

# Hidden Sectors, String/M theory and Two Predictions On Dark Matter

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25<sup>th</sup> International Supersymmetry Conference  
"SUSY '17"

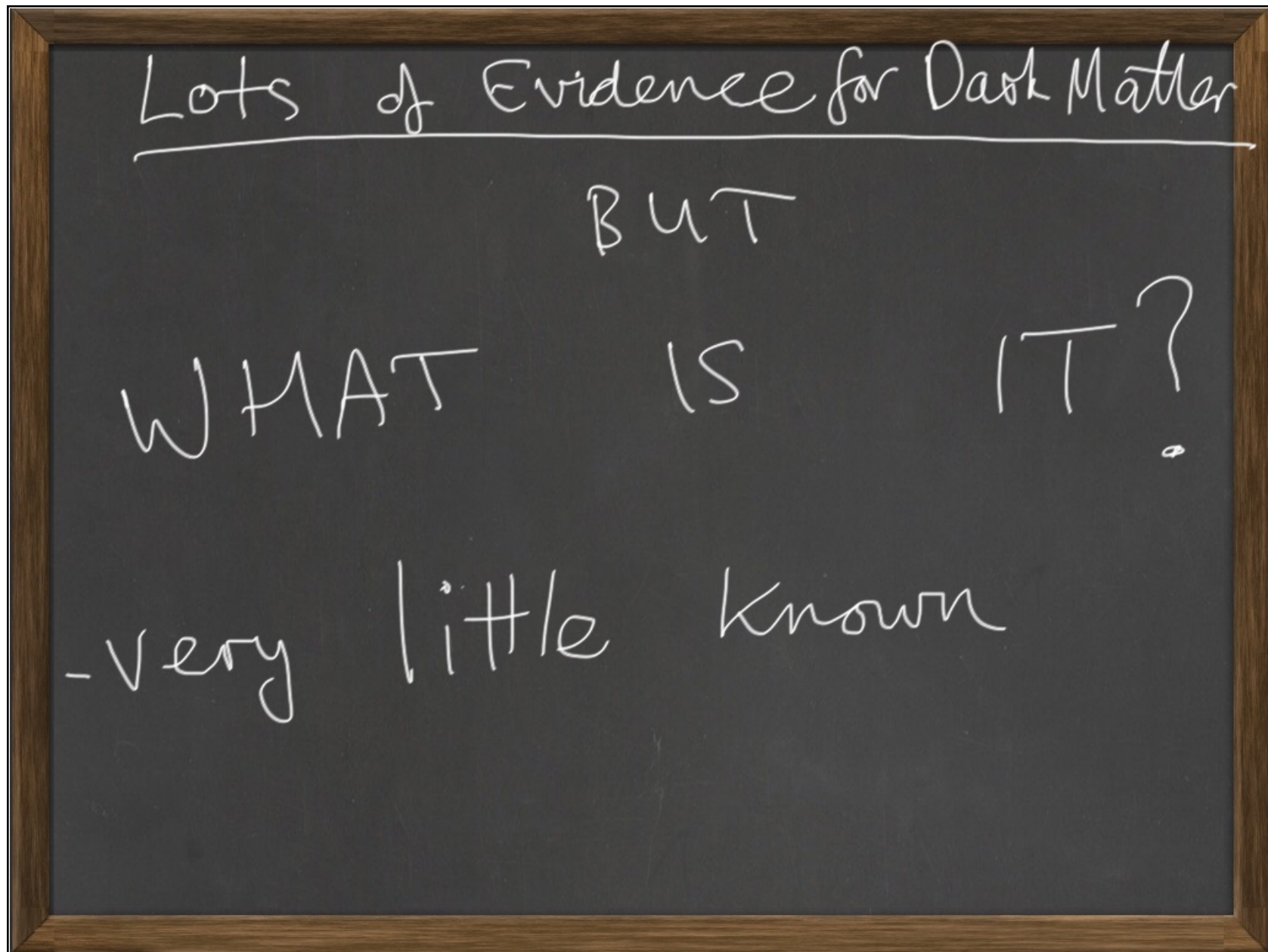
TIFR, Mumbai, 11-15 DEC 2017

A pleasure to be back at Tata  
Institute.

~ 17 years ago, the Strings  
meeting here made a big  
impact on my research and  
life . . . .









- It behaves like matter
  - assume it's particular.
- it doesn't have much charge (could have some)
- doesn't self interact much

That's not very much info.

Spin? Mass?

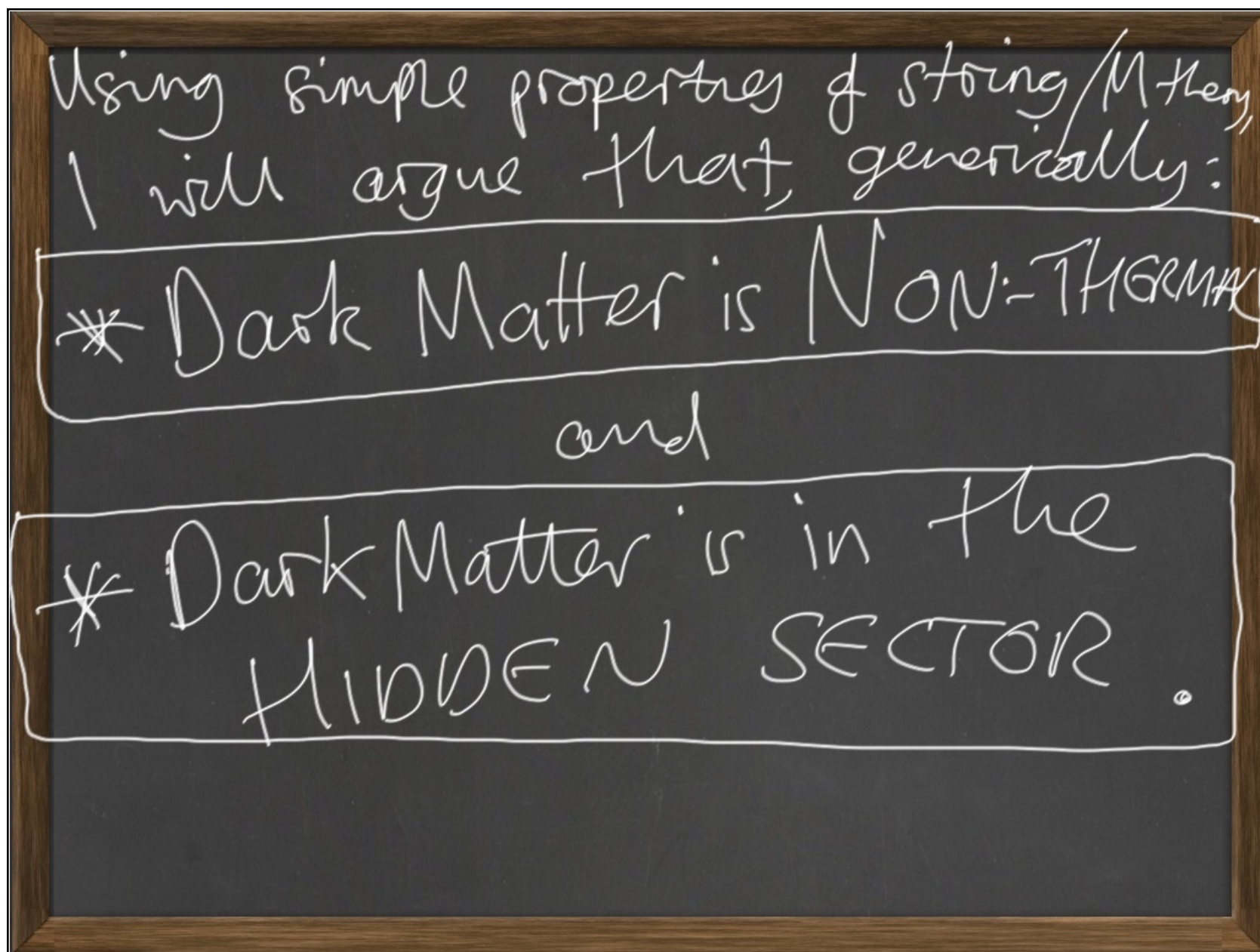
Completely unknown.

