

Jet substructure shedding light on Heavy Neutrino at the LHC

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Neutrino puzzle

* SM neutrino is massless

- ♦ No right handed counter part for Dirac term
- ♦ Lepton no violation/ Gauge invariance with both neutrino, Higgs from doublet => No Majorana term
- ♦ Accidental $B-L$ symmetry prevent mass term in all order

* But, (already) confirmed BSM signatures from neutrino oscillation => Having tiny but non-zero mass and mixing

* Extended Standard Model

Seesaw

► Simple choice - Seesaw mechanism : Dim 5 operator

- ♦ Adding right handed singlet Majorana neutrino

$\langle v \rangle$

$\mathcal{L} \supset -Y_D^{\alpha\beta} \bar{\ell}_L^\alpha H N_R^\beta - \frac{1}{2} M_N^{\alpha\beta} \bar{N}_R^{\alpha C} N_R^\beta + H.c.$

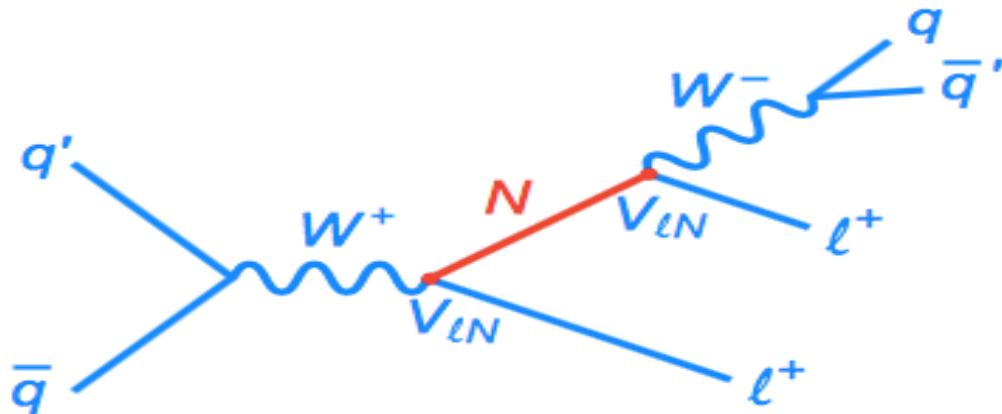
$M_\nu = \begin{pmatrix} 0 & M_D \\ M_D^T & M_N \end{pmatrix} \quad M_D = \frac{Y_D v}{\sqrt{2}}$

$$m_\nu \simeq -M_D M_N^{-1} M_D^T$$

- ♦ Small Yukawa/heavy-light mixing ($M_D M_N^{-1}$) \Rightarrow sub-TeV 'heavy neutrino'
- ♦ Cases-Ibarra parameterisation Yukawa coupling expressed in terms of a orthogonal matrix which remains completely arbitrary and hence can be large.

Seesaw at LHC

- * If heavy neutrino is 'light' enough
& heavy-light mixing 'large' enough
- * Resonant production of heavy Majorana neutrino

