

# *Jet substructure shedding light on Heavy Neutrino at the LHC*

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Partha Konar

Physical Research Laboratory

Ahmedabad, India

Collaborators:

Arindam Das,  
Arun Thalapillil,  
Akanksha Bhardwaj  
ArXiv:1709.09712



# Neutrino puzzle

## \* SM neutrino is massless

- ♦ No right handed counter part for Dirac term
- ♦ Lepton no violation/ Gauge invariance with both neutrino, Higgs from doublet => No Majorana term
- ♦ Accidental  $B-L$  symmetry prevent mass term in all order

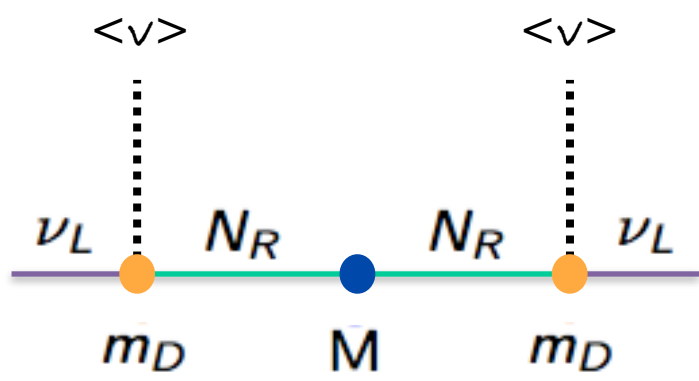
\* But, (already) confirmed BSM signatures from neutrino oscillation => Having tiny but non-zero mass and mixing

## \* Extended Standard Model

# Seesaw

► Simple choice - Seesaw mechanism : Dim 5 operator

♦ Adding right handed singlet Majorana neutrino



$$\mathcal{L} \supset -Y_D^{\alpha\beta} \bar{\ell}_L^\alpha H N_R^\beta - \frac{1}{2} M_N^{\alpha\beta} \bar{N}_R^{\alpha C} N_R^\beta + H.c.$$

$$M_\nu = \begin{pmatrix} 0 & M_D \\ M_D^T & M_N \end{pmatrix} \quad M_D = \frac{Y_D v}{\sqrt{2}}$$

$$m_\nu \simeq -M_D M_N^{-1} M_D^T$$

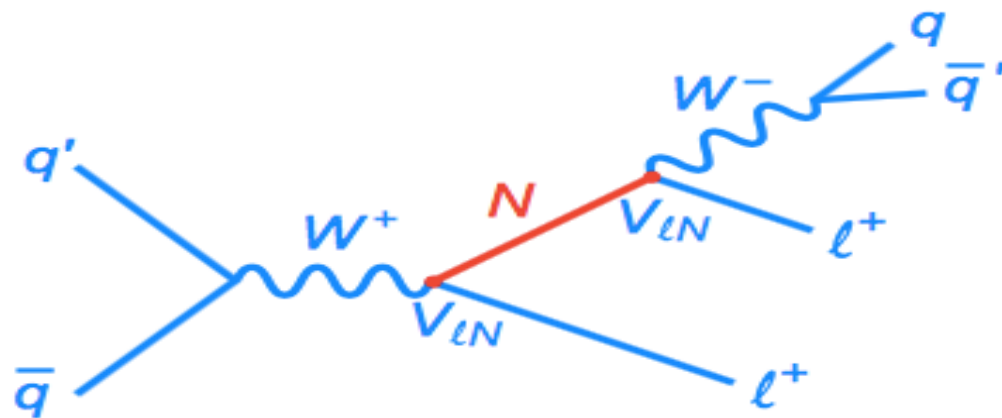
♦ Small Yukawa/heavy-light mixing ( $M_D M_N^{-1}$ )  $\Rightarrow$  sub-TeV 'heavy neutrino'

♦ Cases-Ibarra parameterisation Yukawa coupling expressed in terms of a orthogonal matrix which remains completely arbitrary and hence can be large.

