

Parameter Space of 'fat-brane' UED models after ATLAS di- photon and multi-jet plus missing transverse energy searches

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Outline

- The Minimal Universal Extra Dimension (mUED) model
- The fat-brane scenario with the mUED
- Decays of level-1 Kaluza-Klein (KK) particles in ‘fat-brane’ scenario
- ATLAS multi-jet and di-photon searches
- Parameter space of ‘fat-brane’ model with ATLAS results
- Conclusions

The minimal Universal Extra Dimension (mUED) model

The mUED model is constructed on $M_4 \times S_1$ with orbifolding $M_4 \times S_1/Z_2$

⇒ An infinite number of KK excitations with mass nR^{-1} for each SM particle.
Quasi-degenerate mass spectra at each KK level.

The mUED has two free parameters:

R^{-1} the size of small extra dim.
 Λ the cut-off scale of the model

With this setup the mUED expects:

- 1) The KK-parity $\equiv (-1)^n$ conservation at the vertices.
- 2) Lightest KK (LKP) is stable.

} The level-1 KK quarks and gluons can only be pair produced at the collider experiments.

The collider signal of level-1 KK quarks and gluons are then multi-jet/lepton plus missing transverse energy.

The ‘fat-brane’ scenario with the mUED

In its simplest realization, ‘fat-brane’ = mUED + ADD

- In general, $(3+m+1)D$ manifold $((3+m)$ -brane) of UED can be embedded in $(4+N)D$ bulk.
- The SM fields live in the 4-brane and only the gravity is allowed to access the full bulk.
- Small spatial dimension, accessible for both matter and gravity, can be viewed as the thickness of the SM 3-brane in $(4+N)D$ bulk.
- The sizes of large and small extra dimensions are $r \sim \text{eV}^{-1}$ to $\text{keV}^{-1} \gg R \sim \text{TeV}^{-1}$
- Both SM particles and the graviton get KK excitations with different masses resulting from different compactification scales.
- The gravity induced interactions respect neither KK number nor KK parity.

The 'fat-brane' scenario with the mUED

After fixing Λ there are three free parameters of the model

- $\nearrow R$ size of small extra dim.
- $\longrightarrow r$ size of large extra dim.
- $\searrow N$ number of large extra dim.

$$M_{\text{Pl.}} = M_D^{N+2} \left(\frac{r}{2\pi} \right)^N \quad [\text{ADD Relation}]$$

The gravitational coupling of the matter fields is given by:

$$\mathcal{S}_{\text{int}} = \int d^{4+N}x \delta(x^5) \dots \delta(x^{4+N}) \sqrt{-\hat{g}} \mathcal{L}_{\text{matter}}$$

The $\mathcal{O}(\hat{K})$ term is

$$\mathcal{S}_{\text{int}} \supset -\frac{\hat{K}}{2} \int d^{4+N}x \delta(x^5) \dots \delta(x^{4+N}) \hat{h}^{\hat{\mu}\hat{\nu}} T_{\hat{\mu}\hat{\nu}}$$

The interaction Lagrangian then becomes:

Han, T., et.al. PRD 59, 105006
Macesanu, C. et.al. PRD 68, 084008

$$\kappa^2 \equiv 16\pi G^{(4)} = V_n^{-1} \hat{\kappa}^2$$

$$m_{\vec{n}} = 2\pi \frac{|\vec{n}|}{r}$$

$$\vec{n} = \{n_1, n_2, \dots, n_N\}$$

$$\mathcal{S}_{\text{int}} - \frac{K}{2} \supset \int d^4x \int_0^{\pi R} dy \sum_{\vec{n}} \left[(h_{\mu\nu}^{\vec{n}} + \eta_{\mu\nu} \phi^{\vec{n}}) T^{\mu\nu} - 2A_{\mu 4}^{\vec{n}} T_4^\mu + 2\phi_{44}^{\vec{n}} T_{44} \right] e^{2\pi i \frac{n_4 y}{r}}.$$

Level \vec{n}

Graviton

Gravi-photon

Gravi-scalar

Decays of level-1 KK particles in ‘fat-brane’ scenario

Two types of decays are possible:

Gravity Mediated Decays (GMD):

- Strong dependence on the number of large extra dimensions, N .
- pronounced when $N = 2$.
- GMD comes to play at the beginning of the decay:

$$q^1, g^1 \rightarrow q/g + G^{\vec{n}}$$

- Pair-produced particle's final state is then multi-jet + MET.

KK-number Conserving Decays (KKDC):

- Independent of the number of large extra dimensions, N .
- pronounced when $N = 6$.
- GMD plays part in the final stage of decay:

$$q^1, g^1 \rightarrow \text{soft-jets/leptons}, \gamma^1$$

$$\gamma^1 \rightarrow \gamma + G^{\vec{n}}$$

- Pair-produced particle's final state is then di-photon + MET.

