The Dirty Picture Low Energy QCD at the LHC

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Outline

- Introduction
- 2 Event Shape Variables
- b-quark Fragmentation
- Jet Charge Observables
- Charged Particles



QCD LHC

Large Hadron Collider –

Quantum ChromoDynamics in its full glory & gory details
from hard scattering, ISR, FSR to fragmentation, hadronization
from perturbative/analytical calculations to phenomenological models

Devil is in the detail

- strong coupling
- fragmentation models,
- color reconnection,
- multi-parton interaction (MPI) ...

Good agreement between theory and experiment for the hard scattering part over many orders of magnitudes - jet cross-section, PDF, ... (BM, SG, SSG, ... review talks).



QCD Thali

The limited menu on offer:

- Event Shape Variables
- Heavy Quark Fragmentation
- Jet Charge
- Charged particles multiplicity, p_T and η ,



Why the ESVs

ESV variables are sensitive to various aspects of QCD –

- hard scattering part analytically calculable with good accuracy
 - Complement of Transverse Thrust τ_{\perp} , Thrust Minor $T_{m,\perp}$
 - Third-Jet Resolution Parameter Y₂₃
 - Sphericity S, Transverse Sphericity S⊥, Aplanarity A
- hadronization part described by phenomenological models
 - jet broadening B_T
 - total jet mass ρ_{Tot} , total transverse jet mass ρ_{Tot}^{T}
 - particle multiplicity and p_T in event, jets

Event Shape Variables have been calculated using

- jets clustered charged and neutral particles in an event
- unclustered charged particles in an event



Definitions

Complement of Transverse Thrust:

$$\tau_{\perp} \equiv 1 - T_{\perp} = 1 - \max_{\overrightarrow{n}_{T}} \frac{\sum_{i} |\overrightarrow{p}_{T,i} \cdot \overrightarrow{n}_{T}|}{\sum_{i} p_{T,i}}, \ \ 0 \le \tau_{\perp} \le 1 - 2/\pi$$

It is sensitive to the hard scattering process -

Thrust Minor: A measure of the total $p_{T,i}$ out of the event plane:

$$T_{m,\perp} = \frac{\sum_{i} |\overrightarrow{p}_{T,i} \times \overrightarrow{n}_{T}|}{\sum_{i} p_{T,i}}, \quad 0 \leq T_{m,\perp} < 2/3$$

Third-Jet Resolution Parameter:

$$Y_{23} = \frac{\min(p_{T,3}^2, [\min(p_{T,i}, p_{T,j})^2 \times (\Delta R_{ij})^2 / R^2])}{P_{12}^2} (CMS)$$

$$= \frac{p_{T,3}^2}{H_{T,2}^2} (ATLAS)$$



Definitions

Sphericity, Transverse Sphericity, Aplanarity:

$$M_{xyz} = \sum_{i} \begin{pmatrix} p_{x,i}^{2} & p_{x,i}p_{y,i} & p_{x,i}p_{z,i} \\ p_{y,i}p_{x,i} & p_{y,i}^{2} & p_{y,i}p_{z,i} \\ p_{z,i}p_{x,i} & p_{z,i}p_{y,i} & p_{z,i}^{2} \end{pmatrix}, \quad \lambda_{1} > \lambda_{2} > \lambda_{3}, \quad \sum_{i} \lambda_{i} = 1$$

$$S = \frac{3}{2}(\lambda_2 + \lambda_3), \quad S_{\perp} = \frac{2\lambda_2}{\lambda_1 + \lambda_2}, \quad A = \frac{3}{2}\lambda_3$$

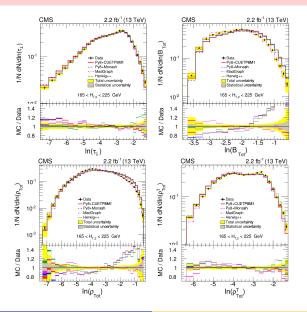
Jet Broadening:

$$\eta_X \equiv \frac{\sum_{i \in X} p_{T,i} \eta_i}{\sum_{i \in X} p_{T,i}} \;, \;\; \phi_X \equiv \frac{\sum_{i \in X} p_{T,i} \phi_i}{\sum_{i \in X} p_{T,i}}$$

$$B_X \equiv rac{1}{2 P_T} \sum_{i \in X} p_{T,i} \sqrt{(\eta_i - \eta_X)^2 + (\phi_i - \phi_X)^2} \;,\; B_T = B_U + B_L$$



CMS ESV with Jets $\sqrt{s} = 13 TeV$



PF jets used $|\eta^{jet}| <$ 2.4, $p_T^{jet} >$ 30 GeV, $N_{iet} \geq$ 3

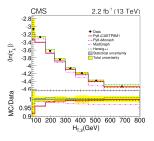
Events divided into 8 ranges of $H_{T,2} = (p_{T,jet1} + p_{T,jet2})/2$

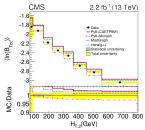
Unfolded data distributions compared with Pythia 8 (CUETP8M1, Monash), Herwig, Madgraph.