How is it produced?

unknown.

When was it produced?

Before BBN most likely?





Based on work done with

• G. Kane, P. Kumar, K. Bobkov, S. Watson  
(Non-thermal) 2006-2013

• S. Ellis, G. Kane, B. Nelson, M. Perry  
(DM is Hidden)

arXiv 1604.05320, PRL 117, 181802, 2016.  
1707.04530

• M. Fairbairn, E. Hardy, arXiv 1704.01804 <sup>JHEP</sup>  
(Hidden glueballs) <sub>1707</sub>

# Hidden Sectors

- Play a crucial role in this talk
- Huge Amount of literature on Hidden Sectors and DM
- This is not a review of Hidden sector models
- Rather, I will present a general picture emphasising those 2 points

## Theorists favourite: WIMPs

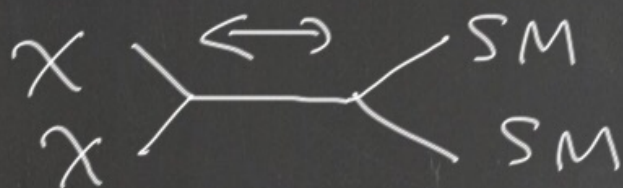
- Overwhelmingly WIMPs are favoured by theorists.
- Also targeted by many search strategies for DM.

Lets review WIMP DM  
(see Ibarra and Belanger talks)



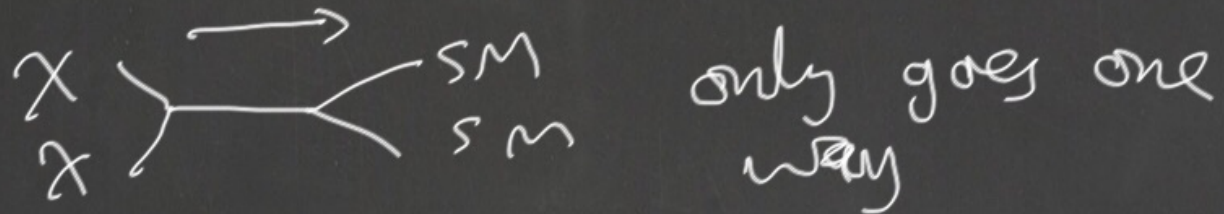
At the end of inflation (or whatever solves the horizon, flatness probs and seeds the CMB!):

- Assume Universe is radiation dominated with a high  $T \gg \underline{M_{EW} \sim 100 \text{ GeV}}$
- Standard Model particles are in equilibrium with WIMPS,  $\chi$



$\chi$  is a stable, electrically neutral particle charged under  $SU(2) \times U(1)_Y$ .

As Universe expands,  $T$  drops.  
 When  $T$  falls below  $m_x$ ,



And  $x$  particles freeze out with

$$\frac{H|_{T \sim \frac{m_x}{\text{few}}}}{\langle \sigma v \rangle_{xx \rightarrow \text{SM}}} \sim n_x$$

$$H \sim \frac{T^2}{m_{Pl}} \quad G \approx \frac{\alpha^2}{M_x^2} \quad \text{X}$$

so 
$$\mu_x \approx \frac{T^2 M_x^2}{\alpha^2 m_{Pl}} \quad \leftarrow$$

$$\rho/s \approx \frac{M_x^3}{g_* \alpha^2 T m_{Pl}} \quad \leftarrow (s \sim g_* T^3)$$

If  $T \sim \frac{M_x}{\alpha} \quad *$

$$\rho/s \sim 10^5 \frac{M_x^2}{m_{Pl}} \sim$$

$$\alpha \sim 10^{-2} \quad \sim 10^{-10} \text{ GeV} \quad \checkmark$$

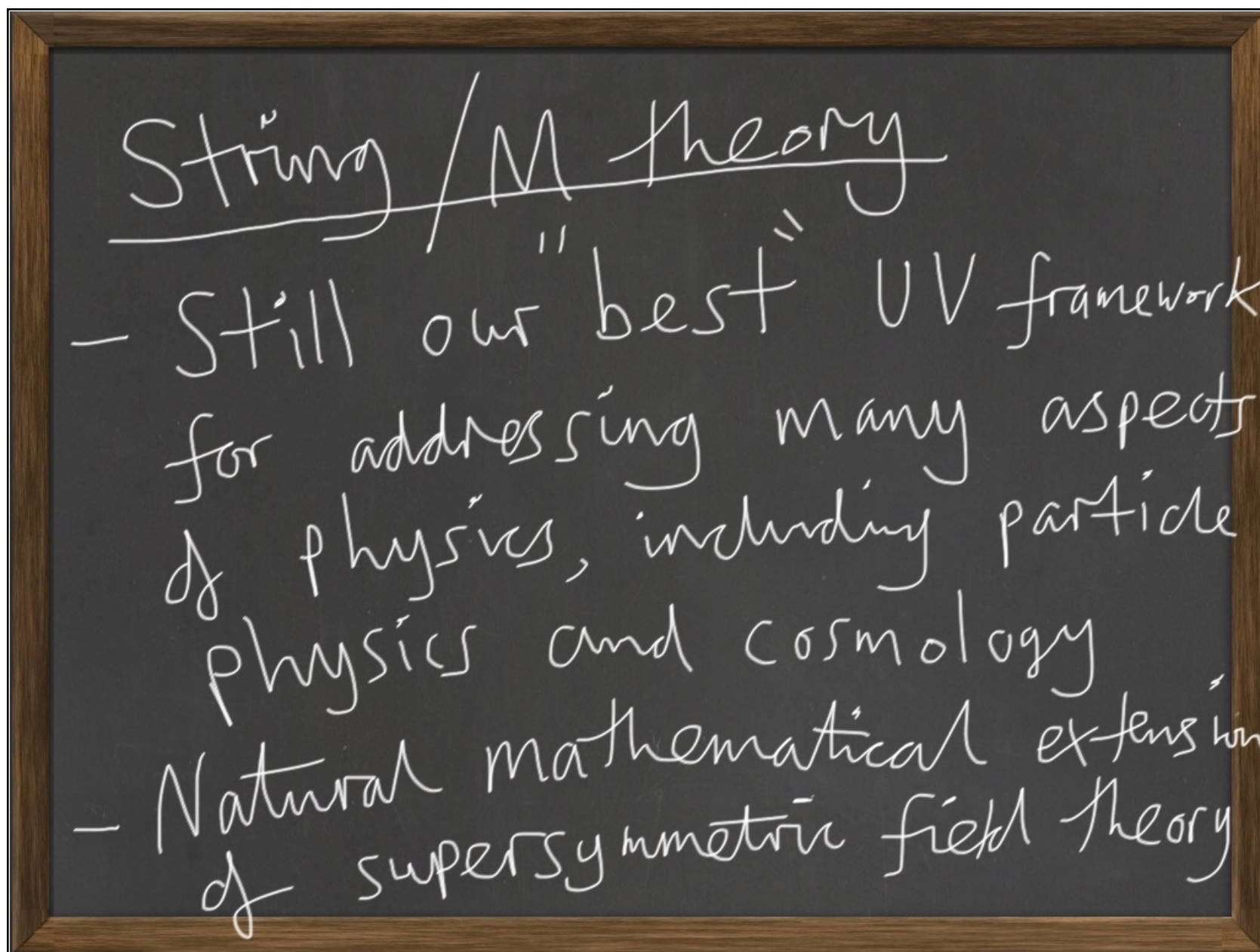


$$\frac{P}{S} \approx \frac{M_x^2}{g_* \alpha^2 m_{Pl}^2}$$

WIMP miracle or coincidence?