# Seesaw at LHC

\* If heavy neutrino is 'light' enough  
& heavy-light mixing 'large' enough

\* Resonant production of heavy Majorana neutrino



♦ Produces multi-lepton signature at the LHC

$$\begin{aligned} pp &\rightarrow \ell_1^+ N, \quad N \rightarrow \ell_2^+ W^- \\ pp &\rightarrow \ell_1^- \bar{N}, \quad \bar{N} \rightarrow \ell_2^- W^+ \end{aligned}$$

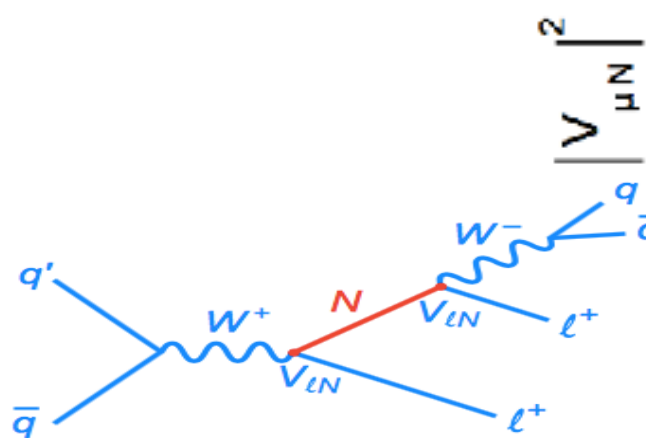