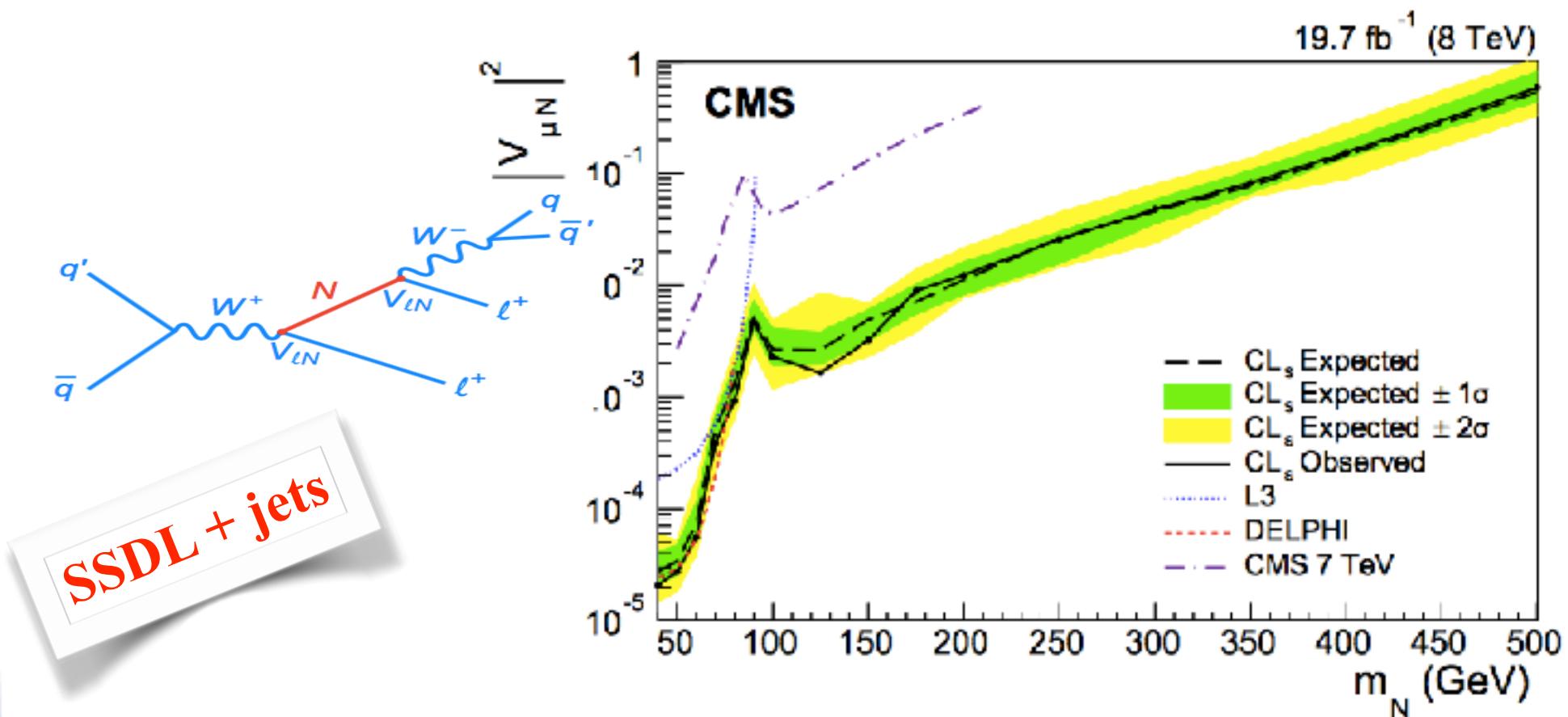


- ◆ Produces multi-lepton signature at the LHC

$$\begin{aligned} pp &\rightarrow \ell_1^+ N, \quad N \rightarrow \ell_2^+ W^- \\ pp &\rightarrow \ell_1^- \bar{N}, \quad \bar{N} \rightarrow \ell_2^- W^+ \end{aligned}$$