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- This is a nice coincidence between the weak scale, Planck scale and the DM abundance.
- But it doesn't have to be a miracle
- Also, it assumes ( $\rho \sim T^4$ ) radiation dominated early Universe. \*





We will consider the low energy limits of solutions of string/M-theory

$\exists$  many solutions of the form:

$$M^{9,1} = \underbrace{\mathbb{Z}^6}_{\text{compact, small}} \times \underbrace{M^{3,1}}_{\text{large}}$$

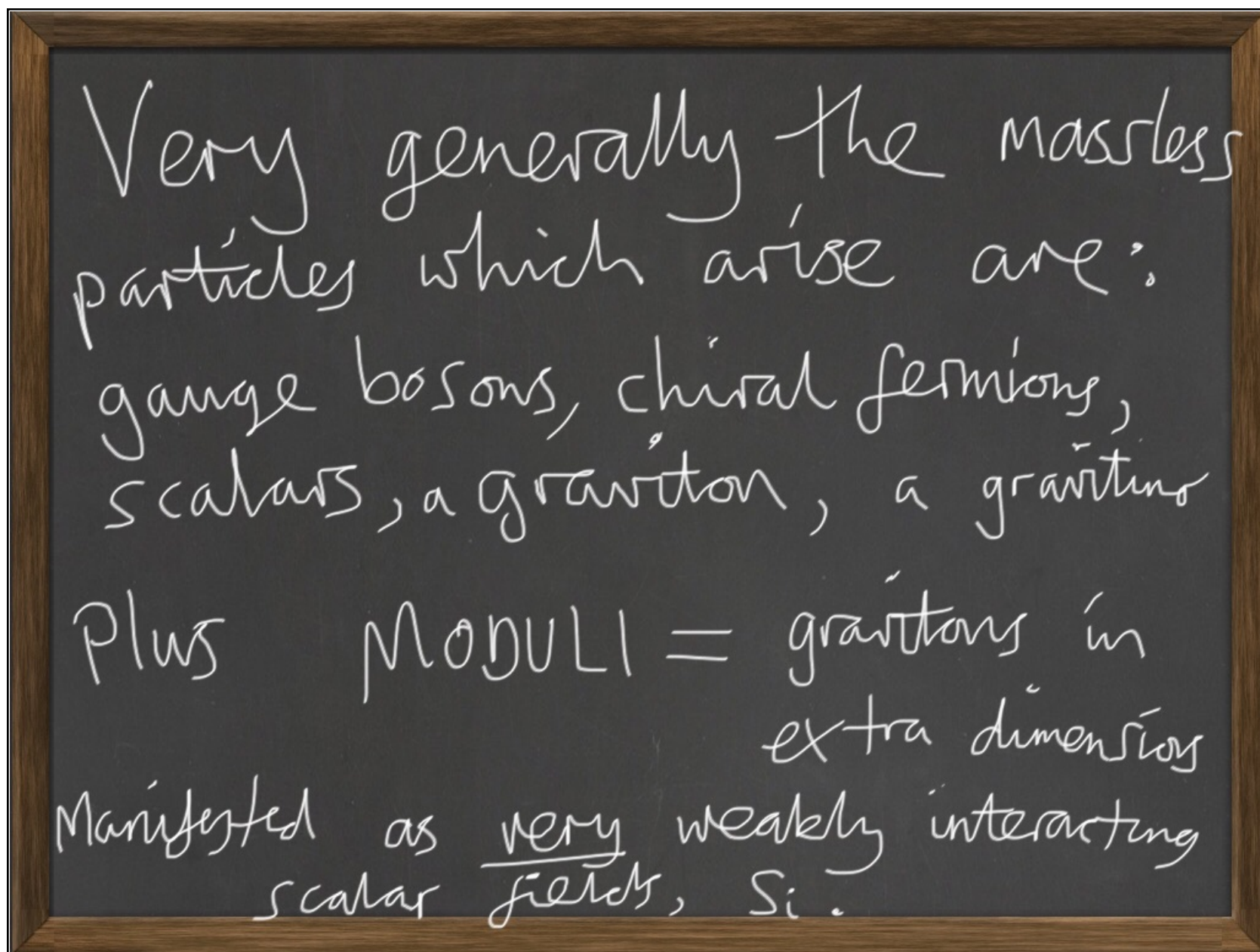
$$M^{10,1} = X^7 \times M^{3,1}$$

$$g(M^{10,1}) \cong g(X) + g(M^{3,1})$$

$$\text{Hol}(g(X)) = \zeta_2!$$

EXTRA

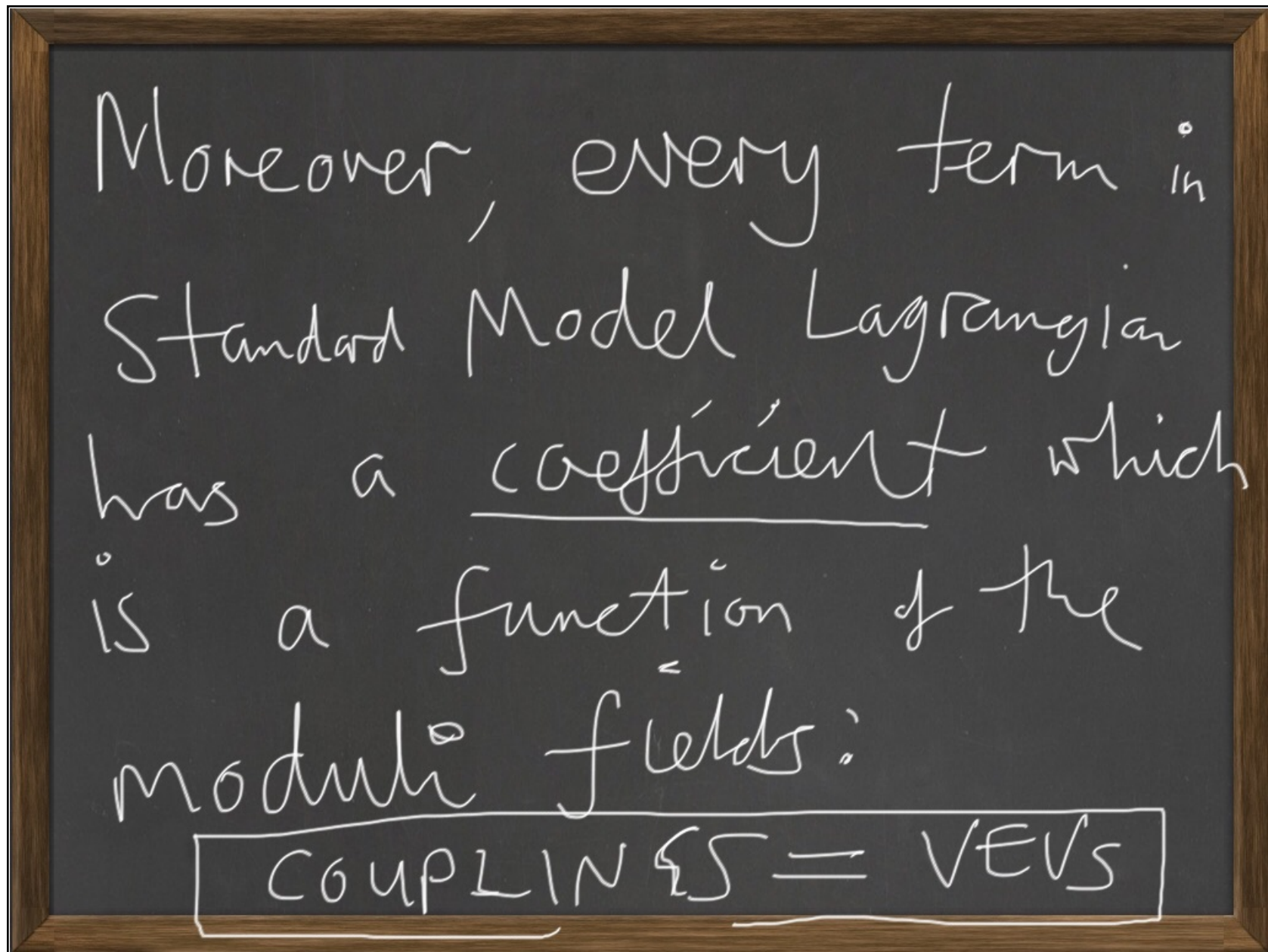
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Low energy,  $d=3+1$  Lagrangian is  
of the form, schematically,

