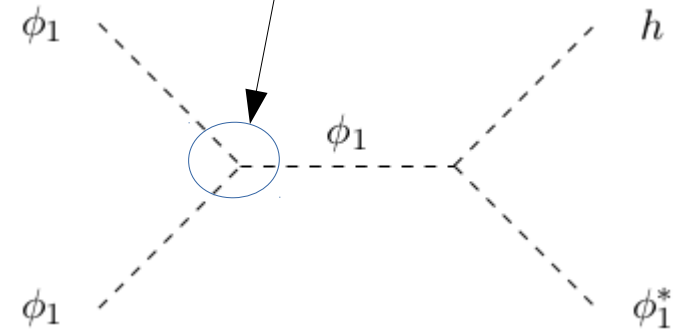
WIMP Higgs portal : Complex scalar under $\mathbb{Z}_3$

➤ Complex Singlet scalar charged under $\mathbb{Z}_3$

$$\begin{aligned}
 -\mathcal{L}_{DM-Higgs} = & -\mu_H^2(H^\dagger H - \frac{v^2}{2}) + \lambda_H(H^\dagger H - \frac{v^2}{2})^2 + m_{\phi_1}^2 \phi_1^* \phi_1 + \frac{\mu_1}{3!}(\phi_1^3 + \text{h.c.}) \\
 & + \lambda_{1s}(\phi_1^* \phi_1)^2 + \lambda_{1h}(\phi_1^* \phi_1)(H^\dagger H - \frac{v^2}{2}).
 \end{aligned}$$



Constrained from
stability of the potential



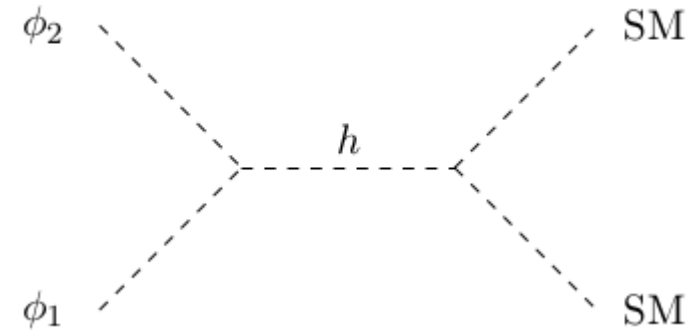
Semi- annihilation

No DD
signature

Multipartite $\mathbb{Z}_2$

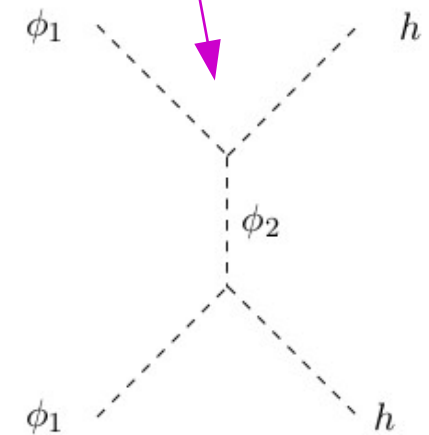
- Two non-degenerate particle ϕ_1 , ϕ_2 under same $\mathbb{Z}_2$ symmetry.
- Annihilation, mediated-annihilation, Co-annihilation operative here.