Decays of level-1 KK particles in 'fat-brane' scenario

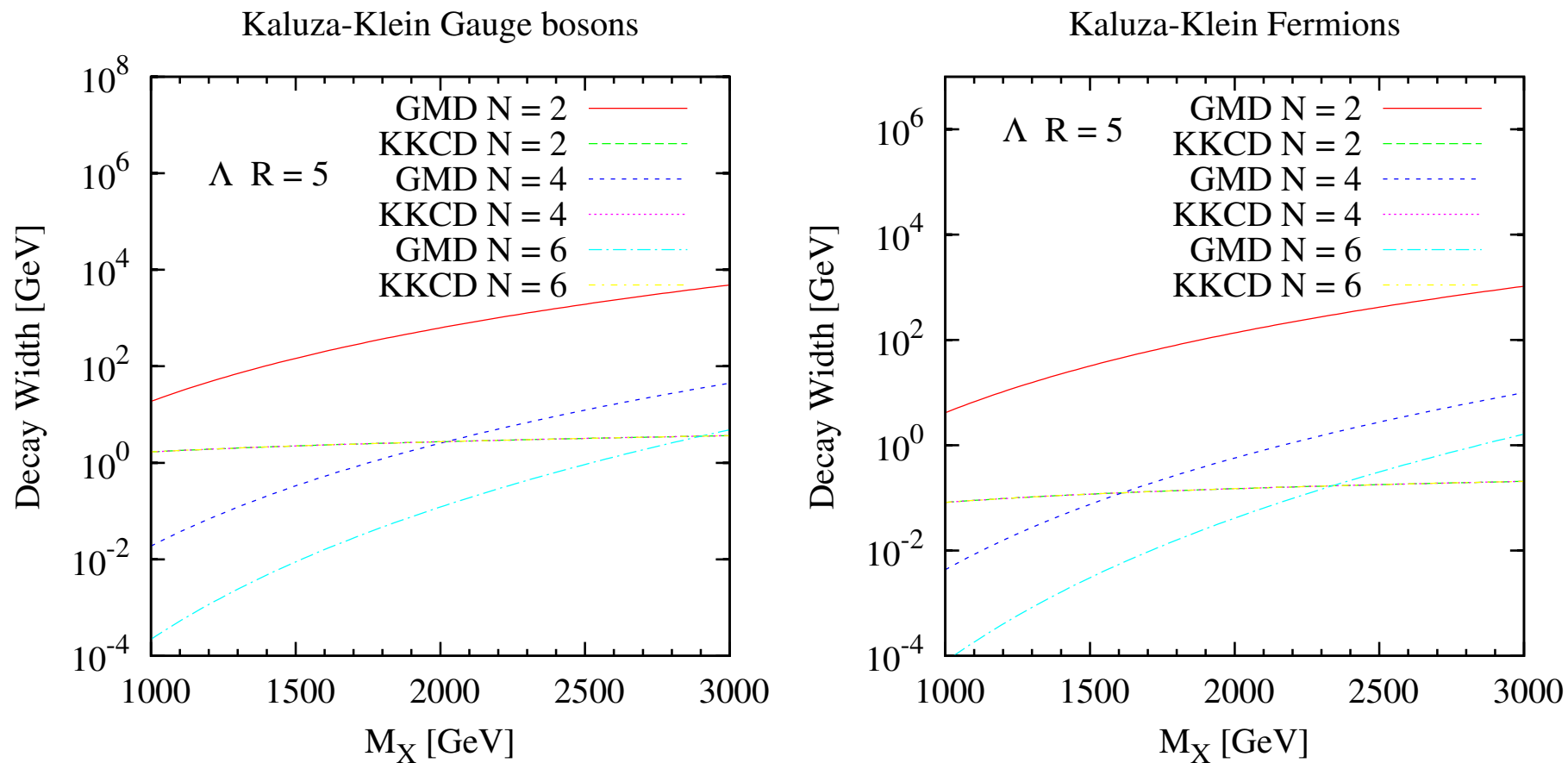


Figure: The decay widths of level-1 KK gauge boson (left) and fermion (right)