$$\begin{aligned}
 -\mathcal{L}_{\text{matter} + \text{grains}} &= \frac{1}{16\pi G_N} \sqrt{-g_{3+1}} R_{3+1} + \frac{1}{g^2} \underline{F_{\mu\nu}^2} \\
 &\quad + i \bar{\Psi} \not{\partial} \Psi + \lambda H \bar{\Psi} \Psi \\
 &\quad + \underline{|\partial H|^2} - \underline{V(H, H^\dagger)} \\
 &\quad + \\
 \underline{\underline{\mathcal{L}_{\text{moduli}}}} &= \underline{\kappa^{ij}(s_i)} \left( \underline{\partial_\mu s_i} \underline{\partial^\mu s_j} + \kappa^{ij}(s) \underline{\partial_\mu a_i} \underline{\partial^\mu a_j} \right) \\
 &\quad - \underline{V(s_i, a_j)} \\
 s_i &= \text{moduli} \quad a_i = \text{axions} \quad + \dots
 \end{aligned}$$





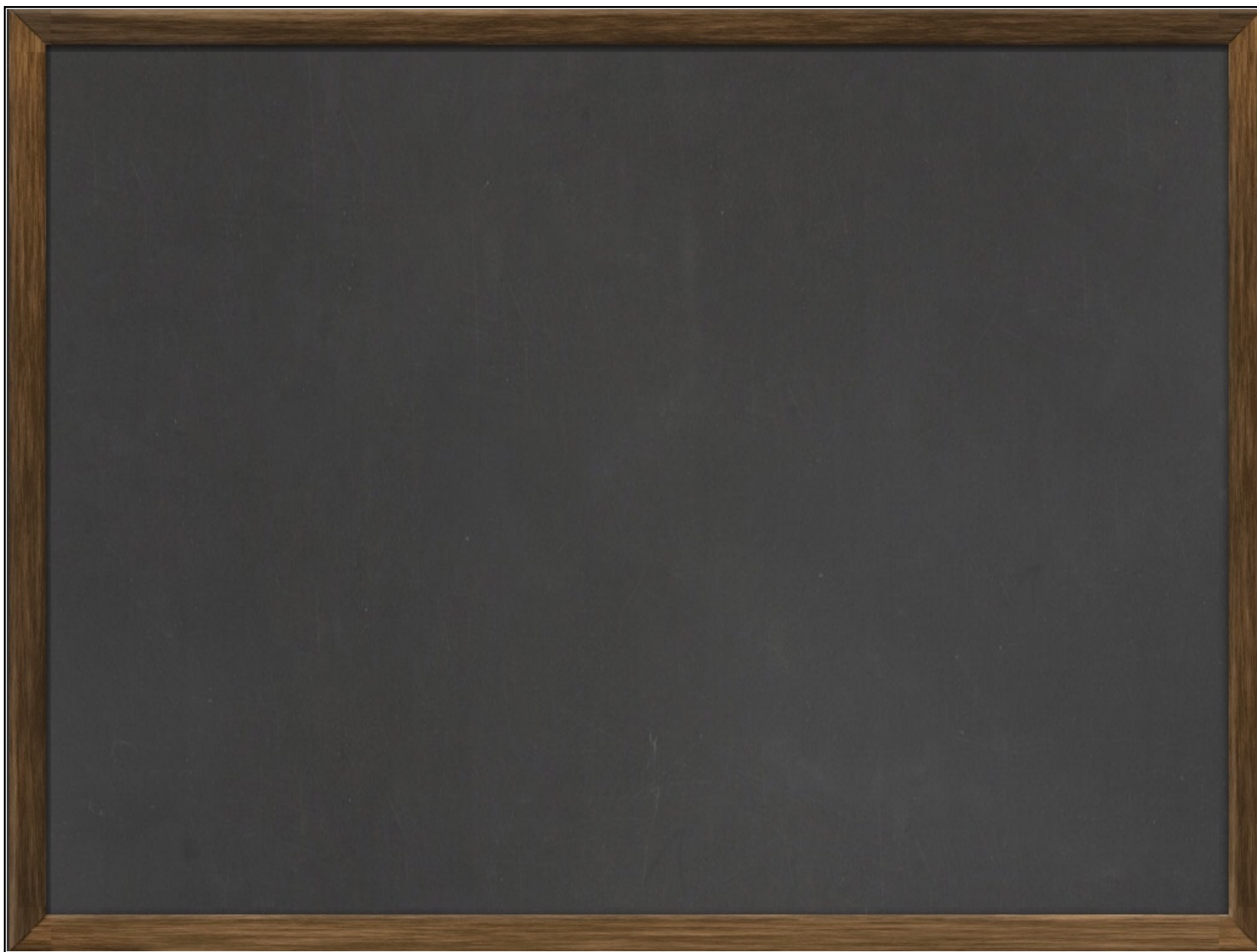
e.g.  $\frac{1}{g^2} F_{\mu\nu}^2$  is really

$S \propto \int d^4x \sqrt{g} F_{\mu\nu}^2$   $\frac{N(S) F_{\mu\nu}^2}{m^2}$ , a dim-5 operator.

Similarly  $\frac{\Lambda^4 \int d^4x \bar{\psi} \psi}{m^2}$  is really

$\rightarrow e^{-\frac{dis(S)}{m^2}} e^{\frac{idis(S)}{f}} \int d^4x \bar{\psi} \psi$

\* The moduli dependence of  $\Lambda$  varies from theory to theory.





- Moduli VEVs control couplings and masses
- Moduli have Planck suppressed couplings to ordinary matter
- Makes sense as Moduli are actually higher dimensional gravitons.

What about the moduli masses?

Supersymmetry:

$\Rightarrow$   $\mathcal{I}_{3\pi}$  is a supergravity theory.

In particular, there is only ONE  
MASS SCALE,  $M_{3/2}$ , the  
gravitino mass.