$m_{\phi_1} < m_{\phi_2} \rightarrow \phi_1$ **DM candidate**

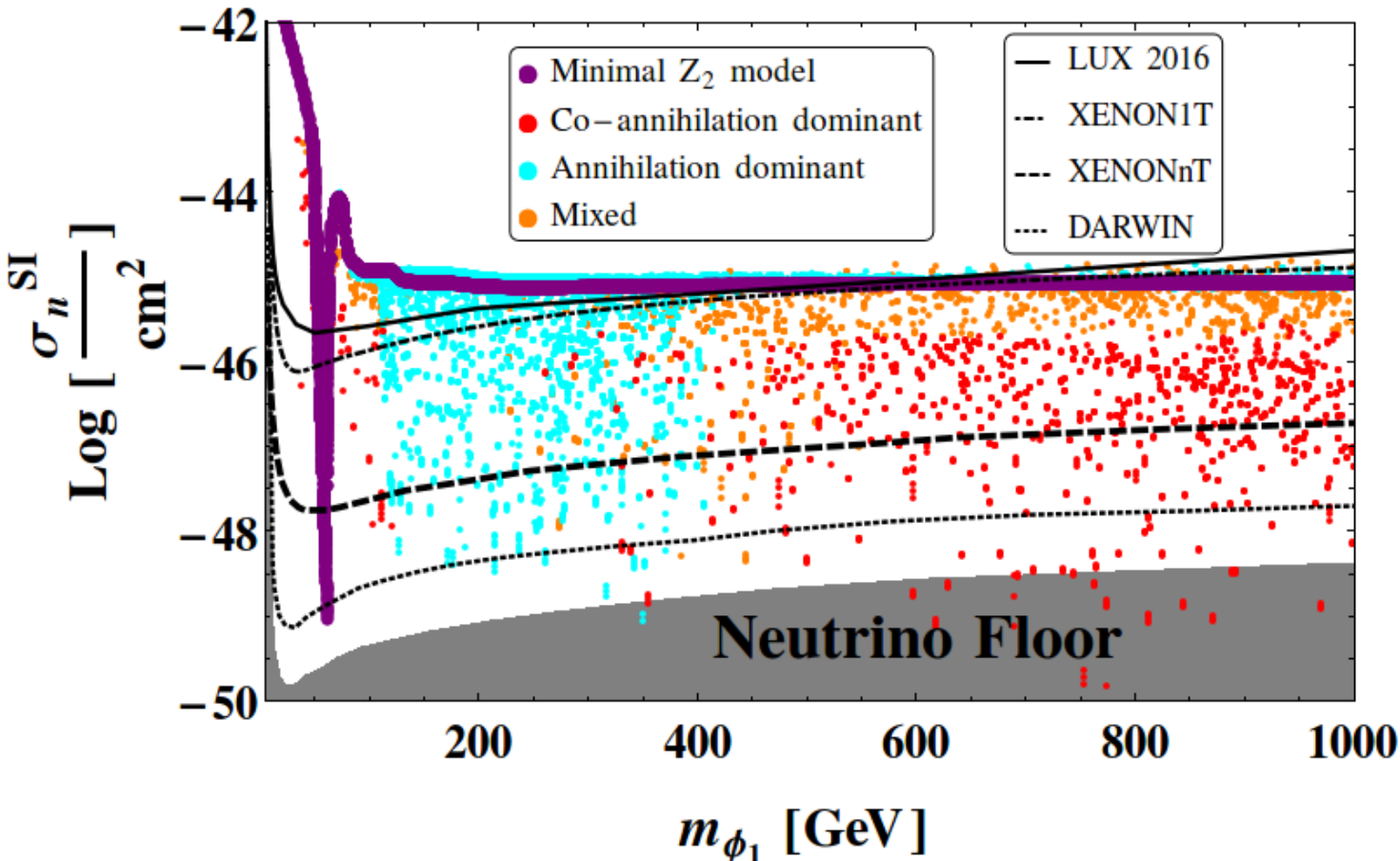


Co-annihilation

Direct detection insensitive

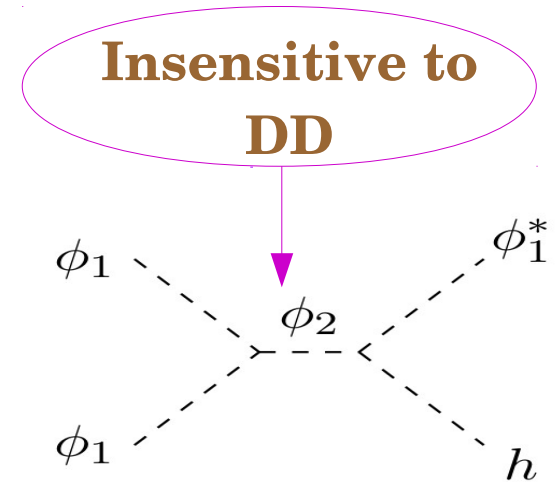


Mediated-annihilation



Multipartite $\mathbb{Z}_3$

- Two non-degenerate particle ϕ_1 , ϕ_2 charged under same $\mathbb{Z}_3$ symmetry.
- Anihilation, co-annihilation, semi-annihilation and Mediated-annihilation operative here.