Decays of level-1 KK particles in 'fat-brane' scenario

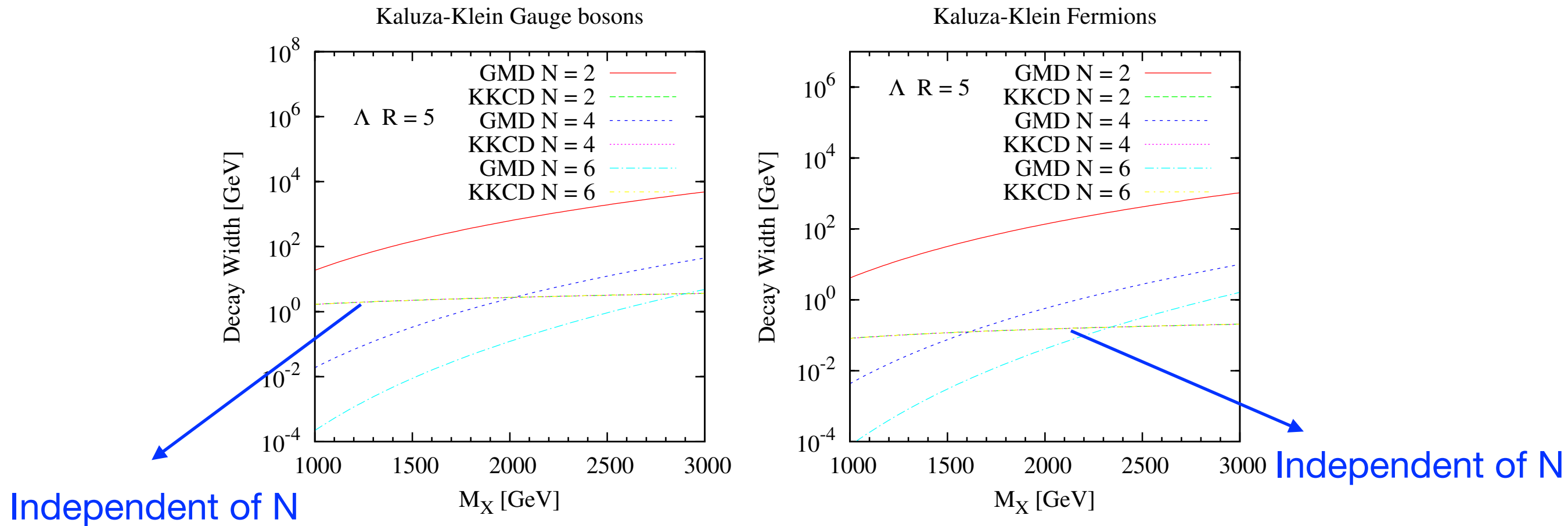


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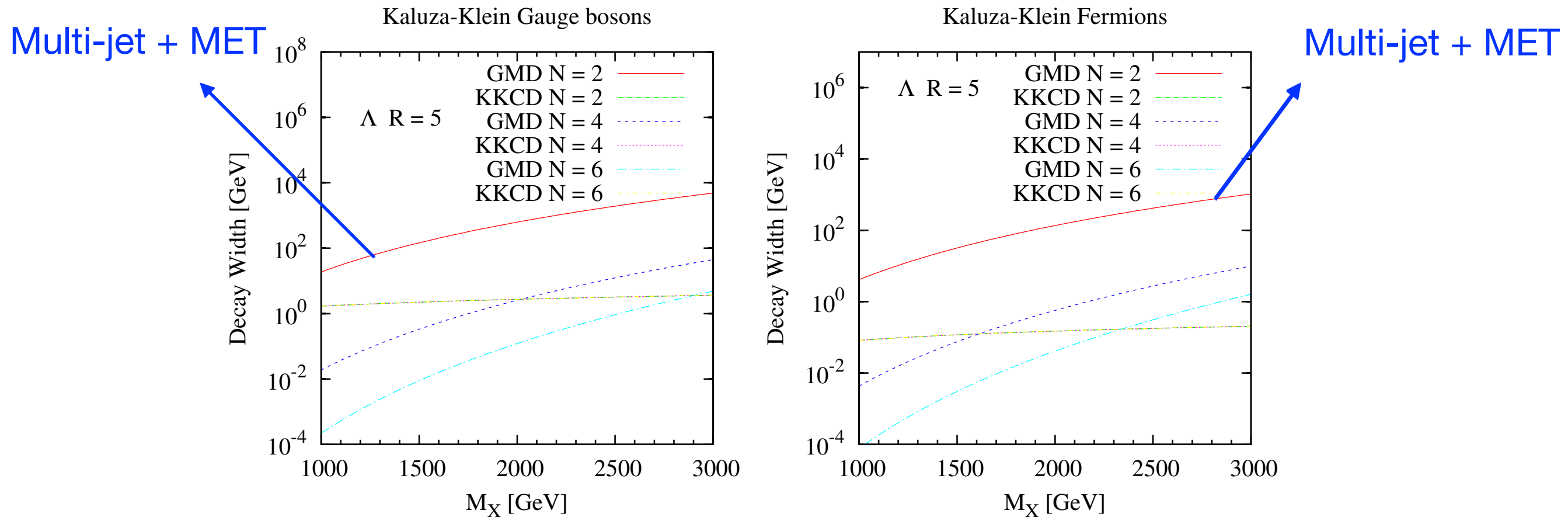


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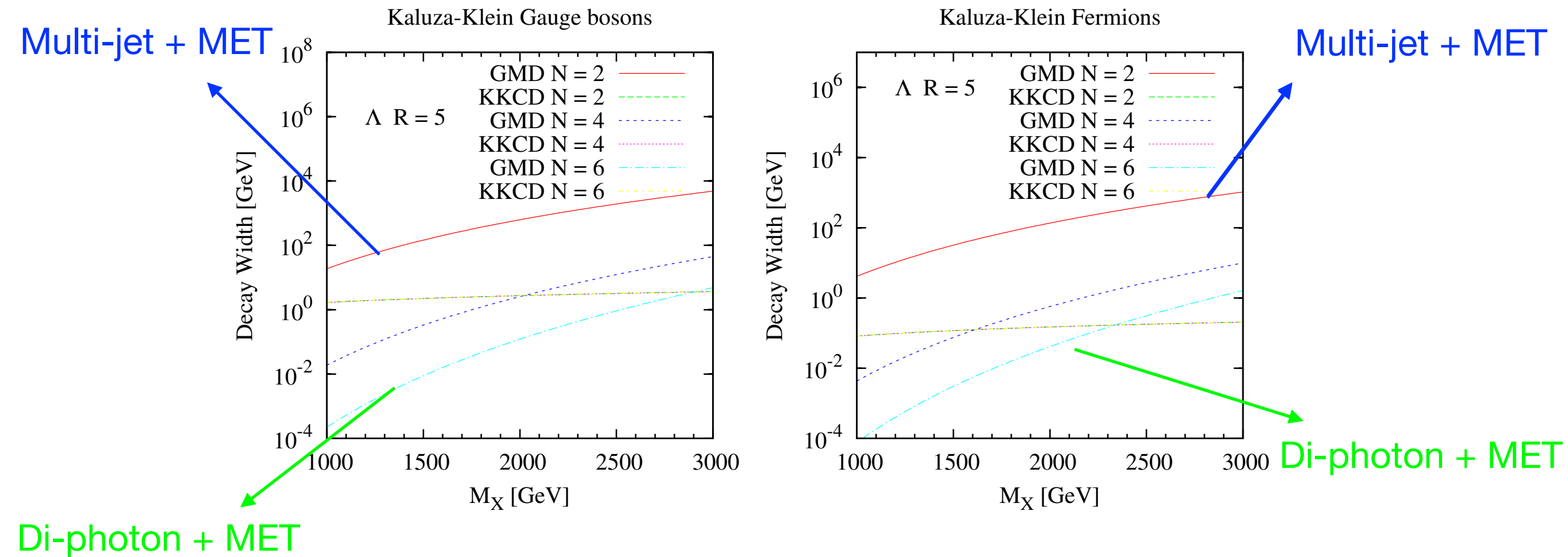