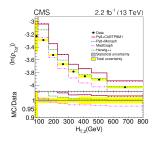
JHEP12(2018)117

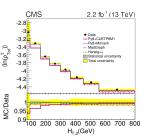
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CMS ESV with Jets $\sqrt{s} = 13 TeV$









Agreement improves with $H_{T,2}$ Agreement good for τ_{\perp} , ρ_{Tot}^{T} Herwig - better agreement for B_{T} , ρ_{Tot} For higher $H_{T,2}$, events become less spherical.

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Low pileup data collected in 2018, $\langle N_{PV} \rangle \sim 1$

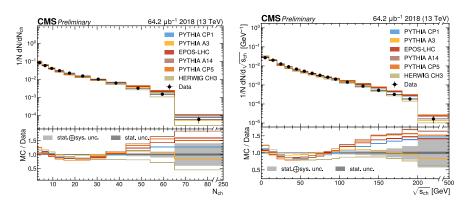
Tracks pT > 0.5 GeV, and $|\eta| < 2.4$, $N_{trk} \ge 2$ No clustering is done.

The unfolded observables are computed using tracks (detector level), or charged particles (particle level)

Results compared with PYTHIA (CP1, A3, A14, CP5), Herwig (CH3), EPOS (LHC), and HERWIG (CH3) tune.

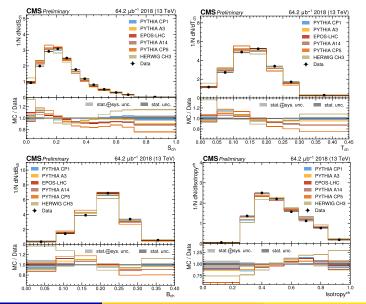
CMS-SMP-23-008-PAS



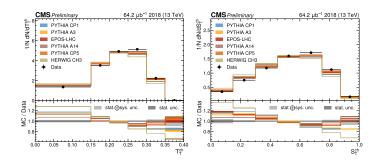


CMS-SMP-23-008-PAS









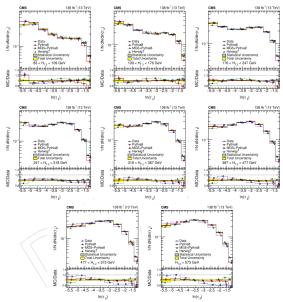


No MC generator shows good agreement.

- For $N_{ch} \lesssim 10$ and low $\sqrt{s_{ch}}$ all generators overestimate the fraction of events.
- For high N_{ch} and high $\sqrt{s_{ch}}$, the predictions of the MC generators diverge.
- For the
- Pythia 8 A3 tune (ATL-PHYS-PUB-2016-017) matches much better for all the ESVs.



CMS ESV with Tracks New Result



New CMS study of ESVs using charged particles in jets has been approved. (Suman Kumar Kundu)



String Fragmentation

Lund Model

$$f(z) = N \frac{(1-z)^a}{z} e^{-(bm^2/z)}$$

For heavy quarks $m \rightarrow m_T$

Lund-Bowler Model

$$f(z) = N \frac{(1-z)^a}{z^{1+r_q b m_q^2}} e^{-(bm_T^2/z)}$$

Z. Phys. C 11 (1981) 169

z fraction of the string momentum taken by the hadron,N normalization constant.

$$m_T = \sqrt{m^2 + p_T^2},$$

a, b parameters of the model m_a mass of the quark,

m mass of the hadron

 r_a shape parameter for the quark



String Fragmentation

- b-jets important for many analyses top quark, Higgs, new particle searches
- heaviest quark that hadronizes large m_b is important in the fragmentation
- MC event generator tuned using e⁺e⁻ data from LEP and SLC no QCD ISR/ UE – less complex
- LHC is a b-factory
 - high statistics data for understanding fragmentation of b-quarks
 - high luminosity pp data is complex PU, MPI,