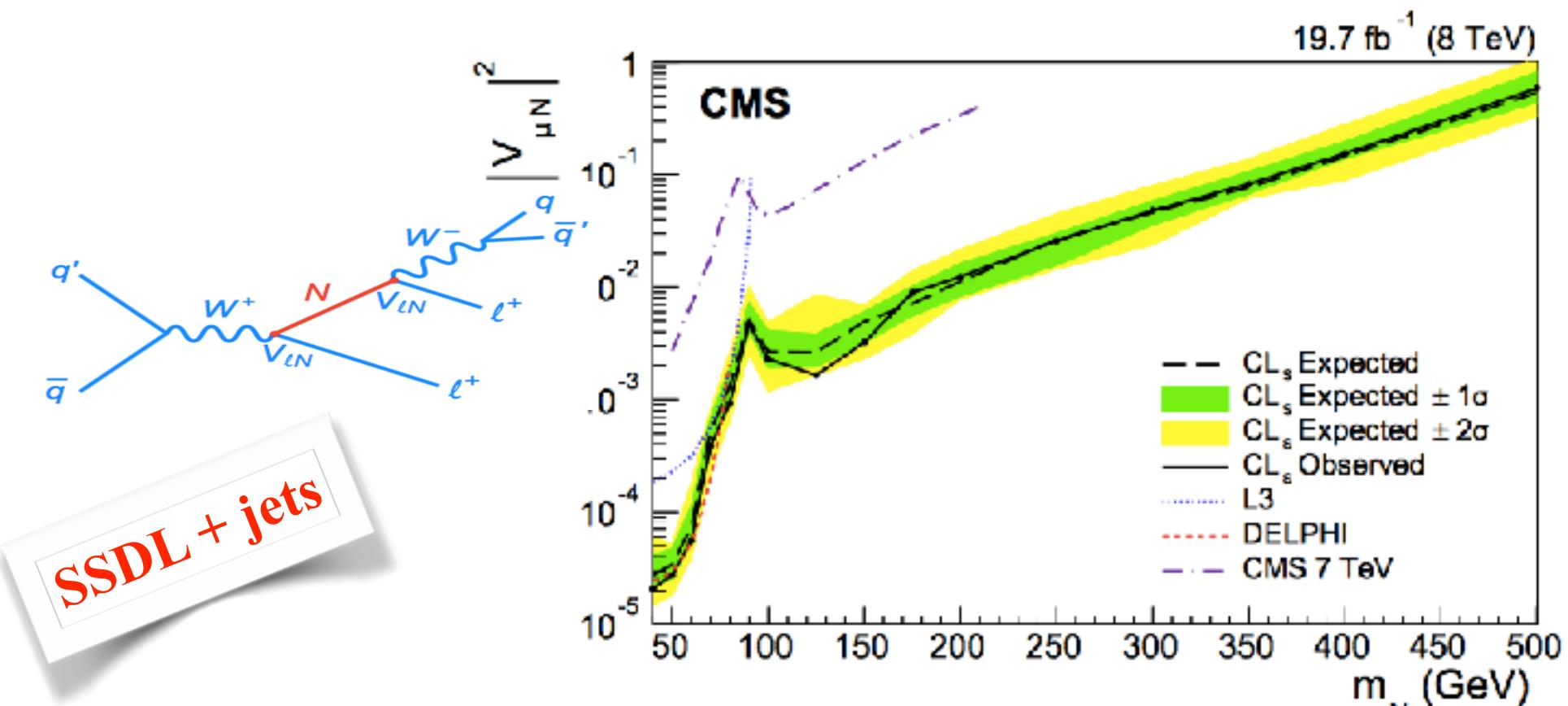
SSDL + jets

Seesaw at LHC



CMS8: 1501.05566, 1603.02248
ATLAS8: 1506.06020

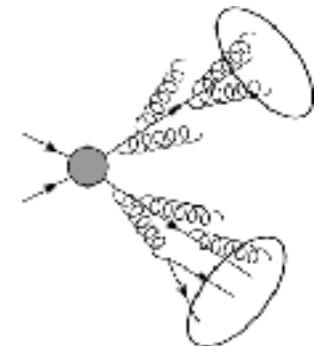
Seesaw at LHC



If we could use the internal feature of these additional jets, to deduce:

- ‘W’-jets
- Production topology
- Control background

"New" jets @ LHC

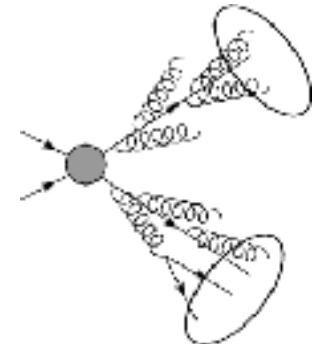


- Increased energy of hadron collider
- Exploration for intermediate to very heavy resonance
 - : if $PT < 2\text{TeV} \sim 100\text{ GeV}$ splitting is distinguishable with cell size 0.1
- Calorimeters have enough finer resolution & improving
 - : $\sim 50 E_{\text{cell}}$ of .1 in $R=0.4$
- Computations in fast sequential algorithms
 - : fast jet algorithm

⇒ production and analysis of boosted objects & fat jet

Decay & fragmentation of boosted object Produce a collimated spray of hadron, standard Algorithm consider it as single jet