So  $M_{\text{moduli}} \sim M_{3/2}$ , without  
fine-tuning

# - Moduli in Cosmology

- Polonyi Problem  
(Coughlan et al 83)

- Cosmological Moduli Problem  
(CMP)

- Carlos, Casas, Quevedo, Roulet 93

Banks, Kaplan, Nelson 93



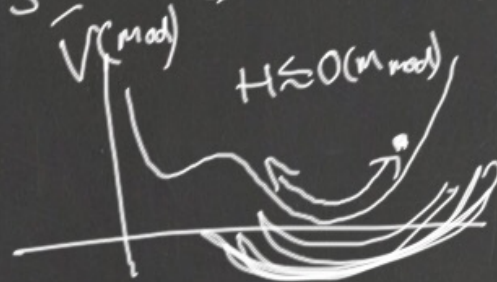
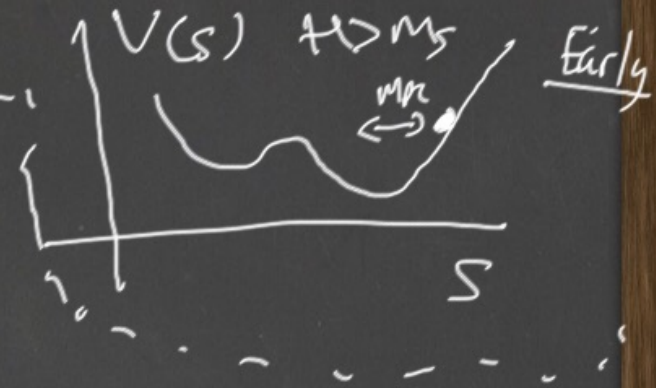
# Cosmology of This Theory

At the end of inflation (or whatever.....)

If  $H \gg m_{3/2} \sim m_{\text{moduli}}$ , the moduli will be stuck at some  $O(1) m_{\text{Pl}}$  place in its potential.

Later,  $H \sim O(m_{\text{moduli}})$

and  $s$  oscillates:



$\rho_{\text{moduli}}$  is a MATTER component  
which QUICKLY dominates over  
 $\rho_{\text{radiation}}$  ( $\frac{1}{R^3}$  vs  $\frac{1}{R^4}$ ).

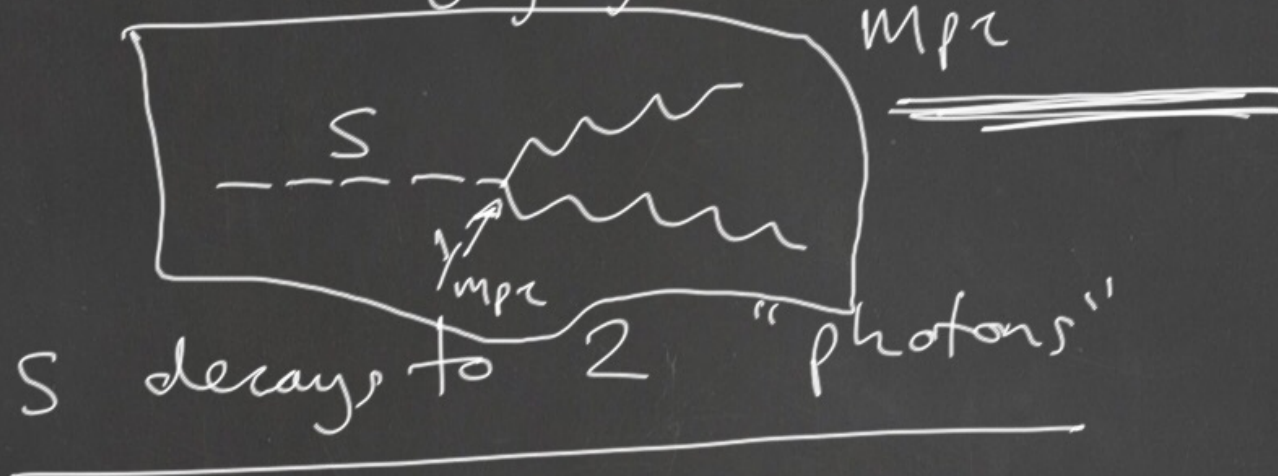
Hence, the Universe becomes  
matter dominated by the  
moduli fields.

GENERIC PREDICTION OF STRINGS/M  
THEORY MODELS WITH LOW ENERGY  
SUSY.



The moduli are unstable particles.  
 (They couple to matter particles fairly  
 'generically' and 'uniformly'.)

(Consider, e.g.,  $\mathcal{L}_{\text{gauge}} \sim \frac{S}{M_{\text{Pl}}^2} F_{\mu\nu}^2$ )





Decay width (or probability) is

$$\Gamma(s \rightarrow \gamma\gamma) \propto |M|^2$$

$$M \sim \frac{1}{m_{Pl}} \cdot$$

$$\Gamma \sim O\left(\frac{1}{m_{Pl}^2}\right) \sim G_N$$

$$\therefore \Gamma(s \rightarrow \gamma\gamma) \approx \frac{M_{moduli}^3}{2 m_{Pl}^2} \times$$

$$\therefore \text{Lifetime } \tau_{moduli} \approx \frac{m_{Pl}^2}{M_{moduli}^3} \leftarrow$$

So, after dominating  $\rho_{\text{universe}}$ , the  
moduli will decay after a time

$$t_{\text{decay}} \sim \frac{M_{\text{Pl}}^2}{m_{\text{moduli}}^3}$$

equivalently  $H_{\text{decay}} \sim \frac{m_{\text{moduli}}^3}{M_{\text{Pl}}^2}$

$$\frac{1}{H_{\text{decay}}} \sim O(1) \text{ sec} \left( \frac{\text{TeV}}{m_{\text{moduli}}} \right)^3$$