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ATLAS multi-jet and di-photon searches

Both searches look for SUSY signal at 13 TeV with 36.1 (multi-jet) and 3.2 fb^{-1} (di-photon) searches:

Multi-jet search (ATLAS-CONF-2017-022):

- 44 inclusive Signal Regions (SR) in M_{eff} search strategy.
- SRs are characterized by increasing number of jets (2j - 6j) and $m_{\text{eff}}(\text{incl.})$

$$m_{\text{eff}}(\text{incl.}) = \sum_{p_T > 50}^N p_T(j) + \text{MET}$$

- No excess of events above the SM background is observed.
- For each SR %95 C.L. upper limits are put on observable BSM cross-section.

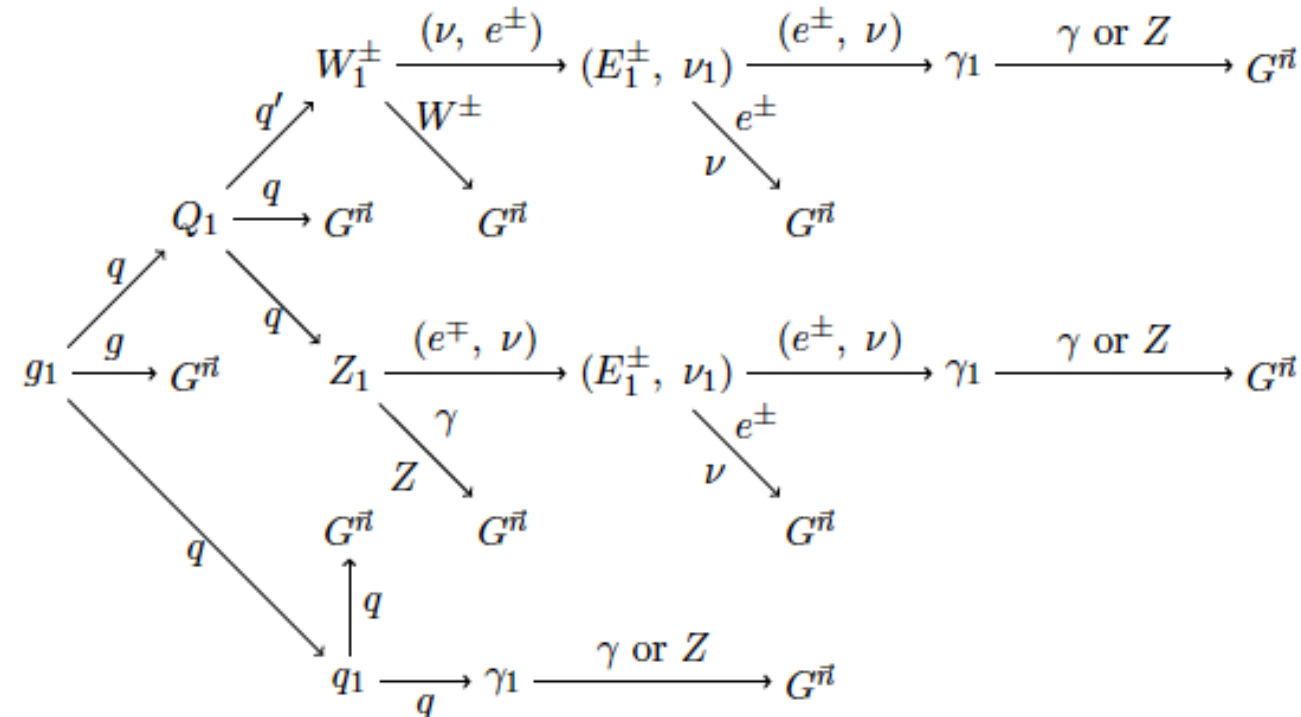
Di-photon search (arXiv:1606.09150):

- 2 tight photons with $p_T > 75$ GeV
- $\Delta\phi_{\text{min}}(\text{jet}, p_T^{\text{miss}}) > 0.5$,
- $E_T^{\text{miss}} > 175$ GeV,
- $m_{\text{eff}} > 1500$ GeV,
- No excess of events above the SM background is observed.
- %95 C.L. upper limit of 0.93 fb on visible BSM cross-section is put.

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1203.1551 [hep-ph]