“New” jets @ LHC

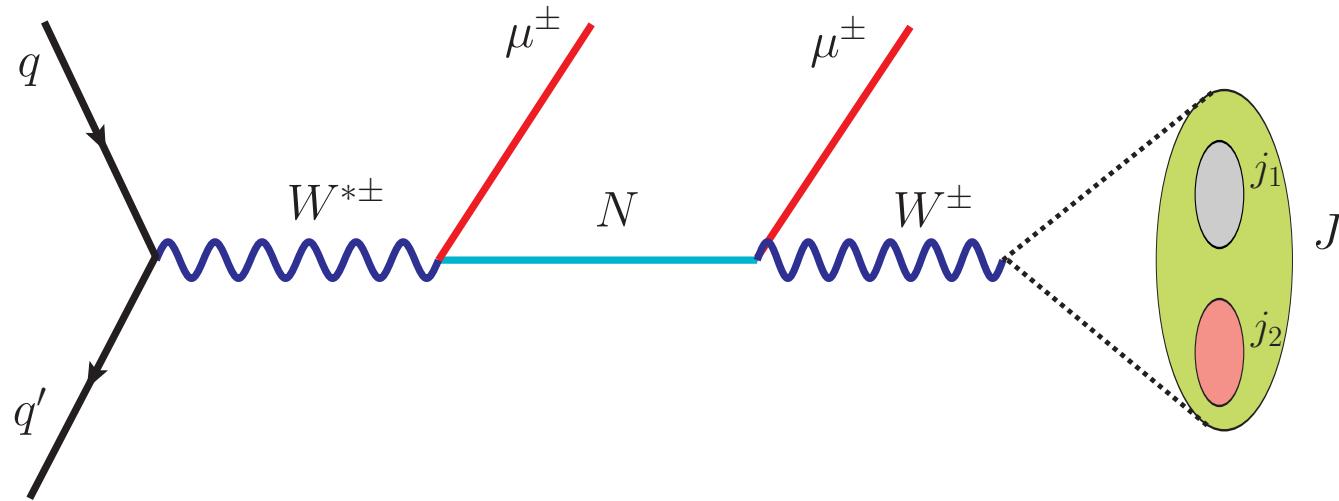


- With same jet mass, boosted hadronic objects have a fundamentally different energy pattern VS to QCD jets
- Clean up smearing effect of jet contamination from IRS, UE, Pileups
- In this work, we exploit the characteristics and kinematic properties for these Boosted W-jets

⇒ production and analysis of boosted objects & fat jet

Decay & fragmentation of boosted object Produce a collimated spray of hadron, standard Algorithm consider it as single jet

Fat jet from Heavy Neutrino



$$pp \rightarrow \ell_1^+ N, \quad N \rightarrow \ell_2^+ W^-$$

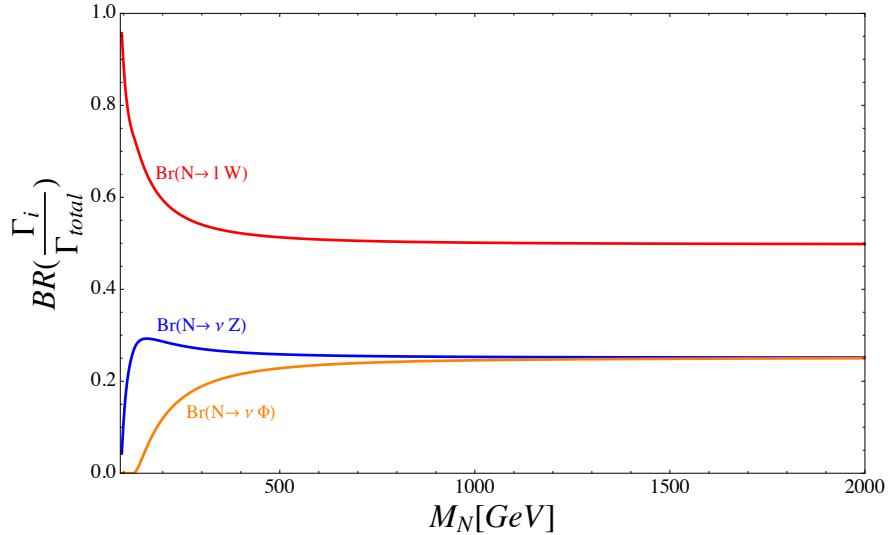
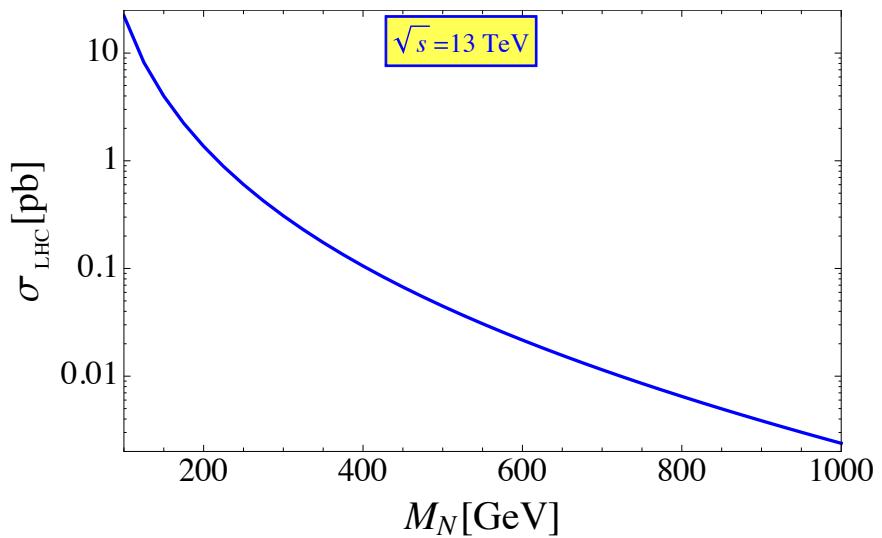
$$pp \rightarrow \ell_1^- \bar{N}, \quad \bar{N} \rightarrow \ell_2^- W^+$$

$$W^- \rightarrow J$$

$$W^+ \rightarrow J$$

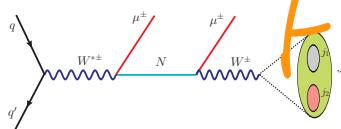
Features of these boosted fat-jets are prominent for large M_N
 ~ 300 GeV and above

Heavy Neutrino at LHC



Total production cross-section and branching ratios of heavy Majorana neutrino as a function of its mass at the LHC with $\sqrt{s} = 13 \text{ TeV}$ and normalised by the $|V_{\mu N}|^2$.