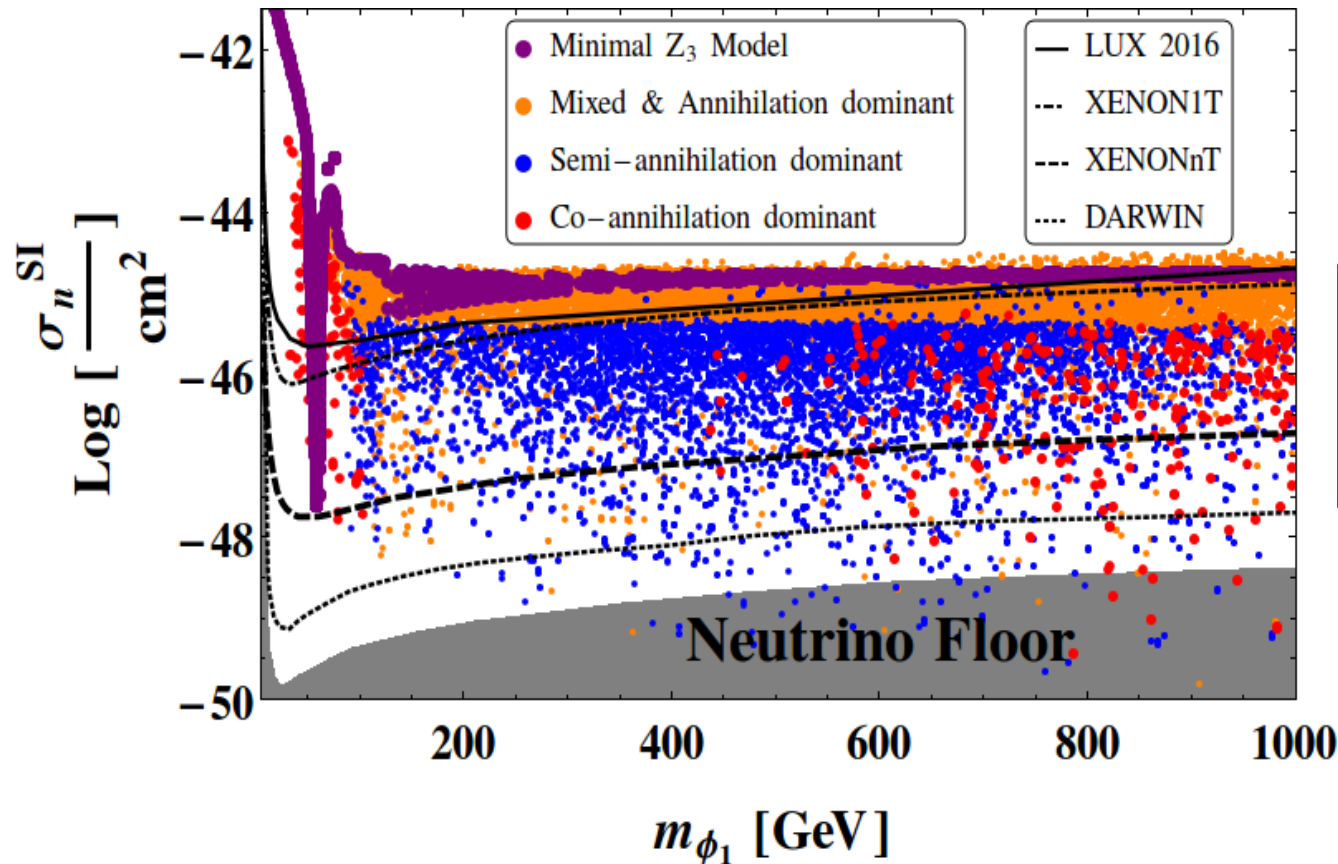


Resonant Semi-annihilation

Resonance Condition

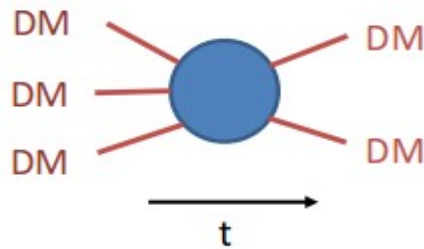
$$m_{\phi_2} = 2m_{\phi_1}$$

$$m_{\phi_1} < m_{\phi_2} \rightarrow \phi_1 \text{ DM}$$



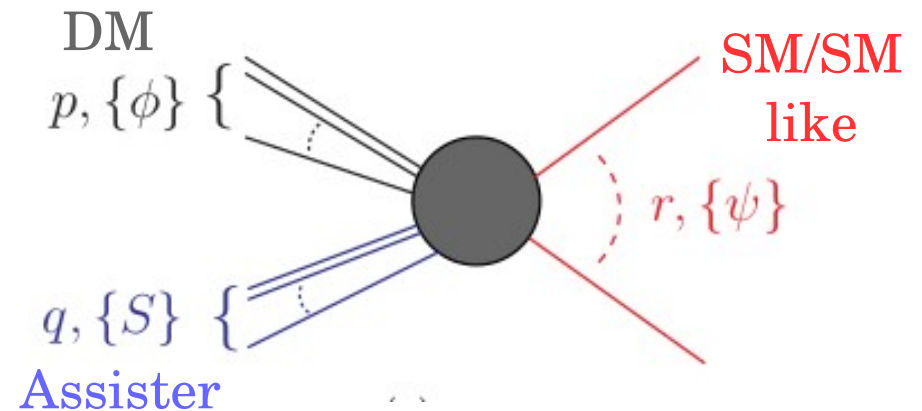
Beyond $2 \rightarrow 2$

- Recently, processes like $3(\text{DM}) \rightarrow 2(\text{DM})$ has been considered. With very tiny interaction with visible sector. [Hochberg et al. PRL 113 171301]

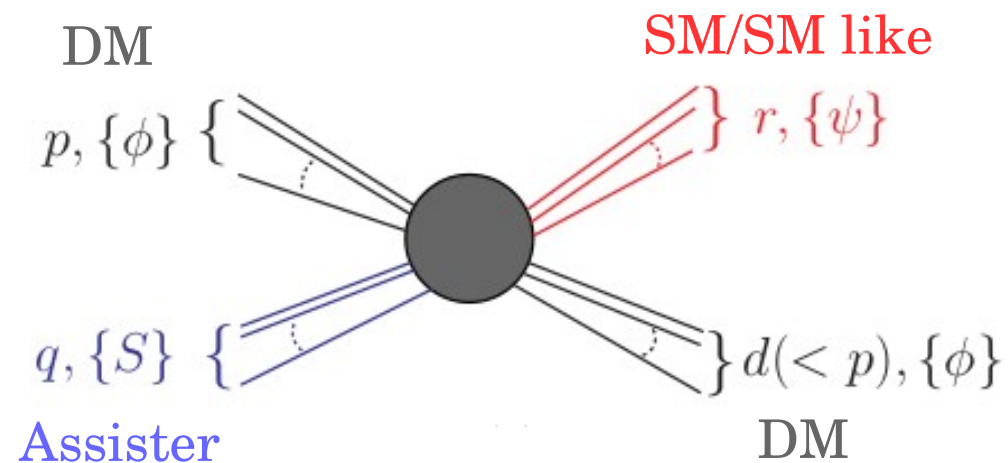


- Topology like $4(\text{DM}) \rightarrow 2(\text{DM})$ also considered. [Bernal et al. JCAP 1504 012 etc.]