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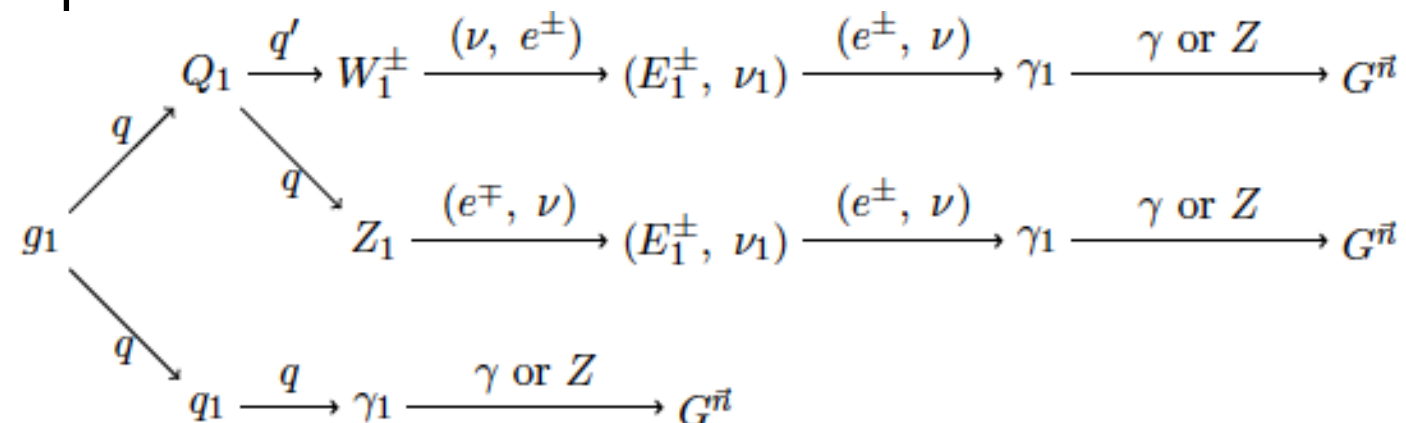
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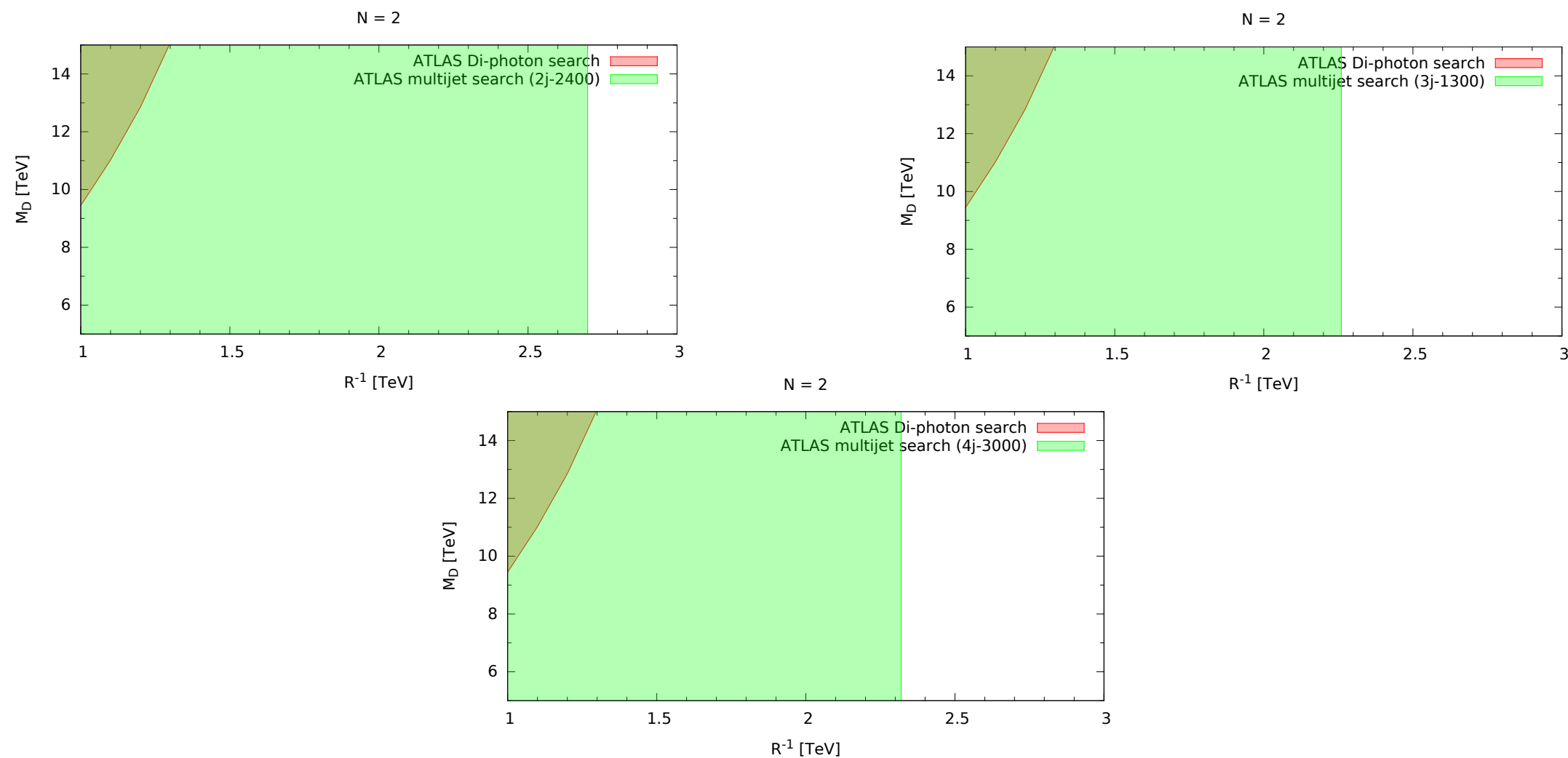
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The parameter space of ‘fat-brane’ model after ATLAS searches

- The model is analyzed for $N = 2, 4$ and 6 and 1 small extra dimension.
- We assumed that $\Lambda R = 5$.
- $q^1 q^1, q^1 g^1$ and $g^1 g^1$ are produced in PYTHIA along with ISR, FSR and fragmentation etc. [PRD 68, 084008 \(2003\)](#)
- GMD's of level-1 KK particles are incorporated. [hep-ph/0510418](#)
- In addition to $\gamma^1 \rightarrow \gamma + G^{\vec{n}}$ dominant decay ($\sim 72\%$), $\gamma^1 \rightarrow Z + G^{\vec{n}}$ ($\sim 28\%$) is also included.