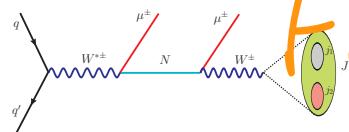
$$BR(N \rightarrow \ell W) : BR(N \rightarrow \nu Z) : BR(N \rightarrow \nu H) \simeq 2 : 1 : 1.$$



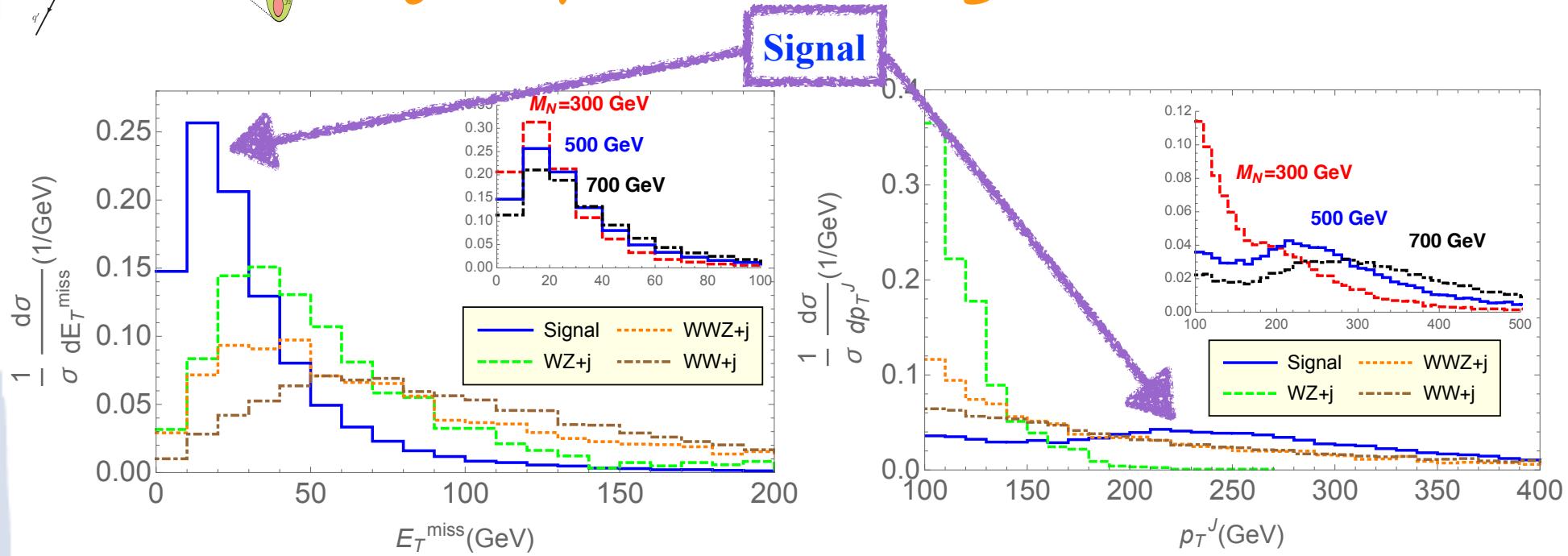
Fat jet from Heavy Neutrino

- ✿ A simplified model - Only light-heavy mixing for the Muons.
- ✿ Backgrounds for Same-sign-di-muon + fat-jet :
 - Electroweak gauge boson decay (Di-boson+j, Tri-boson+j)
 - + “fat-jet” - from a W boson OR mimicked from QCD jets
- ✿ MadGraph5-aMC@NLO => Pythia => Delphes
- ✿ MLM matching on a shower- k_T algorithm with p_T -ordered showers

