

CONSTRAINING WARM INFLATION WITH CMB OBSERVATIONS

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arXiv:1710.10008

Presentation Outline

1 MOTIVATION

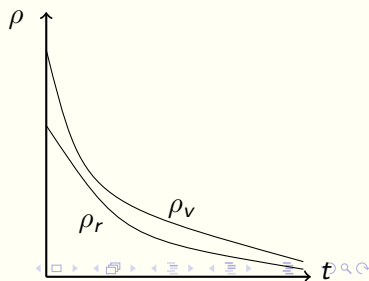
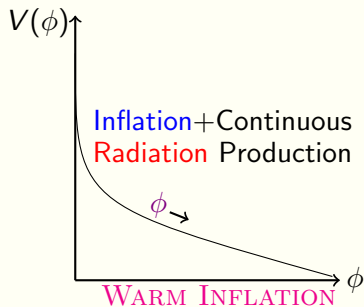
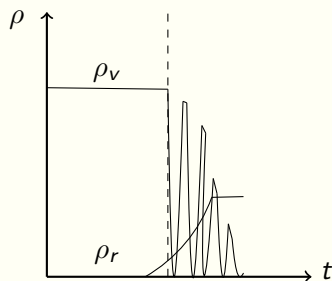
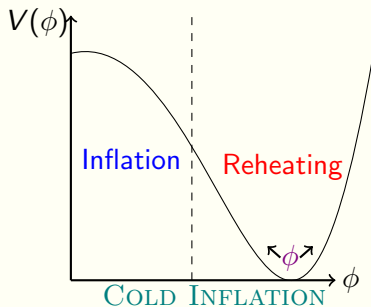
2 THEORY

3 ANALYSIS

4 CONCLUSION

5 REFERENCES

WARM VS COLD INFLATION



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WARM INFLATION

- Dissipation of inflaton (ϕ) energy $\Rightarrow$ Dissipative term : Υ
- Radiation is produced through friction and continuously replenish through the inflaton field decay.

$$\ddot{\phi} + (3H + \Upsilon)\dot{\phi} + V_{,\phi} = 0$$

$$\dot{\rho}_r + 4H\rho_r = \Upsilon \dot{\phi}^2$$

- Inflaton is a pNGB from a broken gauge symmetry (inspired by the **L**ITTLE **H**IGGS **M**ECHANISM).
- Quartic potential $V(\phi) = \lambda\phi^4$ is tested in the linear dissipation regime $\Upsilon = C_T T$.

$$\rho_r = \frac{\pi^2}{30} g_* T^4 = C_R T^4$$

- $Q = \Upsilon/3H$

POWER SPECTRUM

- Q_e and ϕ_e can be computed from:

$$\frac{Q_e^3}{(1 + Q_e)} = \frac{C_T^4}{2^4 \times 3^5 \times C_R \lambda}$$

$$Q^3(1 + Q)^2 = \frac{4}{9} \left(\frac{C_T^4}{C_R \lambda} \right) \left(\frac{m_P}{\phi} \right)^6$$

- $Q(k)$ is determined from:

$$\frac{dQ}{dN} \simeq \frac{Q}{3 + 5Q} (6\epsilon_\phi - 2\eta_\phi)$$

- Scalar power spectrum is given as:

$$P_{\mathcal{R}} = \frac{C_T^4}{4\pi^2 \times 36 C_R} Q_*^{-3} \left[\frac{3Q_*}{C_T} \frac{2\pi Q_*}{\sqrt{1 + 4\pi Q_*/3}} + 1 + 2\mathcal{N}_* \right] \times G[Q_*]$$

- The tensor spectrum is defined to be:

$$P_T = 8 \left(\frac{H_*}{2\pi m_P} \right)^2 = \frac{8\lambda^{1/3}}{4\pi^2} \left(\frac{4C_T^4}{9C_R} \right)^{2/3} \frac{1}{Q_*^2(1 + Q_*)^{2/3}}$$

MODEL PARAMETERS

- Model parameters: C_T , λ , g_* .
- No. of e-folds ($N_e(k)$) is defined as:

$$56.12 - \ln \frac{k}{k_0} + \frac{\ln \frac{2}{3}}{3(1 + \tilde{w})} + \ln \frac{V_k^{1/2}}{V_{end}^{1/2}} + \frac{1 - 3\tilde{w}}{3(1 + \tilde{w})} \ln \frac{\rho_{RH}^{1/4}}{V_{end}^{1/4}} + \ln \frac{V_{end}^{1/4}}{10^{16}}$$

- $\tilde{w} = 1/3$ reduces $N_e(k)$.