**SSDL + jets**

CMS8: 1501.05566, 1603.02248

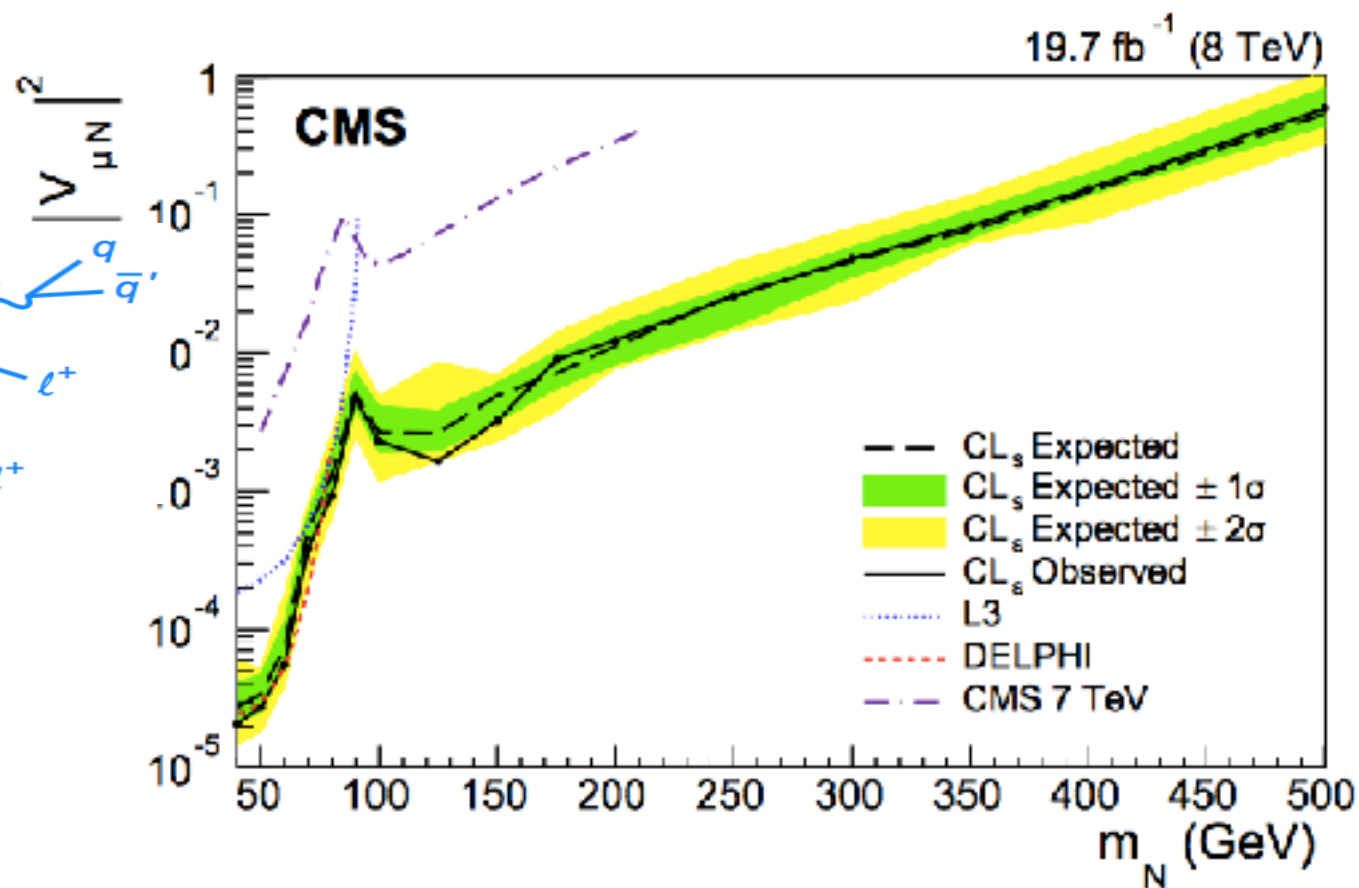
ATLAS8: 1506.06020

Heavy neutrino @ NLO QCD: Das, PK, Majhi

# Seesaw at LHC



**SSDL + jets**

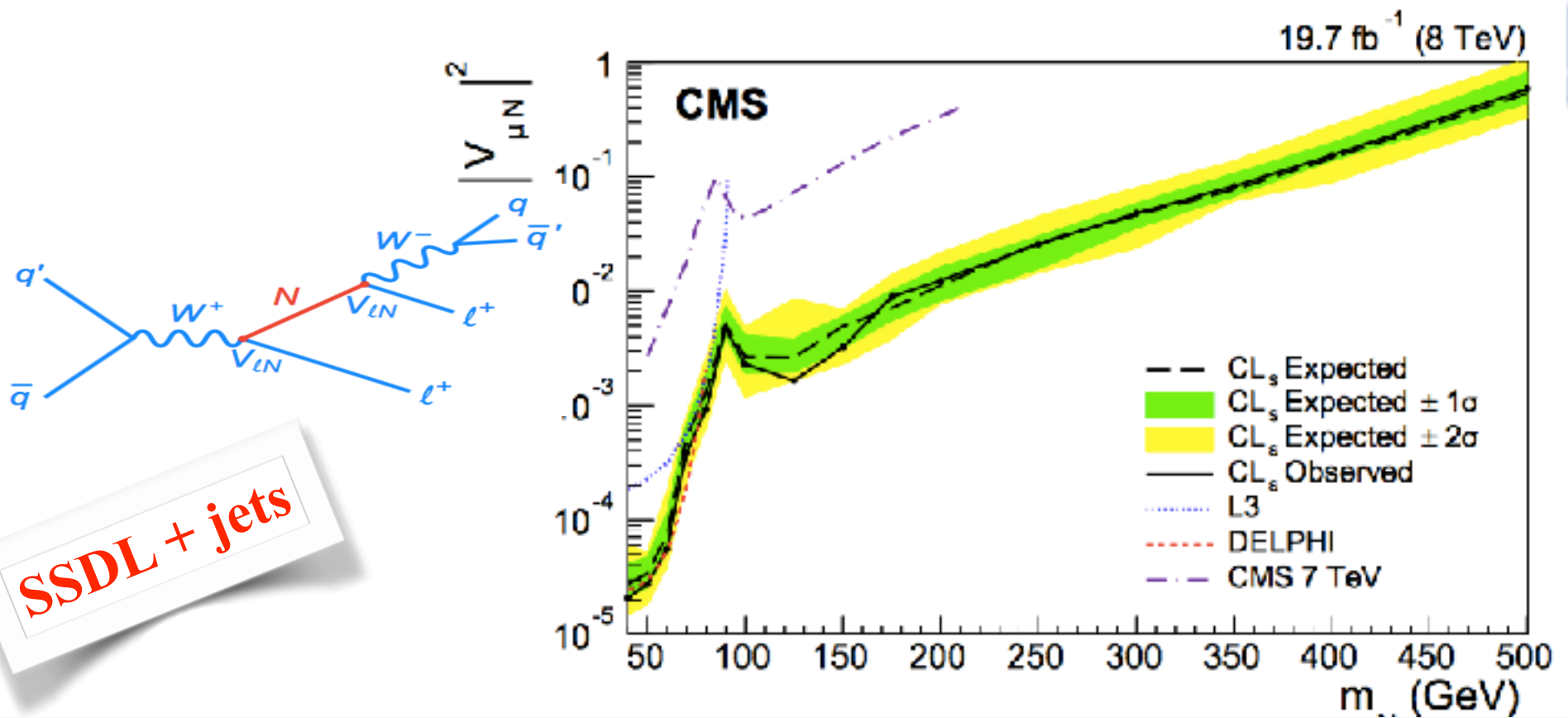


CMS8: 1501.05566, 1603.02248

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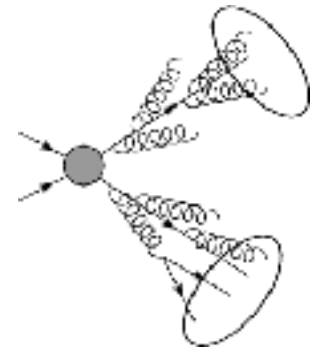
# Seesaw at LHC