CMS: b Jets in $t\bar{t}$ Events

CMS, $\sqrt{s} = 13 \text{ TeV}$, 35.9 fb^{-1} data

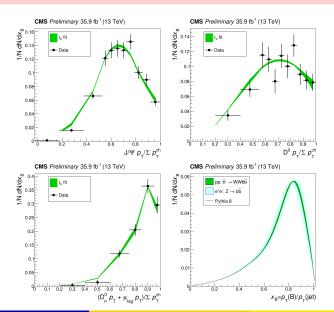
- $t\bar{t}$ events in single lepton (e/μ + jets), opposite sign dilepton ($\ell^+\ell^-$ + jets) channels (isolated high p_T e/μ).
- Charm meson candidates are identified with $J/\psi \to \mu^+\mu^-$ (both muons are non-isolated) $D^0 \to K^\pm\pi^\mp$ opposite sign charged tracks
- Events with candiate D^0 are divided into two sub-samples with soft μ tag and without.
- proxy x_b is calculated for each charm meson,

$$x_b = \frac{p_T(J/\psi, D^0)}{\sum_{jet} p_T^{ch}}$$

the sum includes all tracks with $p_T > 0.5$ GeV in the same jet.

• $t\bar{t}$ event simulation – NLO (POWHEG V2), $m_t=172.5~\text{GeV}+\text{PYTHIA 8}$ (CUETP8M2T4), Lund-Bowler fragmentation

CMS: b Jets in $t\bar{t}$ Events





CMS: b Jets in $t\bar{t}$ Events

- template fit to x_b distributions for each of the meson samples $r_b = 0.858 \pm 0.037 \ (st) \pm 0.031 \ (sys)$
- f(z) does not show any environment dependence e^+e^- , pp.

LHC can compete with LEP in precision measurements.



ATLAS: b Jets in tt Events

$t\bar{t} \rightarrow e\mu\nu\nu b\bar{b}$ events:

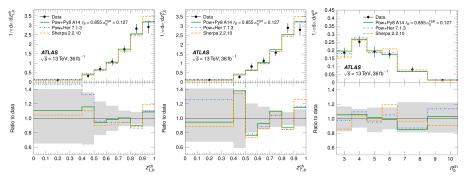
- reconstructed jets = 2, b-tagged jet ≥ 1
- exactly two leptons with opposite sign $e^{\pm}\mu^{\mp}$
- relatively low contamination from light-flavor jets.
- $e\mu$ channel particularly small contributions from non- $t\bar{t}$ processes.

Charged particles with $p_T > 500 \text{ MeV}$ are used to form $\vec{p}_b^{ch} = \sum_{b-hadron} \vec{p}^{ch}, \ \vec{p}_{iet}^{ch} = \sum_{iet} \vec{p}^{ch}$

$$z_{T,b}^{ch} = \frac{p_{T,b}^{ch}}{p_{T,jet}^{ch}}, \ \ z_{L,b}^{ch} = \frac{\vec{p}_b^{ch} \cdot \vec{p}_{jet}^{ch}}{|\vec{p}_{jet}^{ch}|^2}, \ \ \rho = \frac{p_{T,b}^{ch}}{p_T^e + p_T^\mu}, \ \ n_b^{ch}$$



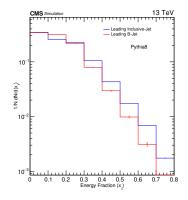
ATLAS: b Jets in $t\bar{t}$ Events



Predictions of MC generators tuned to a combination of lepton and hadron collider measurements show improved agreement with the observed LHC data.



Energy fraction in b Jets



Energy fraction $x_1 = p_{T,1}/p_T^{jet}$ in the leading jet (b-jet) of an event.



Jet Charge Observables

Three jet charge observables are studied – Default, longitudinal and transverse

$$\begin{aligned} Q_{J}^{\kappa} &=& \frac{1}{\left(\boldsymbol{p}_{T}^{jet}\right)^{\kappa}} \sum_{i} q_{i} \left(\boldsymbol{p}_{T}^{i}\right)^{\kappa} \\ Q_{J,L}^{\kappa} &=& \frac{\sum_{i} q_{i} \left(\boldsymbol{p}_{\parallel}^{i}\right)^{\kappa}}{\sum_{i} \left(\boldsymbol{p}_{\parallel}^{i}\right)^{\kappa}}, \ \boldsymbol{p}_{\parallel}^{i} = \frac{\vec{p}^{i} \cdot \vec{p}_{jet}}{|\vec{p}_{jet}|} \\ Q_{J,T}^{\kappa} &=& \frac{\sum_{i} q_{i} \left(\boldsymbol{p}_{\perp}^{i}\right)^{\kappa}}{\sum_{i} \left(\boldsymbol{p}_{\perp}^{i}\right)^{\kappa}}, \ \boldsymbol{p}_{\perp}^{i} = \frac{\vec{p}^{i} \times \vec{p}_{jet}}{\vec{p}_{jet}|} \end{aligned}$$

The parameter κ controls the relative weightage given to the low energy particles.