MODEL PARAMETERS

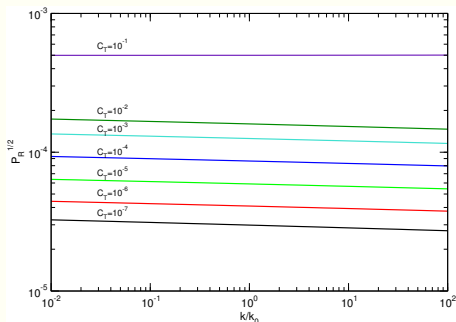
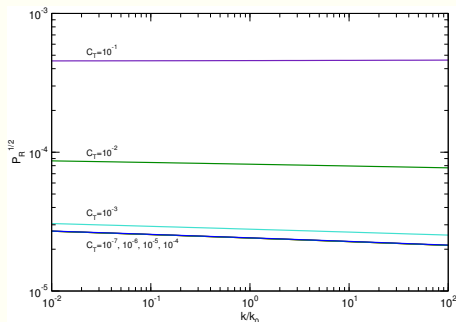
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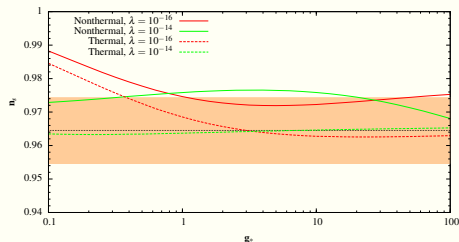
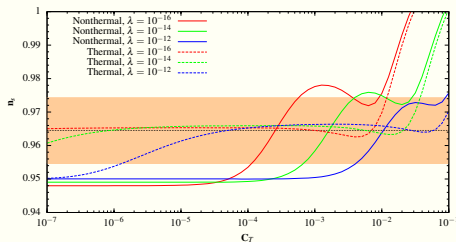
$$N(k) = 56.02 - \ln \frac{k}{k_0} + \ln \frac{V_k^{1/2}}{V_{end}^{1/2}} + \ln \frac{V_{end}^{1/4}}{10^{16} \text{ GeV}}$$

SCALAR POWER SPECTRUM



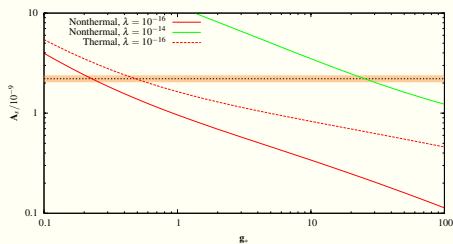
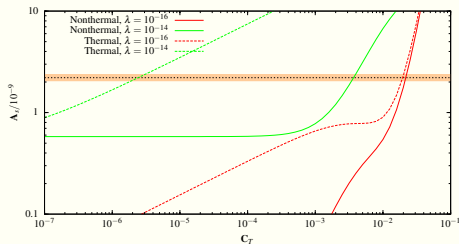
Primordial spectrum as a function of k/k_0 , for different values of the parameter $C_T = 10^{-7}, 10^{-6}, \dots, 10^{-1}$ and for fixed $\lambda = 10^{-14}$, $g_* = 12.5$. LHS is for a non-thermal inflaton, i.e., $\mathcal{N}_* = 0$ and RHS is for a thermal inflation, i.e., $\mathcal{N}_* \neq 0$.

BACKGROUND DEPENDENCE



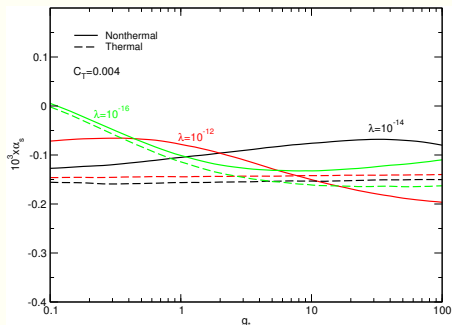
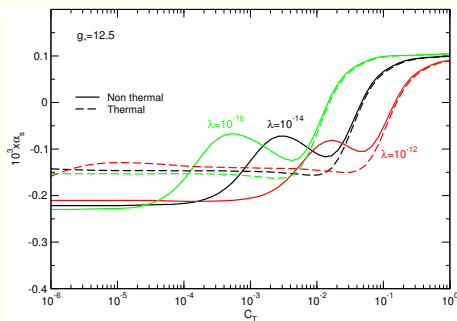
Spectral index (n_s) as a function of C_T with $g_* = 12.5$ in LHS and as a function of g_* with $C_T = 0.004$ in RHS. The solid lines are for $\mathcal{N}_* = 0$ and the dashed lines are for $\mathcal{N}_* \neq 0$.

BACKGROUND DEPENDENCE



Amplitude of spectrum (A_s) as a function of C_T with $g_* = 12.5$ in LHS and as a function of g_* with $C_T = 0.004$ in RHS for different values of λ as indicated in the plot. The solid lines are for $\mathcal{N}_* = 0$ and dashed lines are for $\mathcal{N}_* \neq 0$.