If we could use the internal feature of these additional jets,  
to deduce:

- ‘W’-jets
- Production topology
- Control background

# "New" jets @ LHC



- Increased energy of hadron collider

- Exploration for intermediate to very heavy resonance

: if  $PT < 2\text{TeV} \sim 100\text{ GeV}$  splitting is distinguishable with cell size 0.1

- Calorimeters have enough finer resolution & improving

:  $\sim 50 E_{\text{cell}}$  of .1 in  $R=0.4$

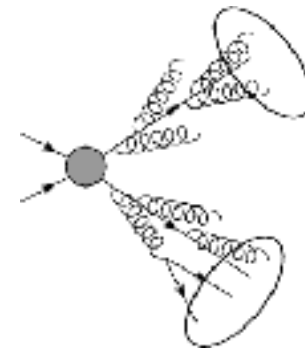
- Computations in fast sequential algorithms

: fast jet algorithm

=> production and analysis of boosted objects & fat jet

# Decay & fragmentation of boosted object Produce a collimated spray of hadron, standard Algorithm consider it as single jet

# "New" jets @ LHC



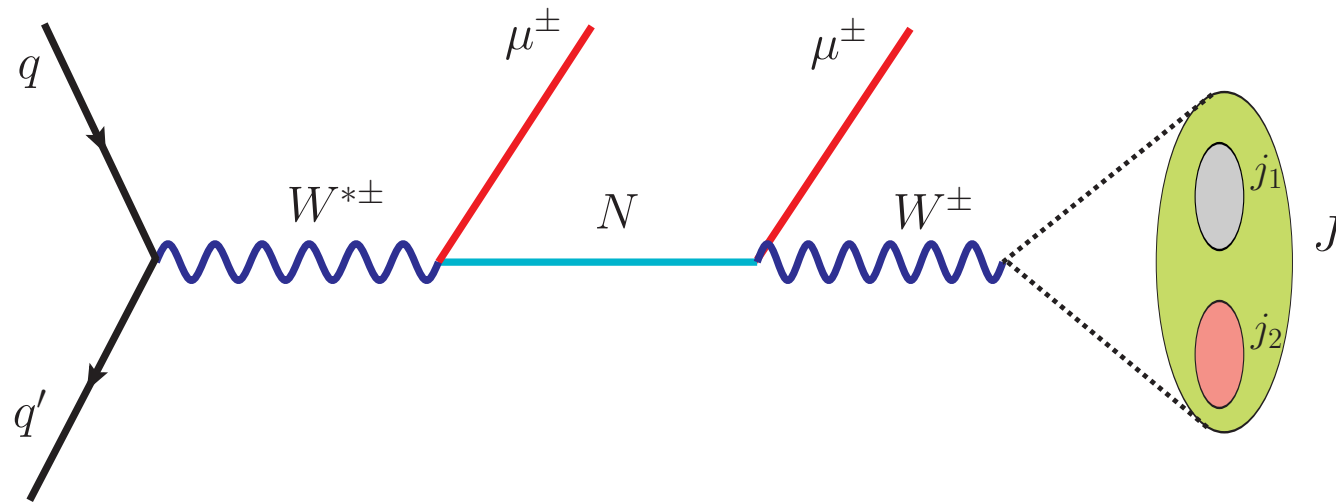
- With same jet mass, boosted hadronic objects have a fundamentally different energy pattern VS to QCD jets
- Clean up smearing effect of jet contamination from IRS, UE, Pileups
- In this work, we exploit the characteristics and kinematic properties for these Boosted W-jets

=> production and analysis of boosted objects & fat jet

# Decay & fragmentation of boosted object Produce a collimated spray of hadron, standard Algorithm consider it as single jet