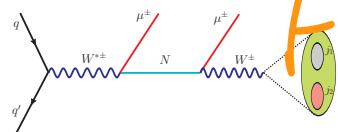
- Same sign μ^\pm with $p_T^\mu > 10 \text{ GeV}$, $|\eta^\mu| < 2.4$
- Candidate fat-jets are identified - $\mathcal{R} = 0.8$ CA jet with $|\eta^j| < 2.4$
- Hardest fat-jet $\sim W^\pm$ candidate jet (J) having $p_T^J > 100 \text{ GeV}$



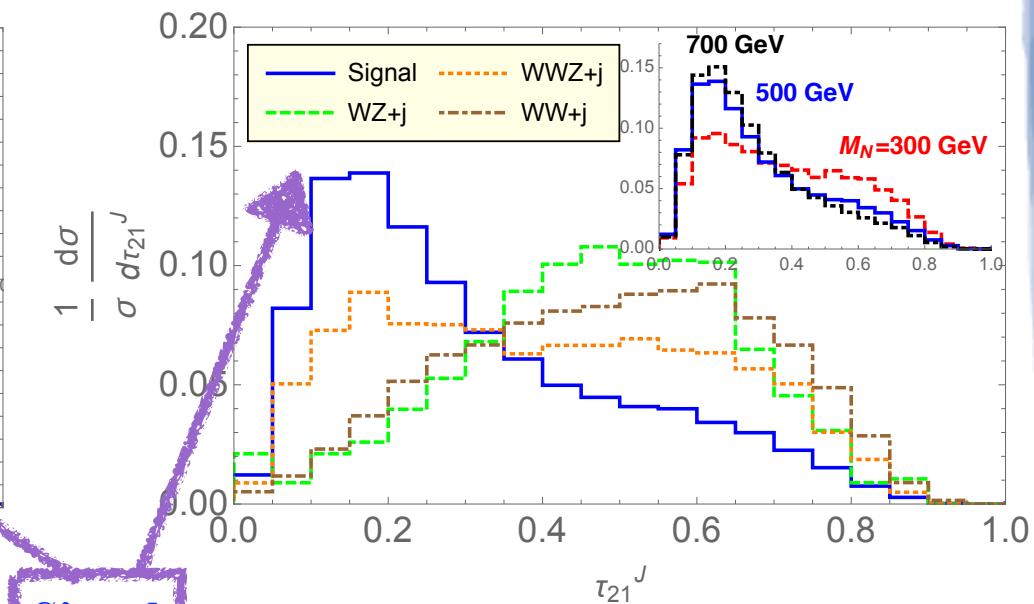
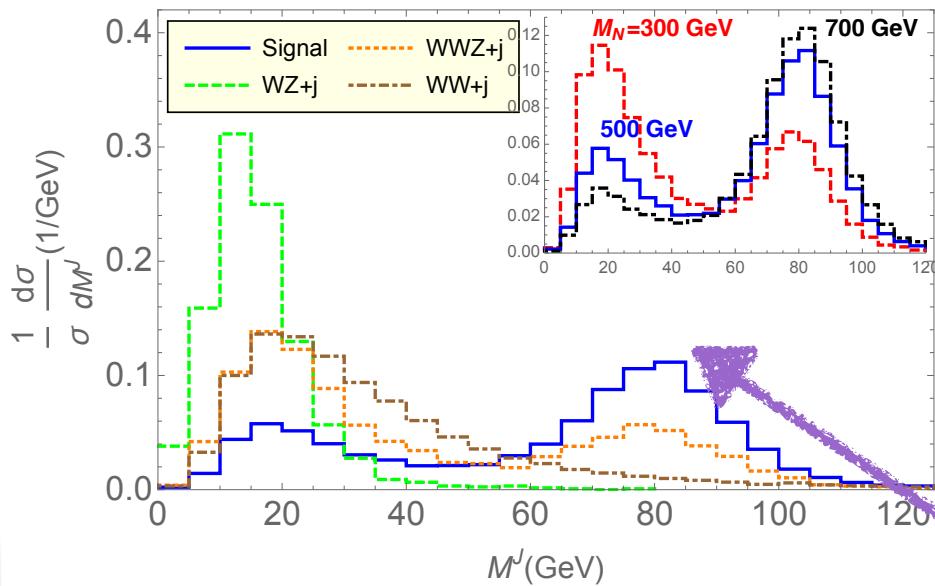
Fat jet from Heavy Neutrino



Normalised distribution as a function of MET & p_T^J after the application of the basic selection cuts; including $p_T^J > 100 \text{ GeV}$