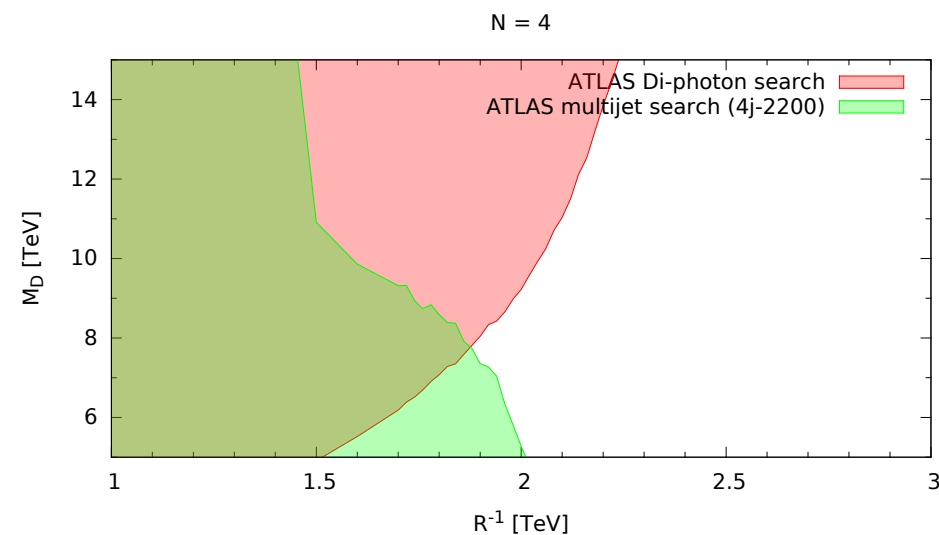
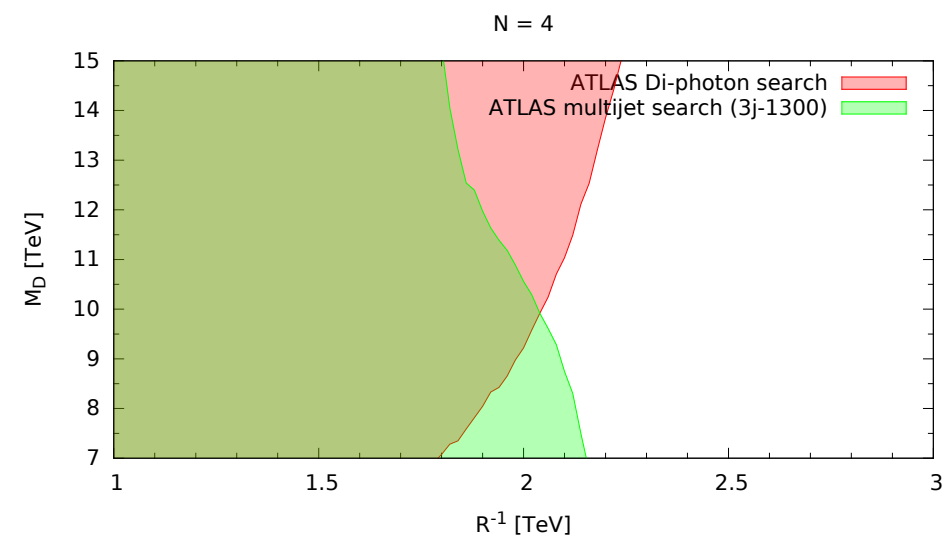
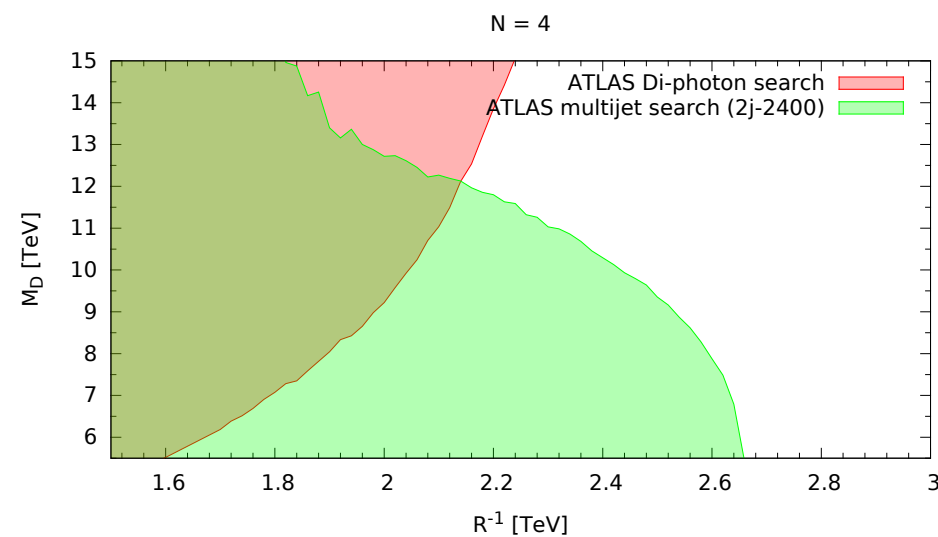
The parameter space of 'fat-brane' model after ATLAS searches [N = 2]

Exclusions from ATLAS Multi-jet and di-photon searches for N = 2 at 13 TeV in 2j, 3j and 4j Signal Regions.