# Fat jet from Heavy Neutrino



$$pp \rightarrow \ell_1^+ N, \quad N \rightarrow \ell_2^+ W^-$$

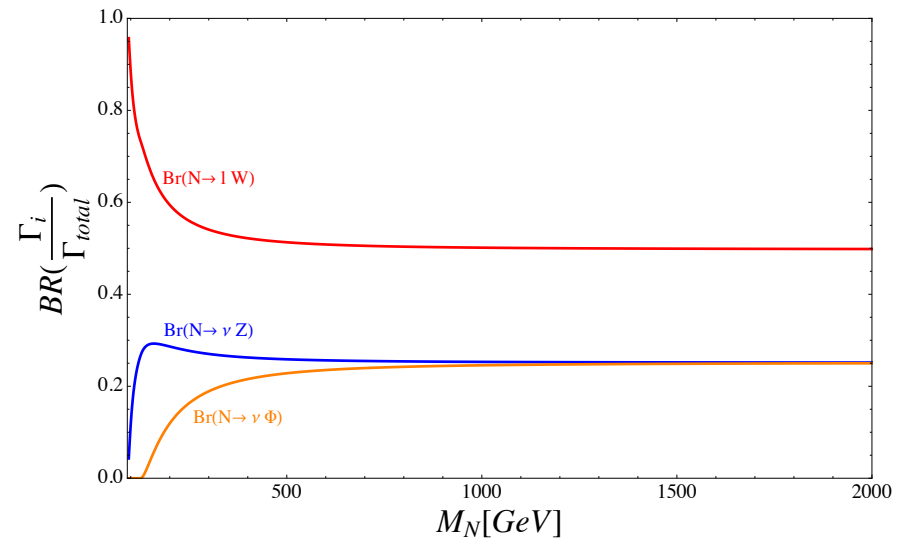
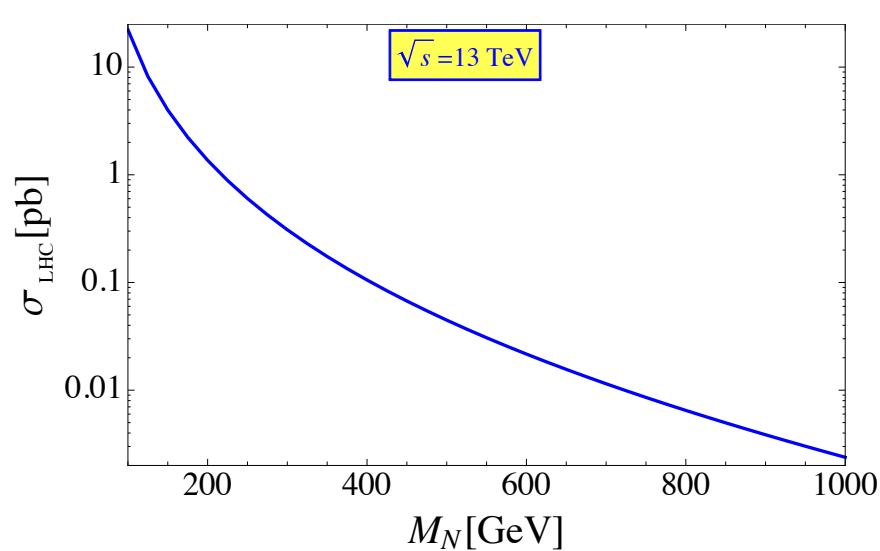
$$pp \rightarrow \ell_1^- \bar{N}, \quad \bar{N} \rightarrow \ell_2^- W^+$$

$$W^- \rightarrow J$$

$$W^+ \rightarrow J$$

Features of these boosted fat-jets are prominent for large  $M_N$   
 $\sim 300$  GeV and above

# Heavy Neutrino at LHC



**Total production cross-section and branching ratios of heavy Majorana neutrino as a function of its mass at the LHC with  $\sqrt{s} = 13$  TeV and normalised by the  $|V_{\mu N}|^2$ .**

$$BR(N \rightarrow \ell W) : BR(N \rightarrow \nu Z) : BR(N \rightarrow \nu H) \simeq 2 : 1 : 1.$$