Fat jet from Heavy Neutrino

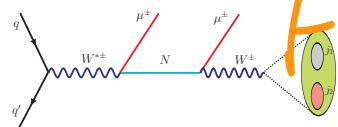


Signal

Normalised distribution as a function of **invariant-mass M_J** & **N -subjet ratio τ_2/τ_1** of fat jet after the application of the basic selection cuts; including $p_T^J > 100$ GeV

$$\tau_N = \frac{1}{N_0} \sum_i p_{i,T} \min \{ \Delta R_{i1}, \Delta R_{i2}, \dots, \Delta R_{iN} \} \quad N_0 = \sum_i p_{i,T} R_i$$

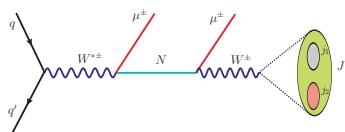
– Inclusive jet shape variable quantify if the original jet consists of N daughter subjets



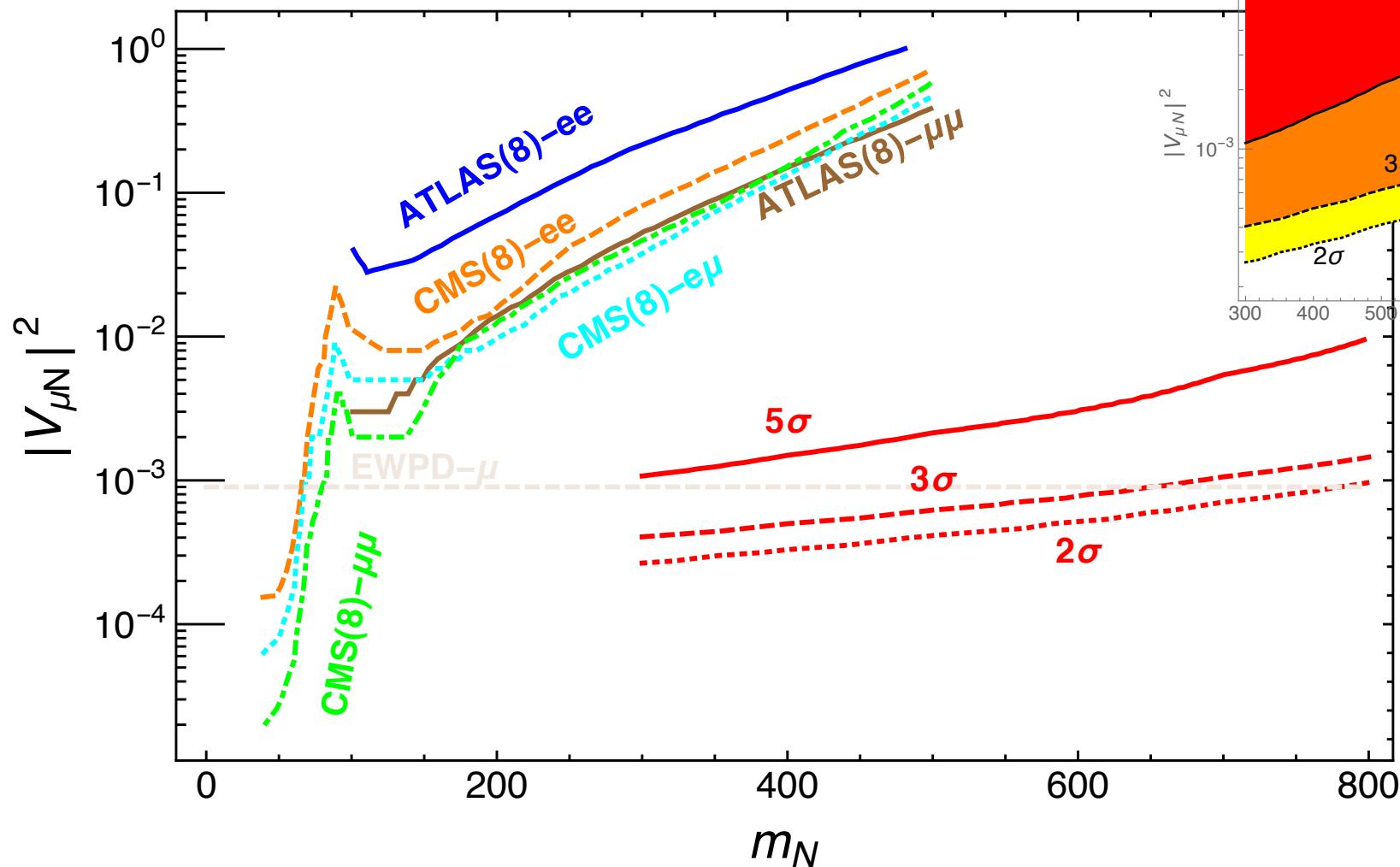
Fat jet from Heavy Neutrino

Cut	Signal for M_N			Background		
	300 GeV	500 GeV	700 GeV	$WW+j$	$WZ+j$	$WWZ+j$
Pre-selection + $\mu^{\pm}\mu^{\pm} + J$ $p_T^J > 100$ GeV	$60.0 + 33.0$ [100%]	$26.7 + 17.4$ [100%]	$14.0 + 9.5$ [100%]	$3282.0 + 3136.5$ [100%]	$5474.4 + 4234.5$ [100%]	$126.05 + 120.2$ [100%]
$p_T(\mu_{1,2})$, $m_{\mu\mu}$	$58.0 + 29.0$ [94%]	$24.1 + 14.8$ [88%]	$11.4 + 6.7$ [77%]	$2724.30 + 2575.03$ [83%]	$3045.2 + 2796.2$ [60%]	$104.0 + 96.7$ [82%]
$E_T^{\text{miss}} < 35$ GeV	$48.84 + 20.0$ [74%]	$20.8 + 13.2$ [77%]	$7.3 + 5.6$ [55%]	$314.04 + 197.1$ [7.9%]	$105.2 + 104.2$ [2.2%]	$12.1 + 9.8$ [8.9%]
$p_T^J > 150$ GeV	$25.6 + 15.0$ [44%]	$15.2 + 10.5$ [58%]	$6.0 + 4.3$ [44%]	$181.02 + 110.4$ [4.5%]	$20.2 + 15.08$ [0.4%]	$7.07 + 6.2$ [5.3%]