- Assisted annihilation



- Assisted Semi-annihilation



Novel features of Assisted Annihilation

$$\begin{array}{ccc} \text{DM} & \text{Assister} & \text{SM/SM like} \\ \underbrace{\phi, \dots, \phi}_p, \underbrace{S, \dots, S}_q & \rightarrow & \underbrace{\psi, \dots, \psi}_r \end{array}$$

- For simplicity, let's assume $r = 2$

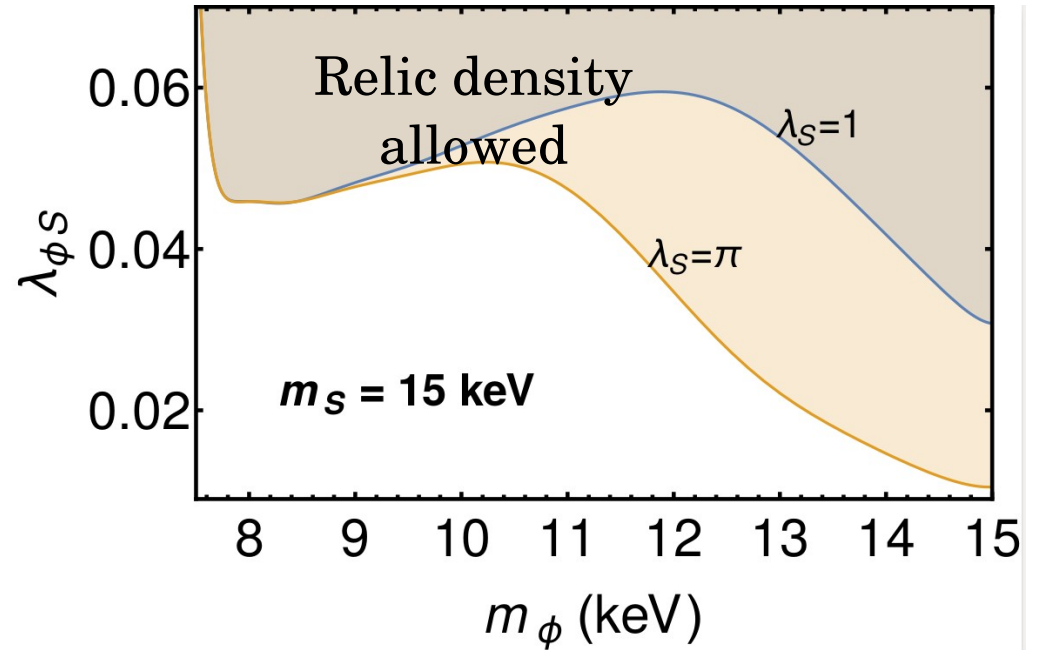
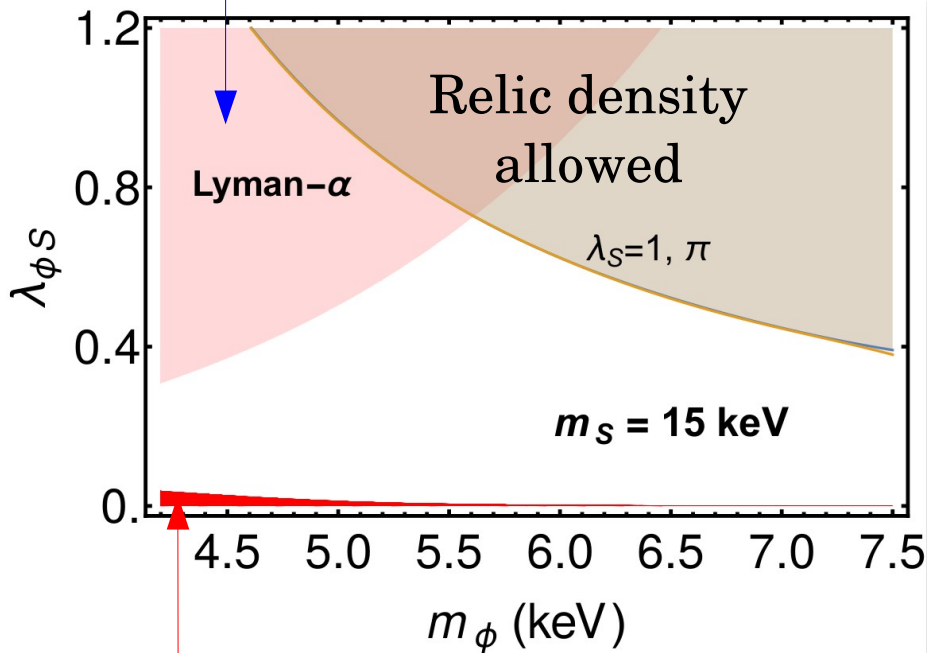
$$\langle \sigma_{\mathbf{p}+\mathbf{q} \rightarrow 2\mathbf{v}^{\mathbf{p}+\mathbf{q}-1}} \rangle \sim \left[\mathbf{M}^{-3(\mathbf{p}+\mathbf{q})+4} \right]$$

- This suppression in cross section naturally leads to light DM.
- No direct detection, indirect detection signal.
- Unlike coannihilation here assister is not charged under the symmetry which makes DM stable.
- Region of interest : $\mathbf{m}_\phi < \mathbf{m}_S, \mathbf{m}_\psi$

$$\frac{dY_\phi}{dx} = -\frac{x s^{p+q-1}}{H(m_\phi)} N_{\text{Bolt}} \langle \sigma v^{p+q-1} \rangle \left[(Y_\phi)^p (Y_\phi^{\text{eq}})^q - (Y_\phi^{\text{eq}})^{p+q} \right], \quad N_{\text{Bolt}} = e^{qx(1-\epsilon)} \epsilon^{3q/2}$$