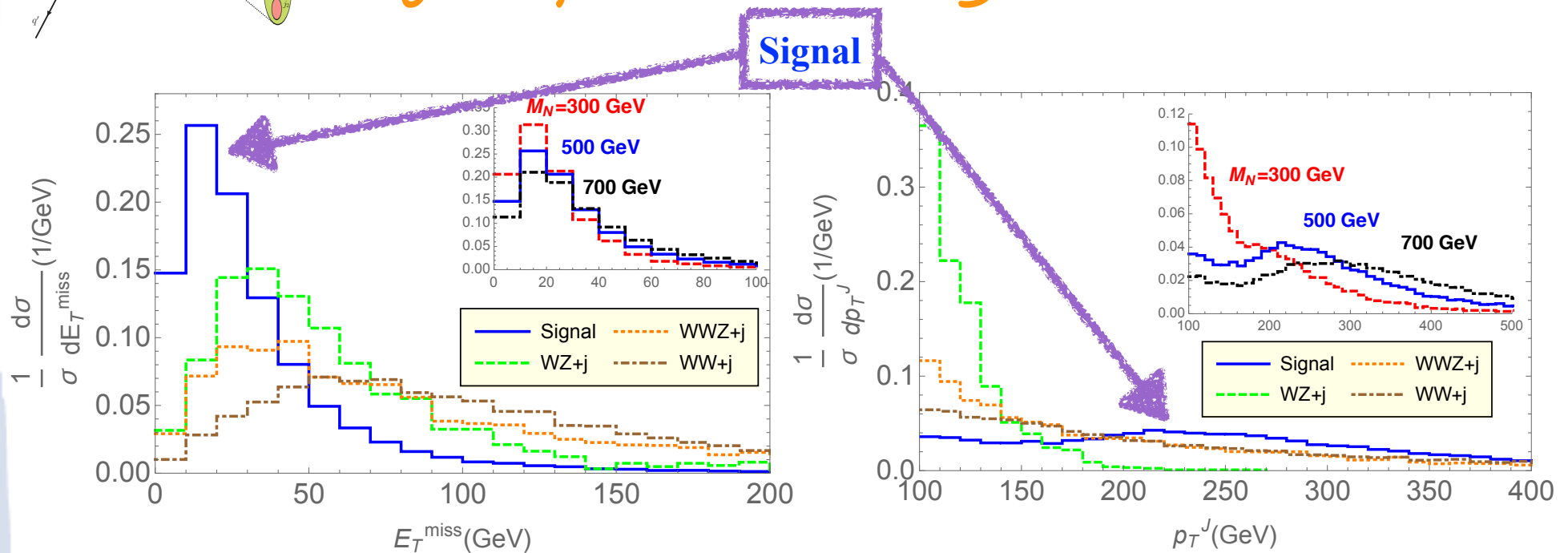


# Fat jet from Heavy Neutrino

- ✱ A simplified model - Only light-heavy mixing for the Muons.
- ✱ Backgrounds for Same-sign-di-muon + fat-jet :  
Electroweak gauge boson decay (Di-boson+j, Tri-boson+j)  
+ “fat-jet” - from a W boson OR mimicked from QCD jets
- ✱ MadGraph5-aMC@NLO => Pythia => Delphes
- ✱ MLM matching on a shower- $k_T$  algorithm with  $p_T$ -ordered showers

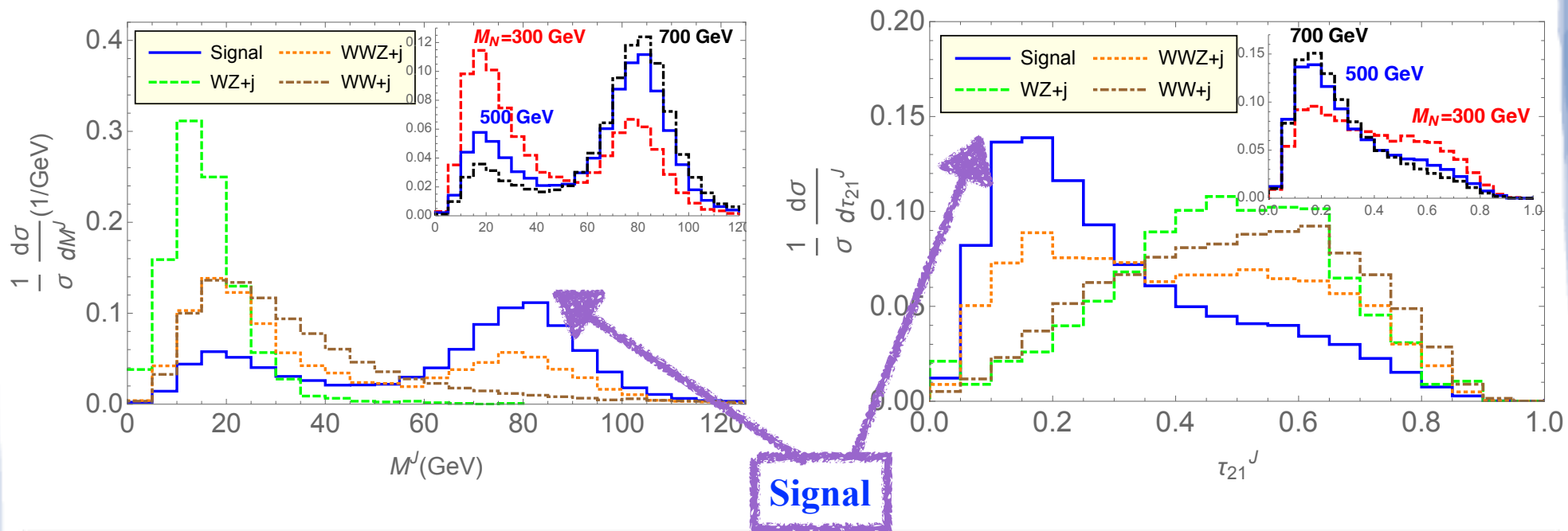
- Same sign  $\mu^\pm$  with  $p_{T^\mu} > 10 \text{ GeV}$ ,  $|\eta^\mu| < 2.4$
- Candidate fat-jets are identified -  $R = 0.8$  CA jet with  $|\eta^J| < 2.4$
- Hardest fat-jet  $\sim W^\pm$  candidate jet (J) having  $p_{T^J} > 100 \text{ GeV}$