$M_J > 50$ GeV	$21.4 + 12.3$ [36%]	$11.2 + 7.4$ [42%]	$4.8 + 3.2$ [34%]	$41.05 + 32.08$ [1.1%]	$6.4 + 4.7$ [0.1%]	$3.3 + 2.5$ [2.3%]
$\tau_{21}^J < 0.5$	$19.5 + 10.0$ [32%]	$9.6 + 5.2$ [34%]	$3.9 + 2.0$ [25%]	$21.1 + 19.3$ [0.6%]	$3.3 + 2.9$ [0.06%]	$1.5 + 1.4$ [1.2%]

The effectiveness of different variables in minimizing backgrounds is illustrated in the form a cut flow. The two numbers correspond to expected events in $\mu^+\mu^+$ and $\mu^-\mu^-$ channels. We adopt a typical mixing angle $|V_{\mu N}| = 0.03$. The numbers are for an integrated luminosity of 3000fb^{-1} , at the 13 TeV LHC.



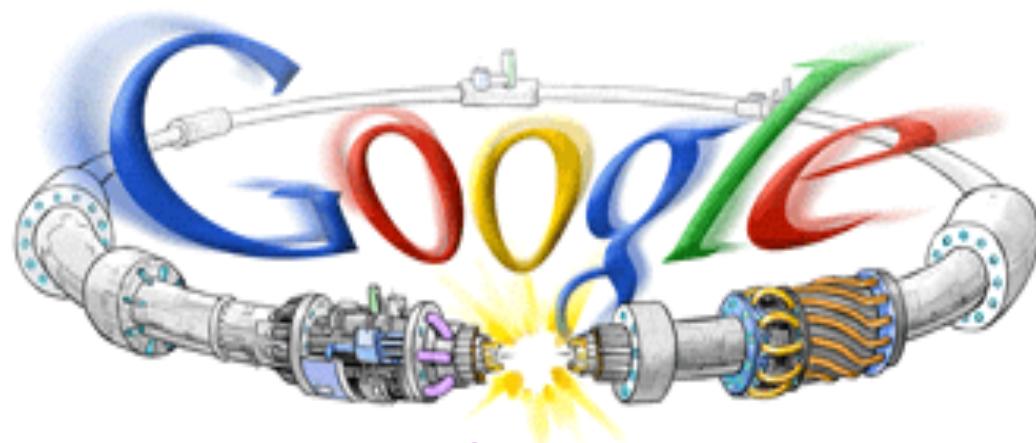
Heavy Neutrino at LHC



Exclusion limit in terms of heavy neutrino mass M_N and $|V_{\mu N}|^2$ at the 13 TeV LHC with other available limits.

Outlook

- Non-zero but Tiny neutrino mass can be realised through the simple type-I mechanisms which involves the Majorana type heavy neutrino.
- TeV scale heavy neutrinos can be produced at the LHC through a large mixing angle with the SM Light neutrinos
- Such production channel is explored using boosted fat jet substructure which is inevitable for probing large heavy neutrino mass.
- Opens up interesting possibility to explore intermediate to heavy masses of RHN



Thank You