BACKGROUND DEPENDENCE



Running of the spectral index (α) as a function of C_T with $g_* = 12.5$ in LHS, and as a function of g_* with $C_T = 0.004$ in RHS, for different values of λ as indicated in the plot. The solid lines are for $\mathcal{N}_* = 0$ and dashed lines are for $\mathcal{N}_* \neq 0$.

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2 THEORY

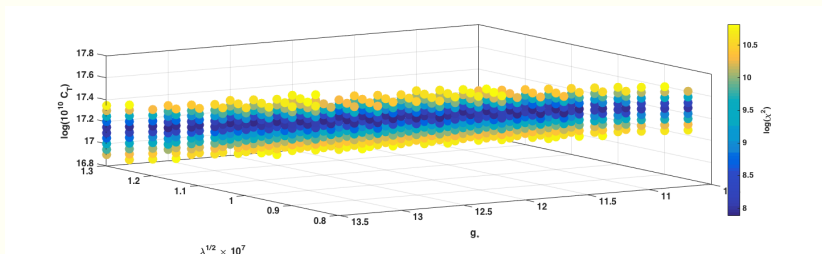
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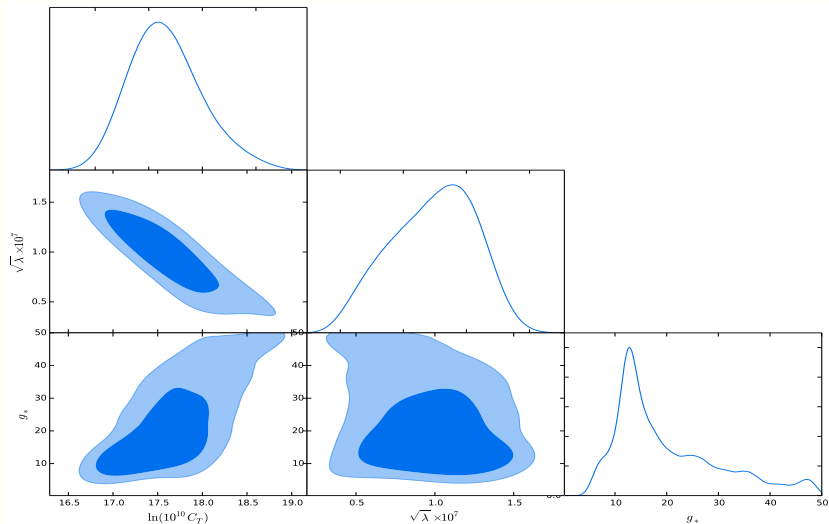
MCMC METHODOLOGY

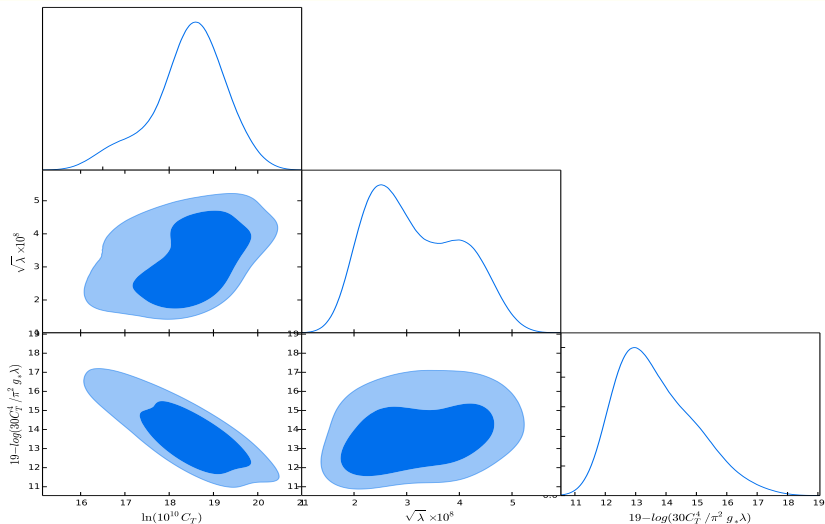
- MCMC is performed using COSMOMC package coupled with CAMB.
- Slow mixing and bad convergence in COSMOMC due to multimodality in the theory.



- Increase the temperature of the chains and changing the standard Metropolis-Hastings algorithm to Wang-Landau algorithm.
- In thermal case, hierarchical centering is employed to solve convoluted multimodality.