4 → 2: keV scale DM

$$T_f < 1 \text{ keV}$$



Semi-relativistic free out

$$\mathcal{L}^{\text{dark}} \supset \underbrace{\frac{1}{2} m_\phi^2 \phi^2}_{\text{DM}} + \underbrace{\frac{1}{2} m_S^2 S^2}_{\text{Assister}} + \frac{\lambda_{\phi S}}{4} \phi^2 S^2 + \frac{\lambda_\phi}{4!} \phi^4 + \frac{\lambda_S}{4!} S^4$$

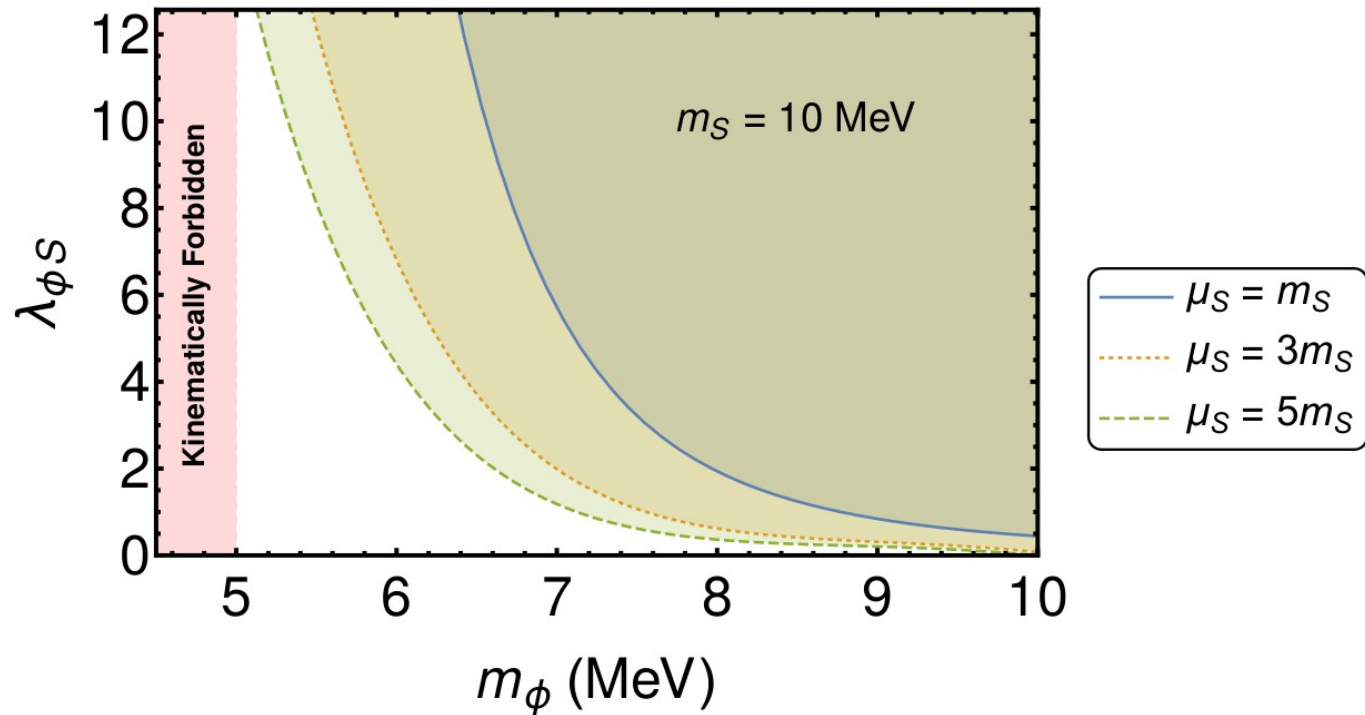
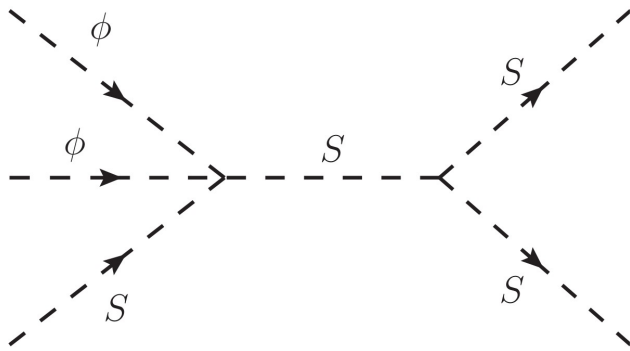
$m_\phi < m_S$

There will be strong constraint on such low mass DM from BBN & CMB.
(In preparation with U K Dey and T S Ray)

3 → 2 : MeV scale DM

➤ Lets introduce a term

$$\mathcal{L}_{3 \rightarrow 2} = \frac{\mu_S}{3!} S^3.$$



$$\frac{m_S}{2} < m_{\phi} < m_S$$

Relic density allowed parameter space

Summary

- Minimal scalar singlet DM models are considerably constrained by recent results of direct detection experiments.
- We show that the non-standard annihilation topologies like co-annihilation, semi-annihilation etc. play crucial role to evade such constraint.
- The interplay of this kind process yields plenty of region in the parameter space that is unconstrained by present and future direct detection experiments.
- Going beyond $2 \rightarrow 2$ scenario, we propose a non-standard annihilation mechanism “assisted annihilation” by which we can get sub-GeV dark matter.
- Such a light cold dark matter with self interaction may solve small scale structure formation issue.