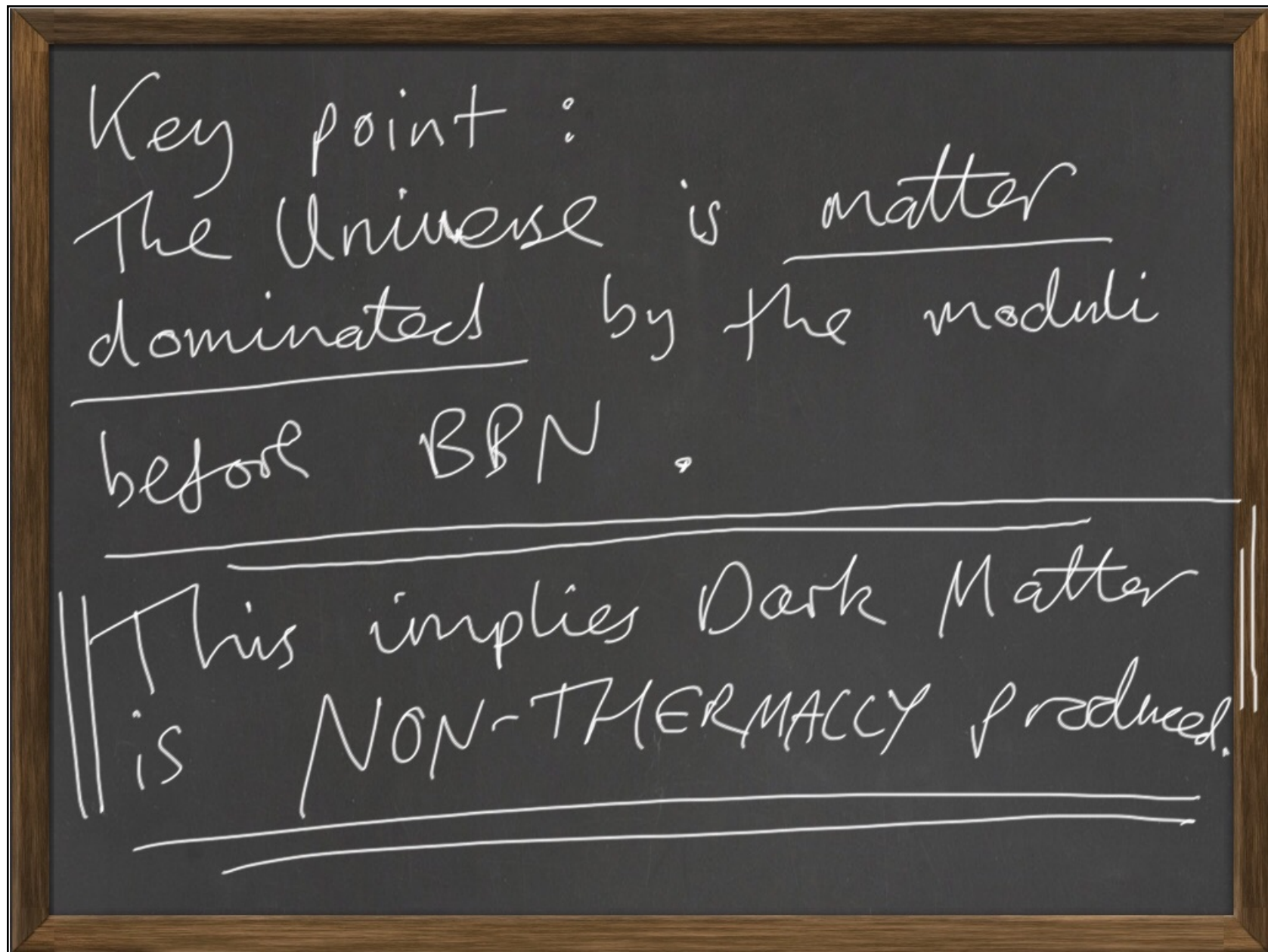
This is in the middle of  
BBN!

So for  $M_{3,2} \sim \text{TeV}$ , moduli decay during BBN. This is bad as they decay into quarks, leptons and gauge bosons.

This injects charged particles and hadrons into the plasma which can dis-associate nuclei and drastically change the successful predictions of BBN.



But, for  $M_{3/2} \sim \mathcal{O}(10) \text{ TeV}$ ,  
the moduli decay before BBN,  
create a radiation dominated  
universe with  $T \sim \underline{10 \text{ MeV}}$   
and this is consistent.

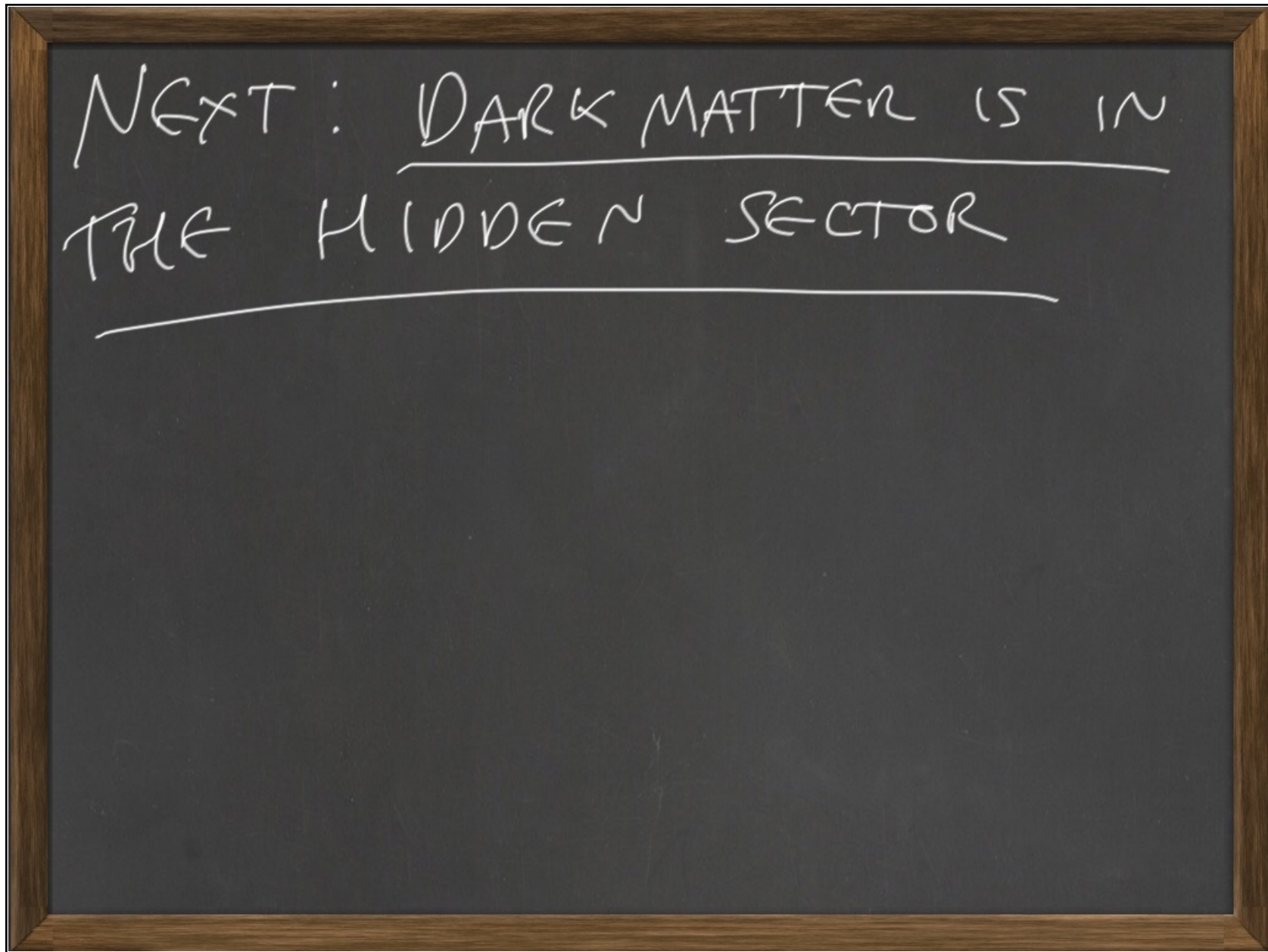


This seems quite a generic conclusion.

### Caveats

- Could assume  $H_{\text{inf}} \ll M_{3/2}$   
(not typical)
- Could arrange a late period of inflation to "get rid of the moduli". (Seems 'tuned?')





## Hidden Sectors

Def<sup>n</sup>: A particle is in the  
 Hidden sector if it has no  
tree level gauge interactions  
 with the Standard Model.  
 ie it has no  $SU(3) \times SU(2) \times U(1)_Y$   
 charge at tree level.

Since we have no idea why the Standard Model has  $G = SU(3) \times SU(2) \times U(1)$  and 45 fermions and a Higgs doublet, there is no reason NOT to consider additional gauge sectors and matter. This is exactly the picture that emerges from string/M theory



## Hidden Sectors in String/M theory

- In Heterotic  $E_8 \times E_8$  theory, one  $E_8$  is "hidden" w.r.t the other.
- In Type II theories, D-branes can be physically separated in the extra dimensions.
- In M/F-theory, singularities supporting gauge symmetries are physically separated. ( $G_2$  or  $C_4$ )

There is no privilege given to the Standard Model.