# Fat jet from Heavy Neutrino



Normalised distribution as a function of **MET** &  **$p_T^J$**  after the application of the basic selection cuts; including  $p_T^J > 100$  GeV

# Fat jet from Heavy Neutrino

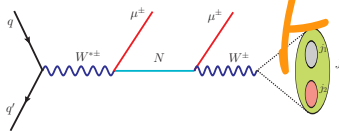


Normalised distribution as a function of **invariant-mass  $M_J$**  & **N-subjet ratio  $\tau_2/\tau_1$**  of fat jet after the application of the basic selection cuts; including  $p_T^J > 100$  GeV

$$\tau_N = \frac{1}{\mathcal{N}_0} \sum_i p_{i,T} \min \{ \Delta R_{i1}, \Delta R_{i2}, \dots, \Delta R_{iN} \} \quad \mathcal{N}_0 = \sum_i p_{i,T} R_i$$

– Inclusive jet shape variable quantify if the original jet consists of N daughter subjets

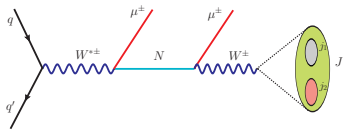
# Fat jet from Heavy Neutrino



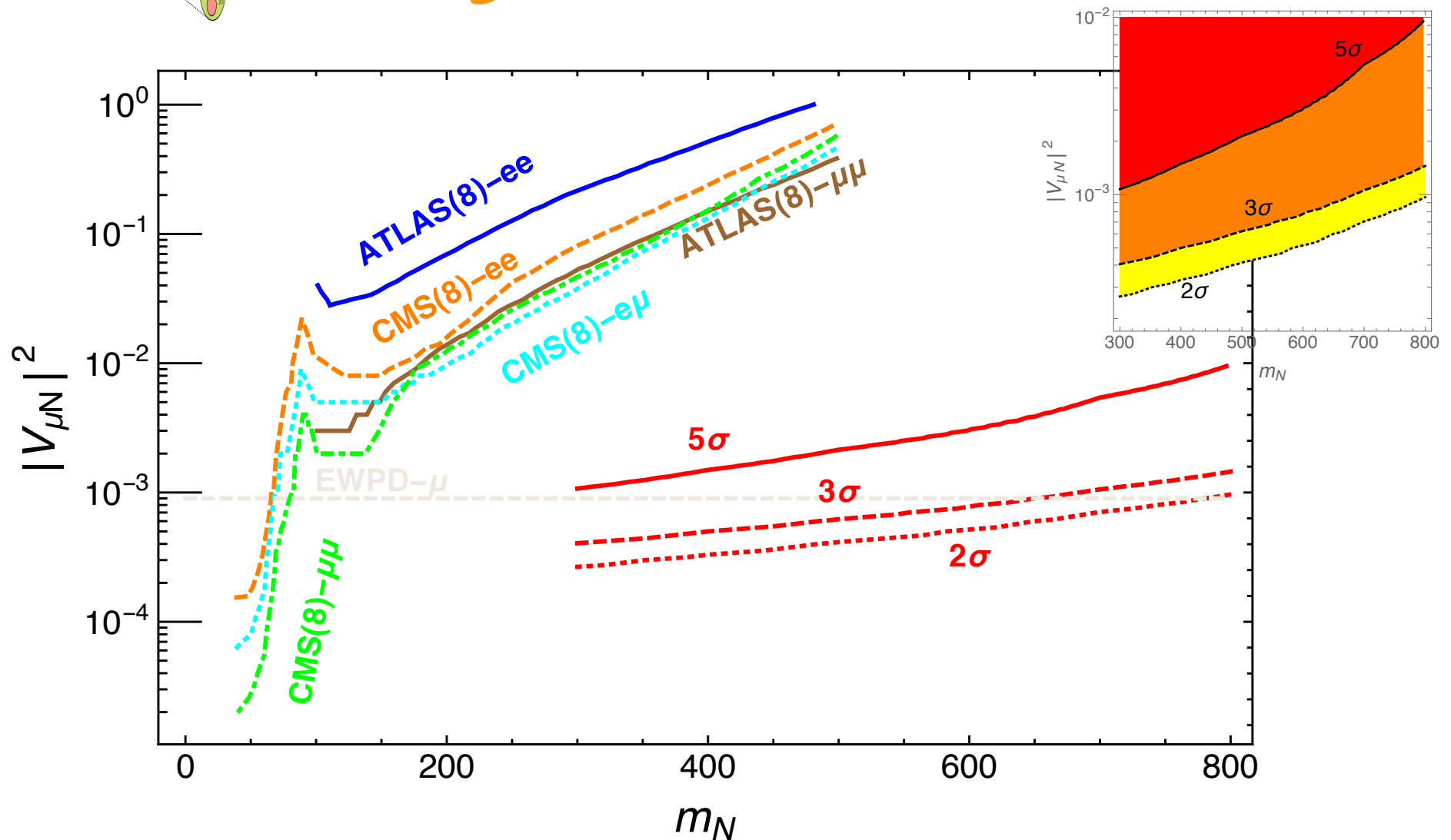
Cut	Signal for $M_N$			Background		
	300 GeV	500 GeV	700 GeV	$WW+j$	$WZ+j$	$WWZ+j$
Pre-selection + $\mu^\pm\mu^\pm + J$ $p_T^J > 100$ GeV	60.0+33.0 [100%]	26.7+17.4 [100%]	14.0+9.5 [100%]	3282.0+3136.5 [100%]	5474.4+4234.5 [100%]	126.05+120.2 [100%]
$p_T(\mu_{1,2}), m_{\mu\mu}$	58.0+ 29.0 [94%]	24.1+ 14.8 [88%]	11.4+6.7 [77%]	2724.30+2575.03 [83%]	3045.2+2796.2 [60%]	104.0+96.7 [82%]
$E_T^{\text{miss}} < 35$ GeV	48.34+20.0 [74%]	20.8+13.2 [77%]	7.3+5.5 [55%]	314.94+197.1 [7.9%]	105.2+104.2 [2.2%]	12.1+ 9.8 [8.9%]