Thank
You

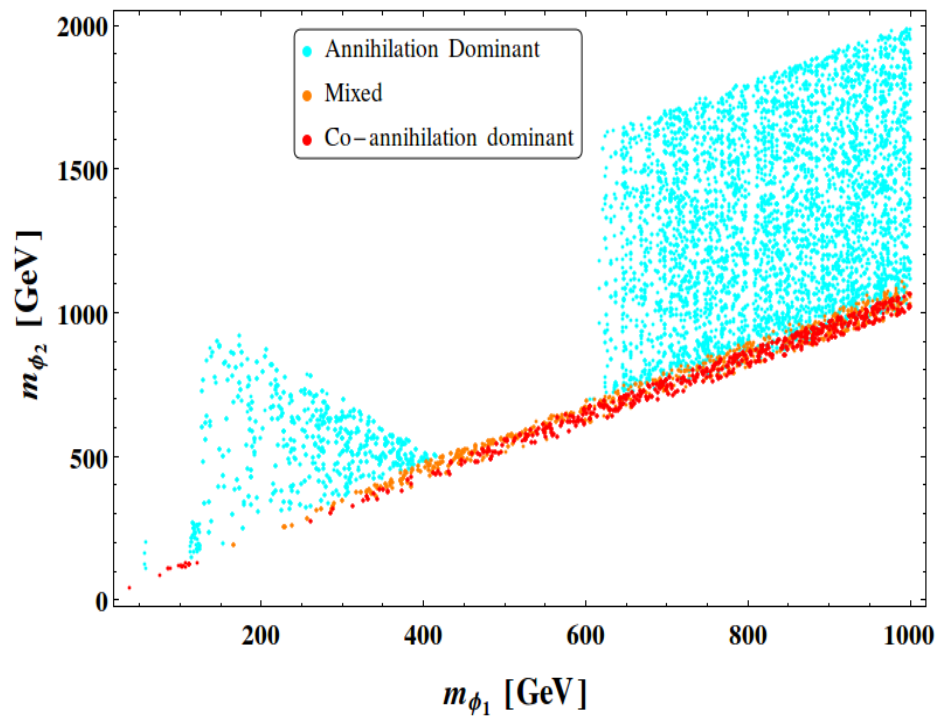
Lagrangian: Multiparticle $\mathbb{Z}_2$

$$\begin{aligned} -\mathcal{L}_{DM-Higgs} = & -\mu_H^2 \left(H^\dagger H - \frac{v^2}{2} \right) + \lambda_H \left(H^\dagger H - \frac{v^2}{2} \right)^2 \\ & + \frac{1}{2} m_{\phi_1}^2 \phi_1^2 + \frac{1}{2} m_{\phi_2}^2 \phi_2^2 + \frac{\lambda_{e_1}}{4} \phi_1^2 \phi_2^2 + \frac{\lambda_{e_2}}{3!} \phi_1^3 \phi_2 + \frac{\lambda_{e_3}}{3!} \phi_1 \phi_2^3 \\ & + \frac{\lambda_{1s}}{4!} \phi_1^4 + \frac{\lambda_{2s}}{4!} \phi_2^4 + \frac{1}{2} \lambda_{1h} \phi_1^2 \left(H^\dagger H - \frac{v^2}{2} \right) \\ & + \frac{1}{2} \lambda_{2h} \phi_2^2 \left(H^\dagger H - \frac{v^2}{2} \right) + \lambda_{12h} \phi_1 \phi_2 \left(H^\dagger H - \frac{v^2}{2} \right). \end{aligned}$$

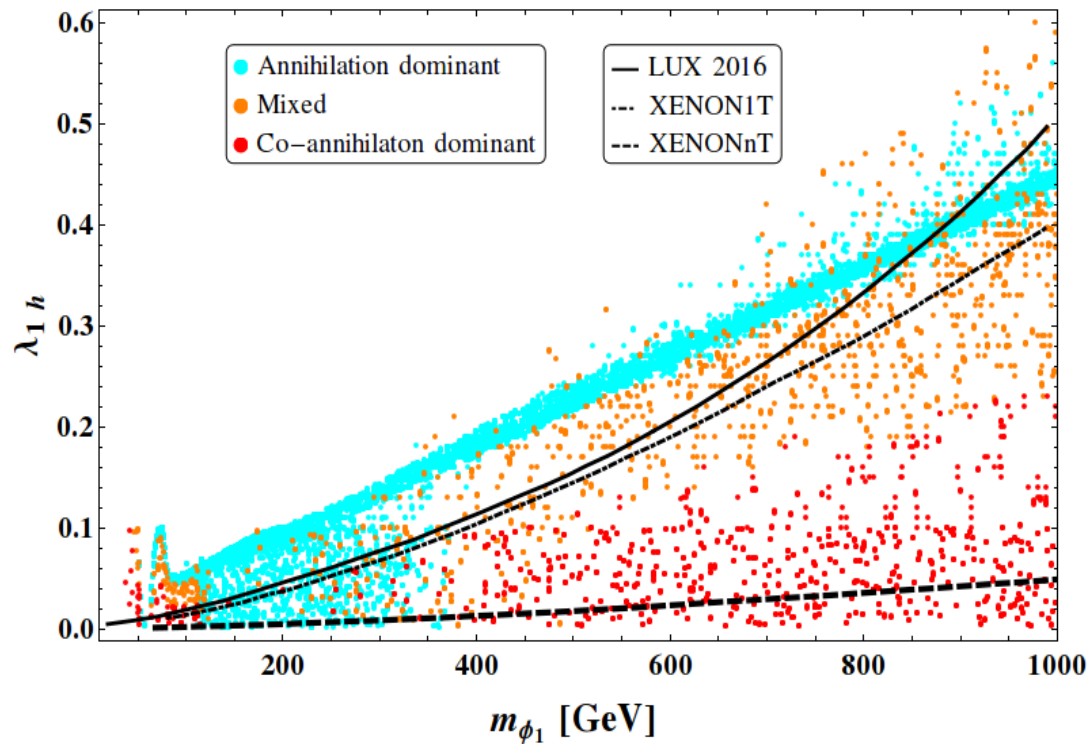
Parameter Space:

$$\begin{aligned} 1 \text{ GeV} \leq m_{\phi_1} \leq 1000 \text{ GeV}, \quad 2 \leq \Delta m \equiv m_{\phi_2} - m_{\phi_1} \\ \leq 1000 \text{ GeV}, \quad 0.001 \leq \lambda_{1h} = \lambda_{2h} \leq 0.1, \quad 0 \leq \lambda_{12h} \leq 1 \end{aligned}$$

Multiparticle : $\mathbb{Z}_2$



Relic density and LUX 2016
allowed points in m_{ϕ_1} - m_{ϕ_2}
plane



Relic density allowed points in
 m_{ϕ_1} - λ_{1h} plane

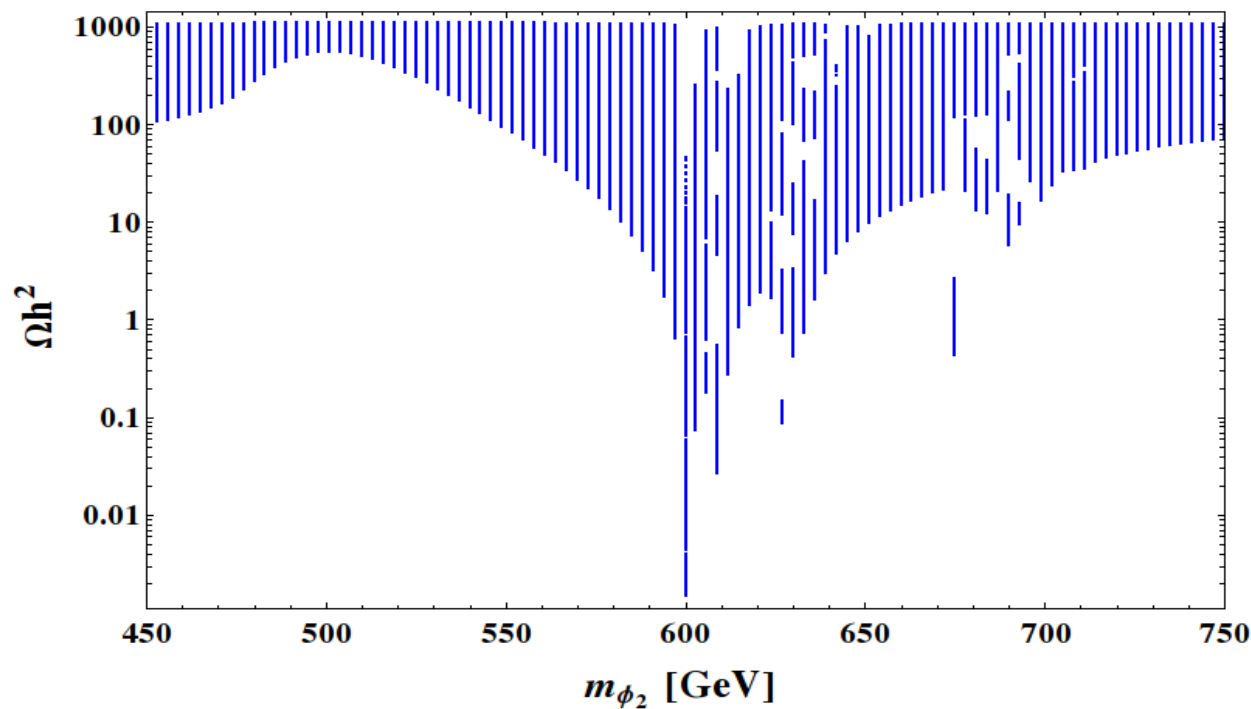
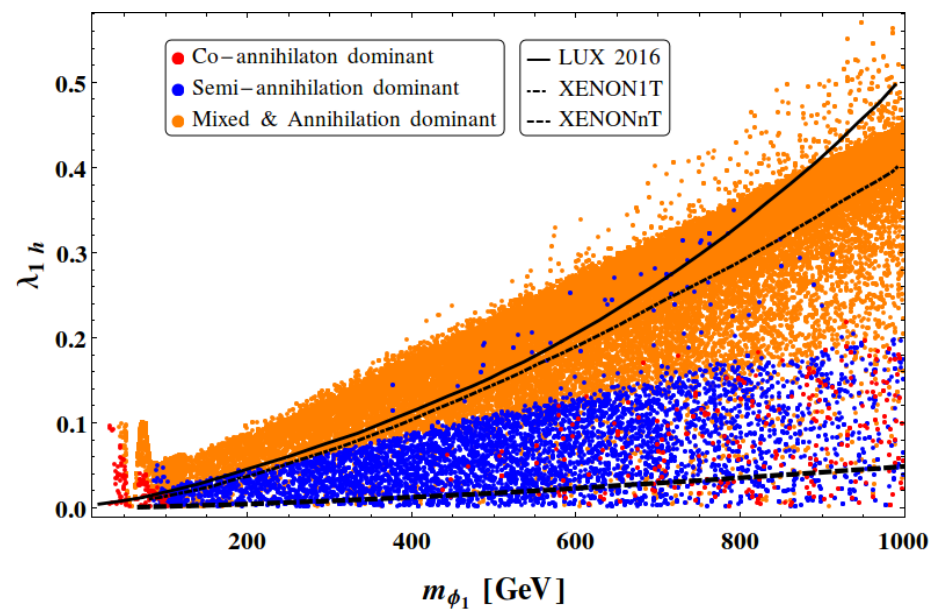
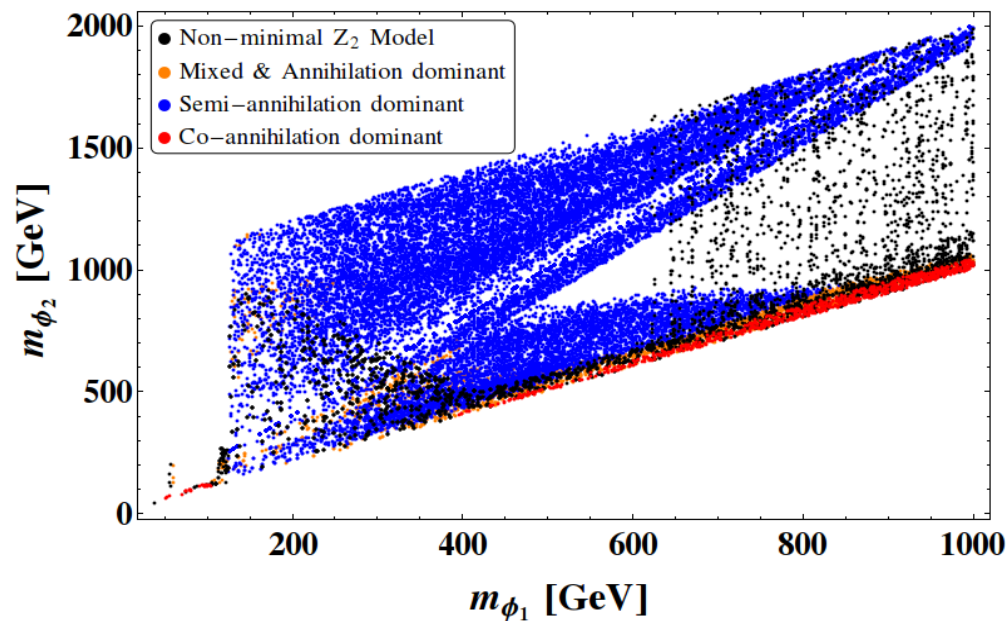
Lagrangian: Multiparticle $\mathbb{Z}_3$