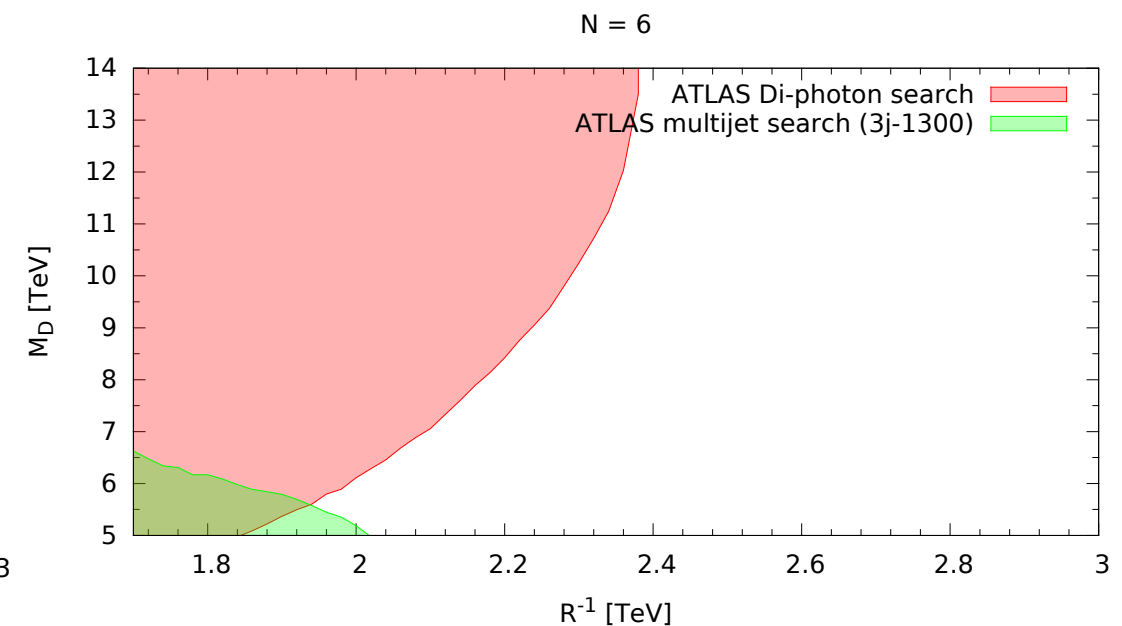
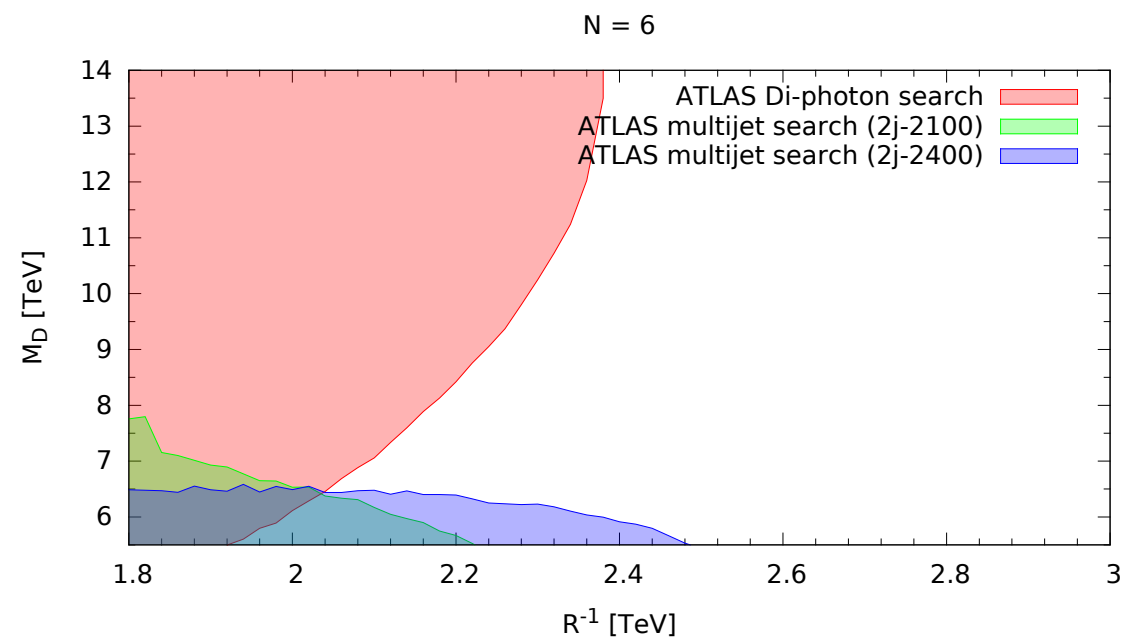
The parameter space of 'fat-brane' model after ATLAS searches [N = 4]

Exclusions from ATLAS Multi-jet and di-photon searches for N = 4 at 13 TeV in 2j, 3j and 4j Signal Regions.



The parameter space of 'fat-brane' model after ATLAS searches [N = 6]

Exclusions from ATLAS Multi-jet and di-photon searches for N = 6 at 13 TeV in 2j, 3j Signal Regions.



Conclusions

- For $N = 2$ the multi-jet search is more pronounced than di-photon in excluding the ‘fat-brane’ model.
- Starting from $N = 4$ the di-photon analysis starts keeping up with the multi-jet results.
- For $N = 6$ exclusion of the parameter space almost dominated by the di-photon analysis alone.
- Multi-jet and di-photon searches are complementary in excluding different parts of the model parameter space.
- For $N = 2, 4$ and 6 $(R^{-1}, M_D) = \left\{ (1.3, 15), (2.1, 12), (2.1, 6.5) \right\}$ TeV are excluded by both searches.