Generically expect additional gauge groups and matter.

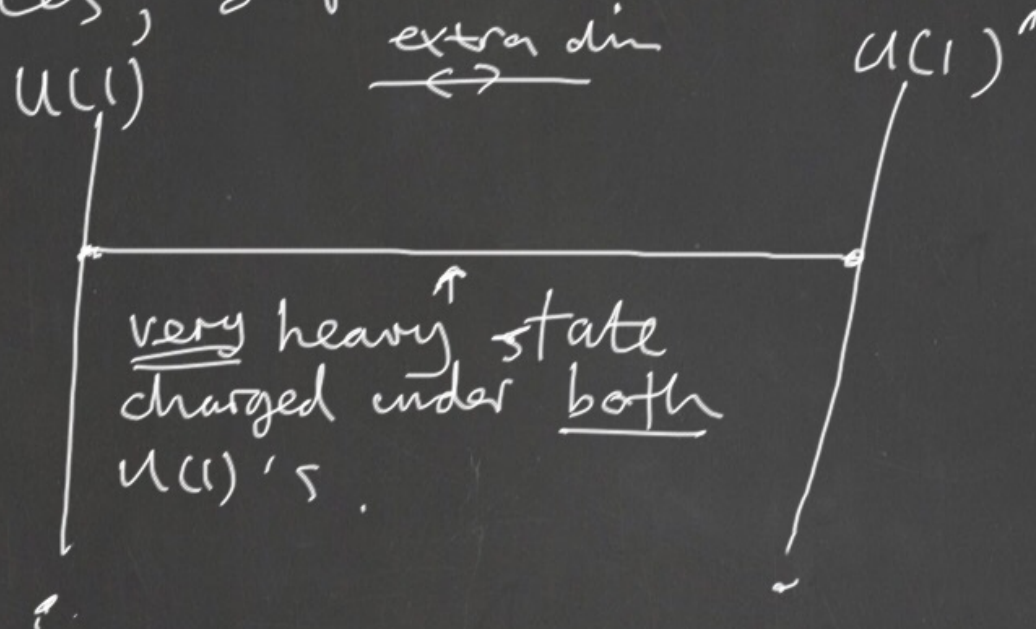
HIDDEN SECTOR MATTER  
IS GENERIC

(Talk yesterday by S. Ellis)

Consider a Type II string model with

$$G = U(1) \times U(1)''$$

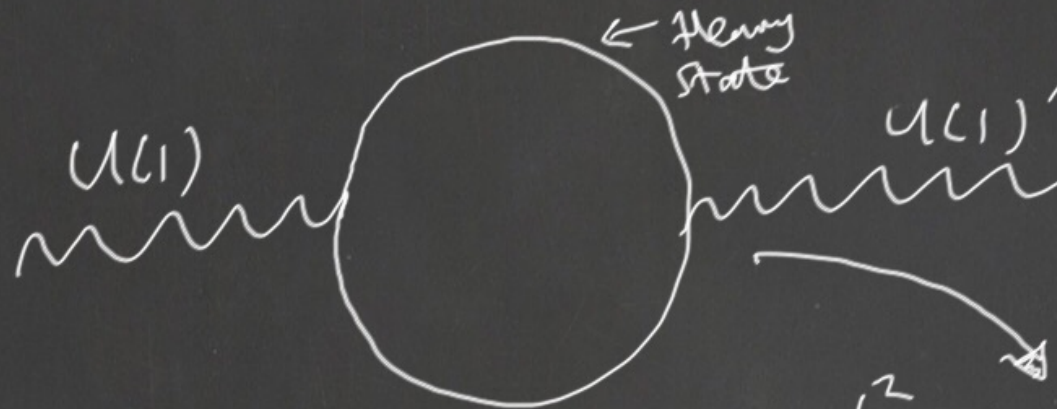
Realise this with two stacks of  
D-branes, separated in extra dim:





Mass, heavy state  $\sim \frac{M_{str}^2 R_{KK}}{M_{KK}}$

H induces a renormalisation of kinetic terms;



$$i \quad F_{\mu\nu}^2 + \tilde{F}_{\mu\nu}^2 \rightarrow F_{\mu\nu}^2 + F_{\mu\nu}'^2 + \epsilon F_{\mu\nu} F_{\mu\nu}'$$

Since  $FF'$  is dim 4,  $\epsilon$  is only log sensitive to UV

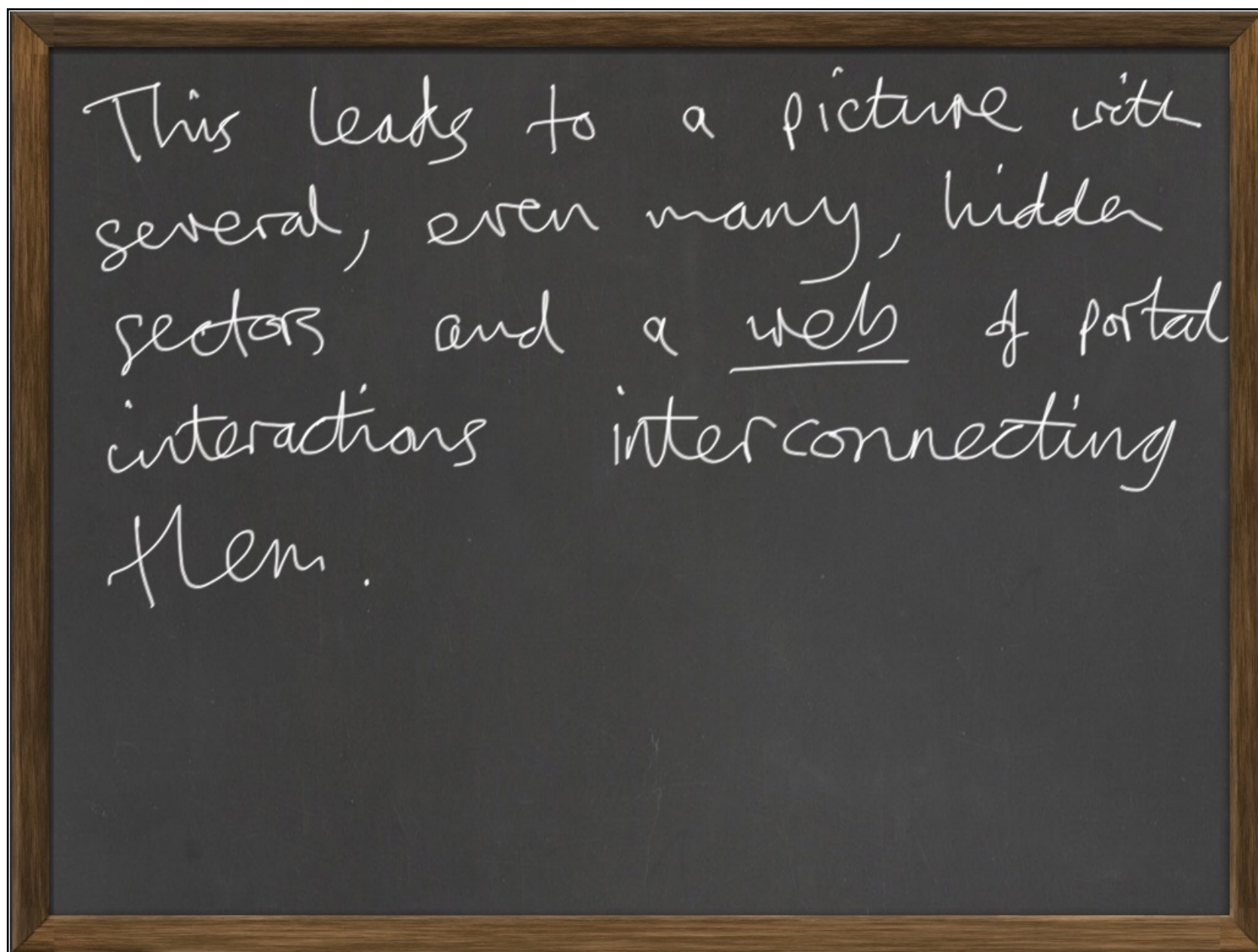
$$\epsilon \sim \frac{gg'}{12\pi^2} \ln\left(\frac{\Lambda}{M}\right).$$