$$\begin{aligned} -\mathcal{L}_{DM-Higgs} = & -\mu_H^2(H^\dagger H - \frac{v^2}{2}) + \lambda_H(H^\dagger H - \frac{v^2}{2})^2 \\ & + m_{\phi_1}^2 \phi_1^* \phi_1 + m_{\phi_2}^2 \phi_2^* \phi_2 + \frac{\mu_1}{3!}(\phi_1^3 + \text{h.c}) + \frac{\mu_2}{3!}(\phi_2^3 + \text{h.c}) \\ & + \frac{\mu_{12}}{2!}(\phi_1^2 \phi_2 + \phi_2^2 \phi_1 + \text{h.c}) + \lambda_{1s}(\phi_1^* \phi_1)^2 + \lambda_{2s}(\phi_2^* \phi_2)^2 \\ & + \lambda_e[(\phi_1^* \phi_1)(\phi_2^* \phi_2) + \{(\phi_1^* \phi_2)^2 + \text{h.c}\}] + \lambda_{1h}(\phi_1^* \phi_1)(H^\dagger H - \frac{v^2}{2}) \\ & + \lambda_{2h}(\phi_2^* \phi_2)(H^\dagger H - \frac{v^2}{2}) + \lambda_{12h}[\phi_1^* \phi_2 + \text{h.c}](H^\dagger H - \frac{v^2}{2}) . \end{aligned}$$

Parameter Space:

$$\begin{aligned} 10 \text{ GeV} \leq m_{\phi_1} \leq 1000 \text{ GeV}, \quad 2 \leq \Delta m \leq 1000 \text{ GeV}, \\ \mu_i \leq 2m_i, \quad \mu_{12} \leq 2m_1, \quad 0.001 \leq \lambda_{1h} = \lambda_{2h} \leq 0.1, \\ 0.1 \leq \lambda_{12h} \leq 1 . \end{aligned}$$

Multiparticle : $\mathbb{Z}_3$



(Top-left) Relic density and xenon 1T allowed points in m_{ϕ_1} - m_{ϕ_2} plane

(Top-right) Relic density allowed points in m_{ϕ_1} - λ_{1h} plane

(Bottom) Resonance plot for $m_{\phi_1} = 300$ GeV