Thank you

Backup

The total decay width is obtained by summing over all possible gravity excitations with mass smaller than the decaying particle:

$$\Gamma = \sum_{\vec{n}} \Gamma_{\vec{n}} = \sum_{\vec{n}} [\Gamma_{h\vec{n}} + \Gamma_{A\vec{n}} + \Gamma_{\phi\vec{n}}]$$

The gravity KK-states are nearly degenerate in mass is given by $\Delta m = 2\pi/r \sim \text{eV to keV}$

The mUED mass spectra:

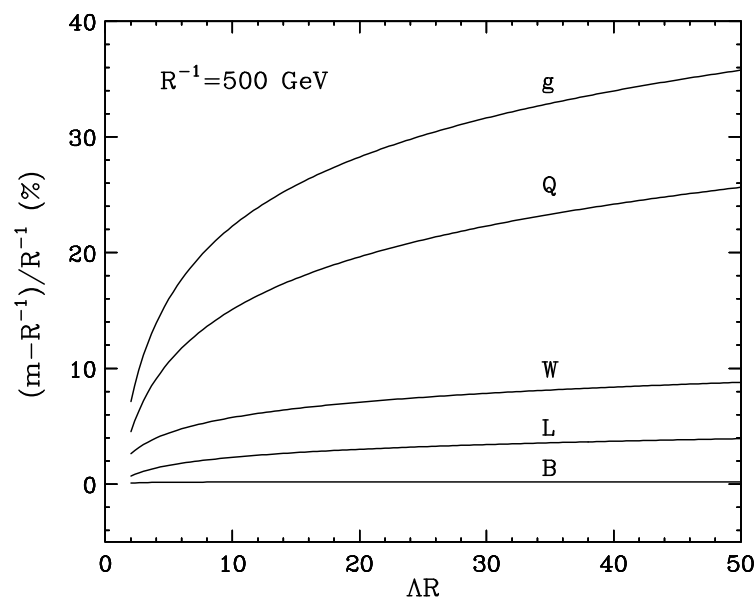


FIG. 2: Radiative corrections (in %) to the spectrum of the first KK level for $R^{-1} = 500 \text{ GeV}$, versus ΔR .

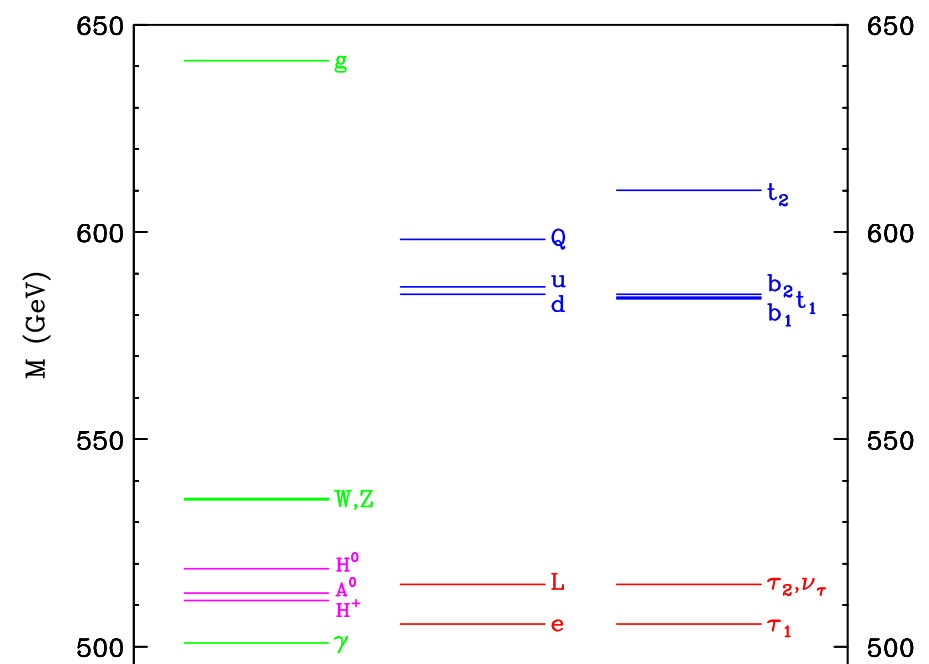


FIG. 1: One-loop corrected mass spectrum of the first KK level in MUEDs for $R^{-1} = 500 \text{ GeV}$, $\Delta R = 20$ and $m_h = 120 \text{ GeV}$.

hep-ph/020531

Backup

The (4+N)D metric, assumed to be approximately flat, is $\hat{g}_{\hat{\mu}\hat{\nu}} = \eta_{\hat{\mu}\hat{\nu}} + \hat{\kappa}\hat{h}_{\hat{\mu}\hat{\nu}}$ where

$$\hat{\kappa}^2 = 16\pi G^{(4+N)}$$

The higher dimensional (4+N) tensor consists of:

$$\hat{h}_{\hat{\mu}\hat{\nu}} = V_N^{-1/2} \begin{pmatrix} h_{\mu\nu} + \eta_{\mu\nu}\phi & A_{\mu i} \\ A_{\nu j} & 2\phi_{ij} \end{pmatrix}$$

The fields can be KK expanded as:

$$h_{\mu\nu}(x, y) = \sum_{\vec{n}} h_{\mu\nu}^{\vec{n}}(x) \exp\left(i\frac{2\pi\vec{n} \cdot \vec{y}}{r}\right),$$

$$A_{\mu i}(x, y) = \sum_{\vec{n}} A_{\mu i}^{\vec{n}}(x) \exp\left(i\frac{2\pi\vec{n} \cdot \vec{y}}{r}\right),$$

$$\phi_{ij}(x, y) = \sum_{\vec{n}} \phi_{ij}^{\vec{n}}(x) \exp\left(i\frac{2\pi\vec{n} \cdot \vec{y}}{r}\right), \quad \vec{n} = \{n_1, n_2, \dots, n_N\}$$