- Such mixings are generically present between U(1)'s.
- This has been known for quite some time (Dienes, Kolda, March-Russell '97)

The  $\mathbb{E}FF'$  interactions (and those related to it by supersymmetry) provides a PORTAL between different hidden sectors.

eg gauge bosons can mix between sectors, as can gauginos, via  $\mathbb{E} \lambda \not{D} \lambda'$ .





Consider now the (supersymmetric) Standard Model sector. This has a (so-called) "Lightest Supersymmetric Particle" which is often the WIMP DM candidate.

(Usually (without hidden sectors) this is stable as it is the lightest particle with non zero R-parity.

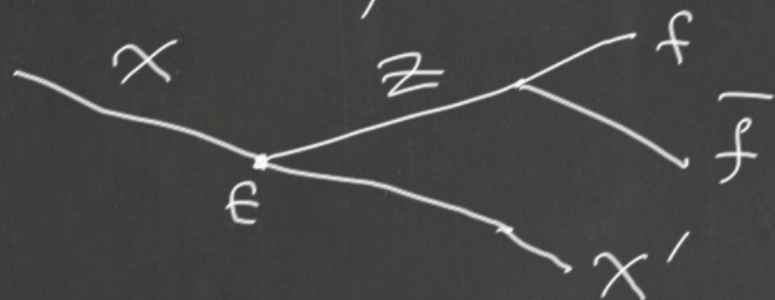
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With multiple hidden sectors, there is NO GOOD REASON why the LVSP\* should be the lightest R-parity charged particle in the theory. It could happen by accident, but is unlikely.

\*LVSP = Lightest Visible Sector Supersymmetric Particle

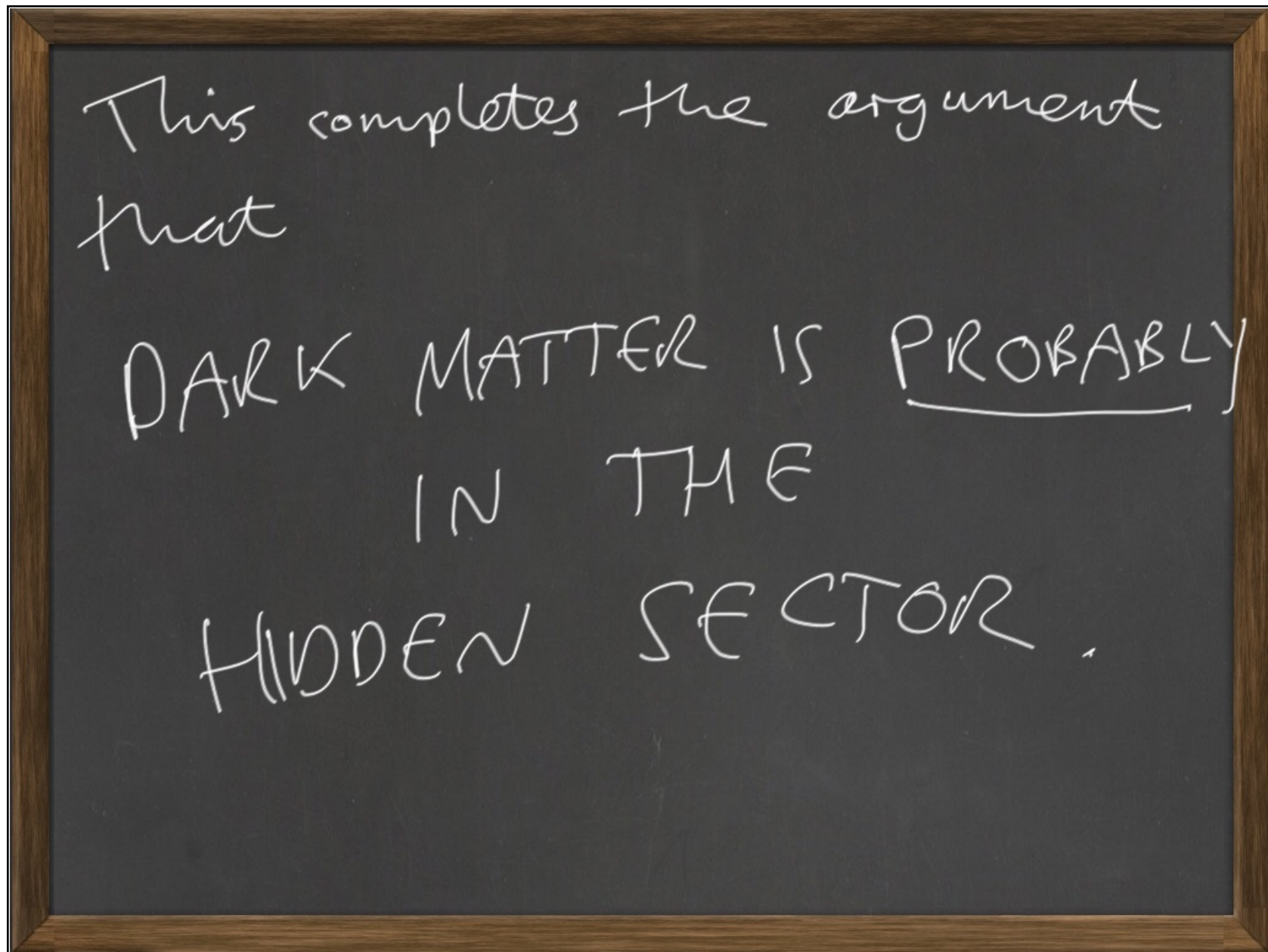


Mixing between Hidden  $U(1)'$  and  $U(1)_Y$  leads to, e.g.



and  $\tau_x \sim 10^{-17} \text{ s} \left( \frac{10^{-3}}{\epsilon} \right)^2 \times \text{mixing angles}$   
for on shell  $Z$

$\tau_x \sim 10^{-9} \text{ s} \left( \frac{10^{-3}}{\epsilon} \right)^2 \left( \frac{50 \text{ GeV}}{m_x - m_{x'}} \right)^4 \times \text{angles}$   
for 3-body decay



1  
The argument relied on  
three ingredients :

- 1 : Hidden Sectors are Generic
- 2 : PORTALS are generic
- 3 : The LUSP is not the lightest super particle.



So, what is Dark Matter?

- Axions are also generic in string/M theory and are very difficult to remove.
- Stable particles produced by moduli decays will also be a component of Dark Matter:
- Light, decoupled (chiral) fermions;
- Hidden sector glueballs, other composites

Conclude:

- Dark Sector is very rich, perhaps richer than the Standard Model.
- DM is more likely to be non-thermally populated
- Rich phenomenology at LHC and elsewhere.