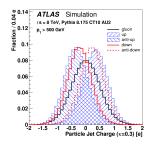
JCO Measurements

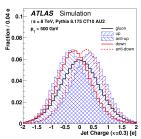
- Distributions of the JCOs are broad, overlapping \Rightarrow discriminating partonic origin $(q/\bar{q}/g)$ of a jet is not possible.
- D0 used $Q_J^{0.6}$ and later CDF used $Q_{J,L}^{0.5}$ to confirm q_t and exclude exotic quarks.

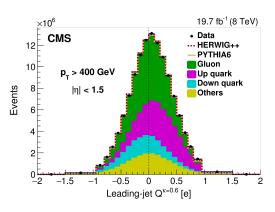
 PRL 98 (2007) 041801, PRD 88 (2013) 032003
- ATLAS studied $Q_{J,L}^{0.5}$ for b-jets in $t\bar{t}$ events to estimate q_t . CMS also studied $Q_{J,L}^{0.7}$
- JCOs not affected by PU ⇒ suitable for the HL-LHC.



JCO Measurements



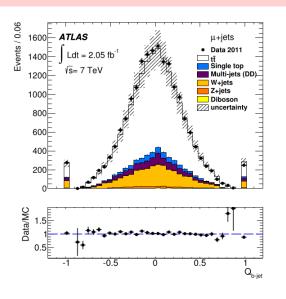


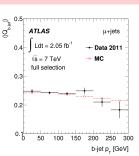


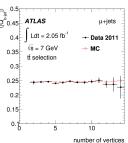
Phys. Rev. D 93 (2016), JHEP 10(2017) 131



ATLAS: JCO Measurements



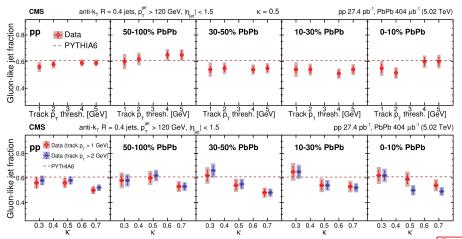






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CMS: JCO Measurements



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CMS: JCO Measurements

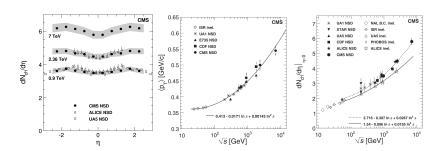
A new CMS study of JCOs is ready for approval process.



- color reconnection
- Power law behaviour negative binomial distribution (NBD) double NBD
- KNO scaling $\Psi(z) = \langle n \rangle P_n, \ z = n/\langle n \rangle$



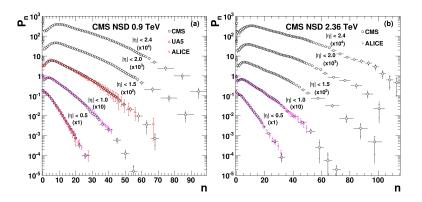
Charged Particles



$$\langle p_T
angle = 0.545 \pm 0.005 \pm 0.15$$
 GeV/c

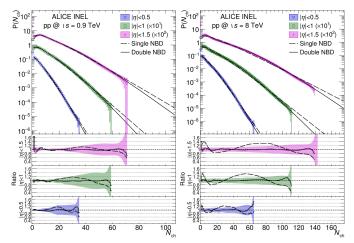
CMS, PRL 105, 022002 (2010)





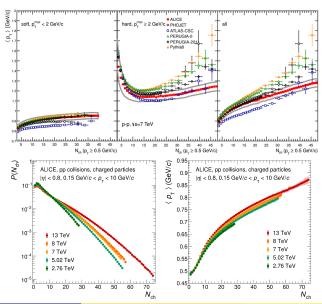
JHEP 01 (2011) 079





Single NBD fails to describe data at higher multiplicities.







Outlook

- Many measurements have been done to understand QCD at low energies - event and jet shapes, charged particle production, ...
- The measurements focus on different aspects of particle production in hadron (and nuclear) collisions
- A huge jigsaw puzzle with many pieces
- A combined and exhaustive picture of the process of particle production is yet to emerge ...

... and we keep colliding, particles and ideas