$p_T^J > 150$ GeV	25.6+15.0 [44%]	15.2+ 10.5 [58%]	6.0+4.3 [44%]	184.02+110.4 [4.5%]	20.2+ 15.08 [0.4%]	7.07+6.2 [5.3%]
$M_J > 50$ GeV	21.4+12.3 [36%]	11.2+ 7.4 [42%]	4.8+3.2 [34%]	41.05+32.08 [1.1%]	6.4+4.7 [0.1%]	3.3+2.5 [2.3%]
$\tau_{21}^J < 0.5$	19.5+10.0 [32%]	9.6+5.2 [34%]	3.9+2.0 [25%]	21.1+19.3 [0.6%]	3.3+2.9 [0.06%]	1.5+1.4 [1.2%]

The effectiveness of different variables in minimizing backgrounds is illustrated in the form a cut flow. The two numbers correspond to expected events in  $\mu^+\mu^+$  and  $\mu^-\mu^-$  channels. We adopt a typical mixing angle  $|V_{\mu N}| = 0.03$ . The numbers are for an integrated luminosity of  $3000\text{fb}^{-1}$ , at the 13 TeV LHC.





# Heavy Neutrino at LHC



Exclusion limit in terms of heavy neutrino mass  $M_N$  and  $|V_{\mu N}|^2$  at the 13 TeV LHC with other available limits.

# Outlook

- Non-zero but Tiny neutrino mass can be realised through the simple type-I mechanisms which involves the Majorana type heavy neutrino.
- TeV scale heavy neutrinos can be produced at the LHC through a large mixing angle with the SM light neutrinos
- Such production channel is explored using boosted fat jet substructure which is inevitable for probing large heavy neutrino mass.
- Opens up interesting possibility to explore intermediate to heavy masses of RHN



*Thank You*