NON-THERMAL CASE ($N_* = 0$)



THERMAL CASE ($N_* \neq 0$)

RESULTS

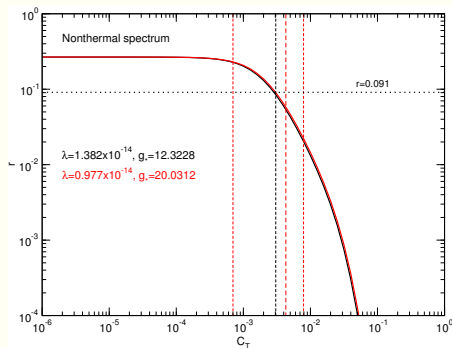
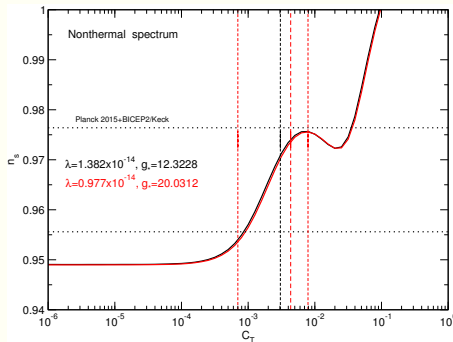
Constraints on cosmological parameters for non-thermal and thermal case compared with Λ CDM + r using Planck 2015+BICEP2/Keck Array observations.

Warm Inflation					Cold Inflation		
	$\mathcal{N}_* = 0$		$\mathcal{N}_* \neq 0$			Λ CDM + r	
parameters	mean value	1σ	mean value	1σ	parameters	mean value	1σ
$\Omega_b h^2$	0.02233	0.00022	0.02224	0.00019	$\Omega_b h^2$	0.02224	0.00017
$\Omega_c h^2$	0.1178	0.0015	0.1194	0.0013	$\Omega_c h^2$	0.1192	0.0016
$100\theta_{MC}$	1.04097	0.00046	1.04088	0.00038	$100\theta_{MC}$	1.04085	0.00034
τ	0.077	0.019	0.068	0.021	τ	0.064	0.018
C_T	0.0043	0.0018	0.0104	0.0077	$\ln(A_s \times 10^{10})$	3.06	0.031
λ	9.77×10^{-15}	5.41×10^{-15}	9.74×10^{-16}	6.78×10^{-16}	n_s	0.966	0.0052
g_*	20.03	10.39	139.91	487.98	r	< 0.07	

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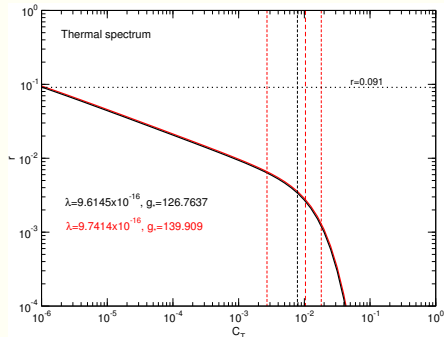
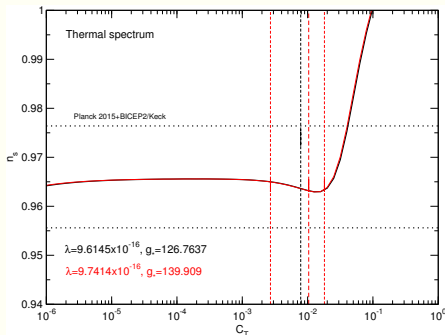
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NON-THERMAL CASE ($N_* = 0$)



The predictions for the **spectral index** and **tensor-to-scalar ratio** for the **best-fit** and **mean value** of parameters for non-thermal case ($N_* = 0$).

THERMAL CASE ($N_* \neq 0$)



The predictions for the **spectral index** and **tensor-to-scalar ratio** for the **best-fit** and **mean value** of parameters for thermal case ($N_* \neq 0$).

DISCUSSIONS

- No reheating is required in warm inflation $\Rightarrow$ Warm exit to the radiation dominated universe.
- Model parameters are kept free and a statistical simulation is carried out to constrain the parameters and inflationary observables.
- The difficulties related to the MCMC simulations are dealt with more sophisticated statistical tweaks.
- Non-thermal case ($N_* = 0$): Spectral index is within the observational bounds but the tensor-to-scalar ratio is close to the upper bound.
- Thermal case ($N_* \neq 0$): tensor-to-scalar ratio is well within the observational bound.

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- S. Bartrum, M. Bastero-Gil, A. Berera, R. Cerezo, R. O. Ramos and J. G. Rosa, Phys. Lett. B 732, 116 (2014) [arXiv:1307.5868 [hep-ph]].
- M. Bastero-Gil, A. Berera, R. O. Ramos and J. G. Rosa, Phys. Rev. Lett. 117 (2016) no.15, 151301 [arXiv:1604.08838 [hep-ph]].
- [READ THE FULL STORY @ arXiv:1710.10008.](#)

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THANK YOU !!